

The Hubble Tension as the Signature of a New Phase Transition



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Nordita

in collaboration with:

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and

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and based on earlier work with:

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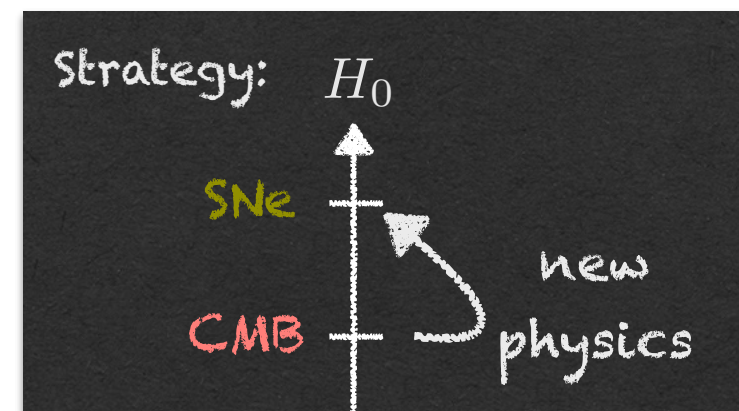
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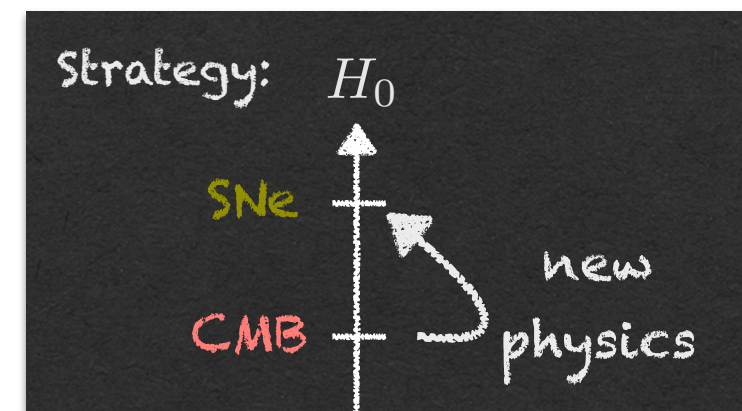
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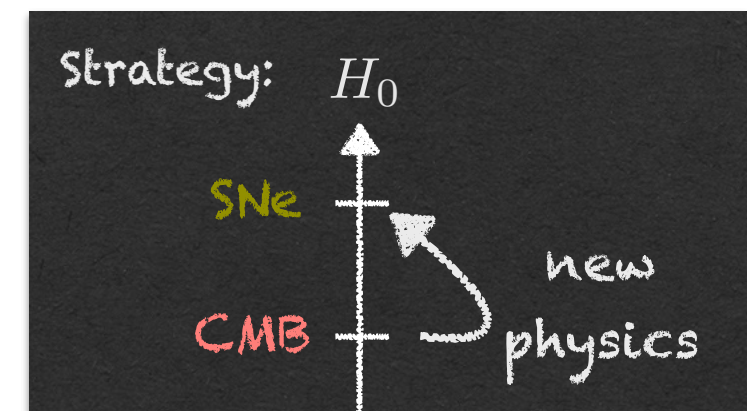
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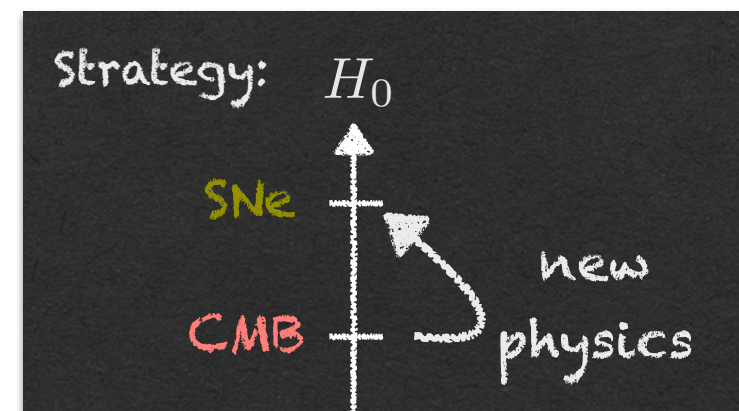
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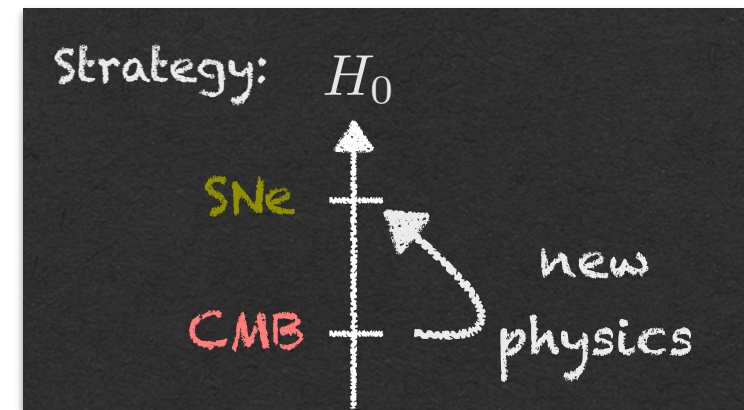
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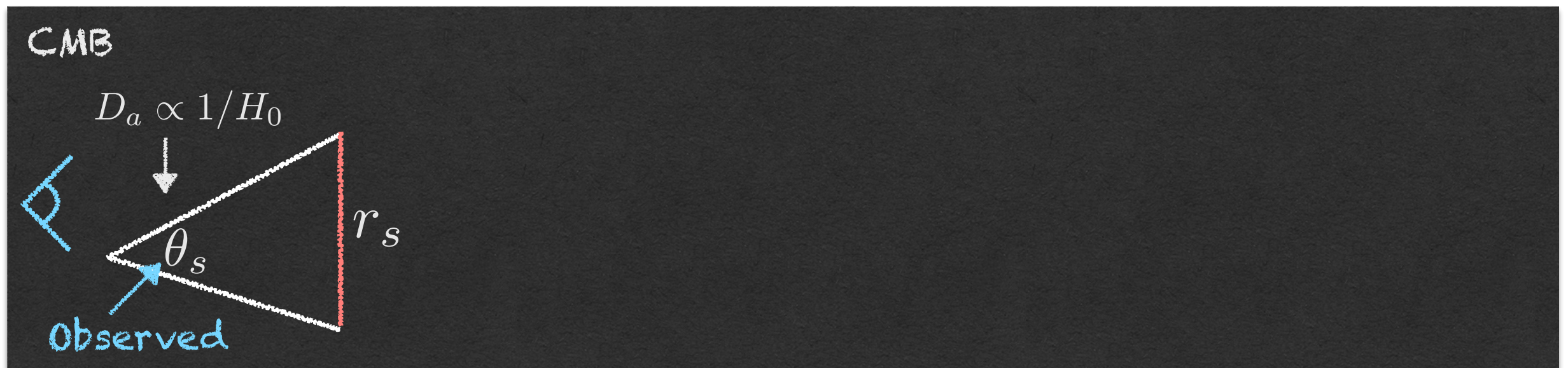
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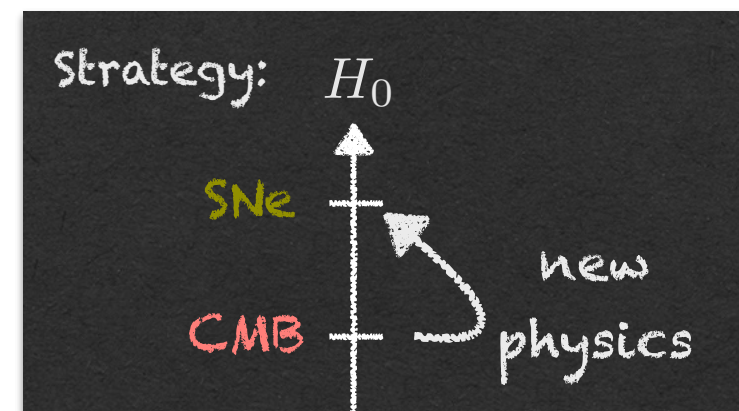
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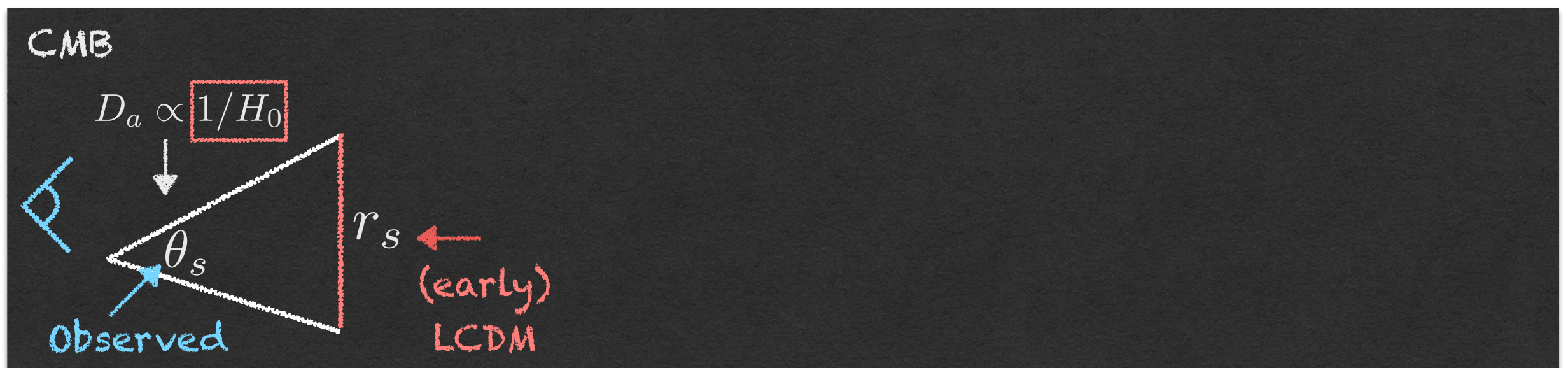
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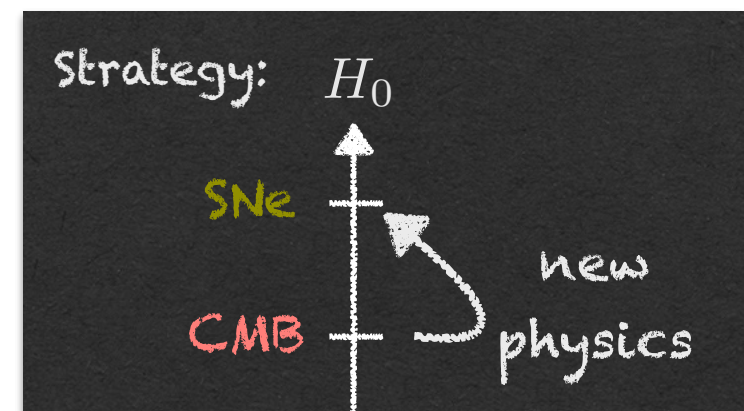
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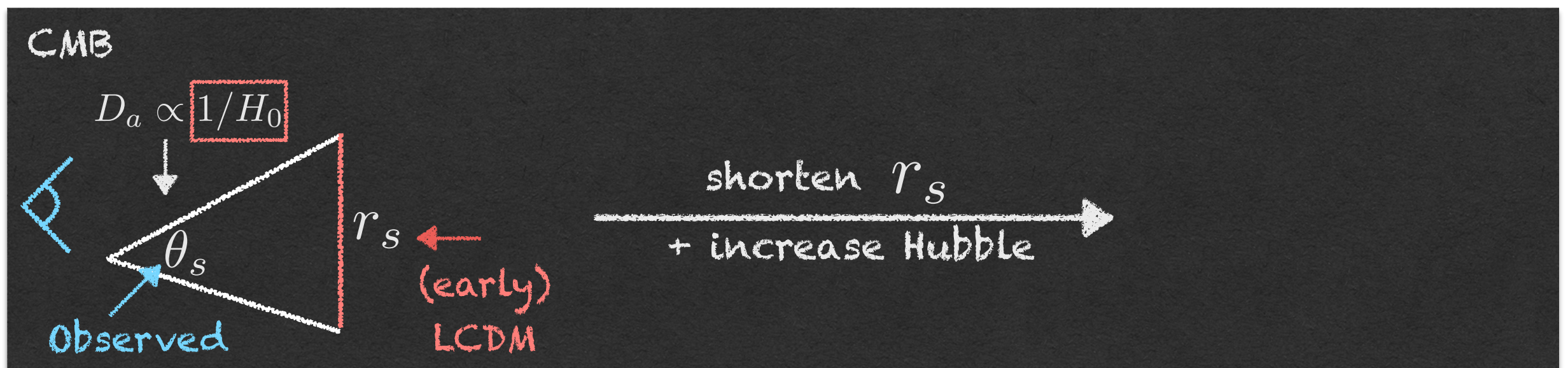
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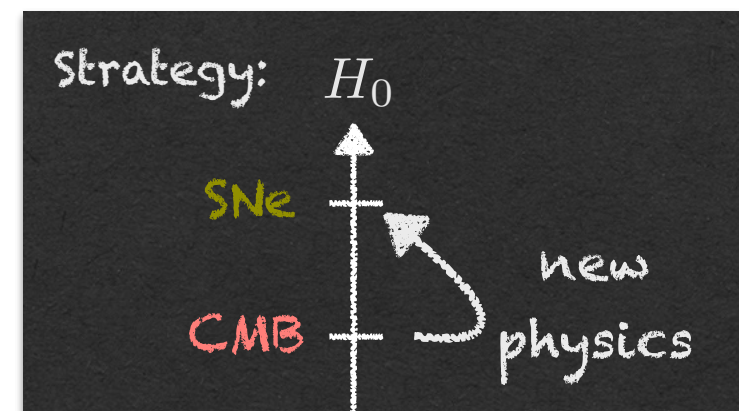
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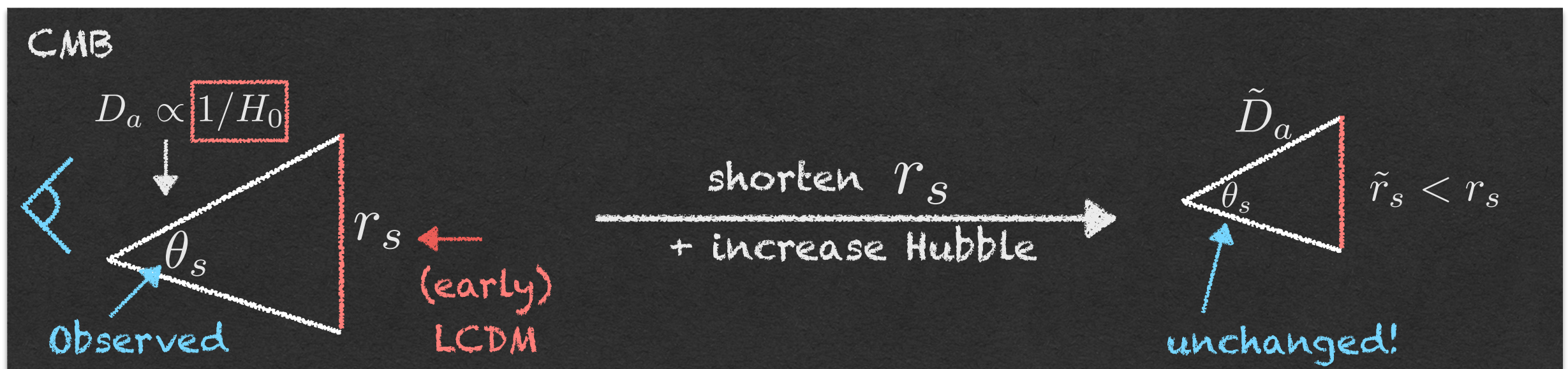
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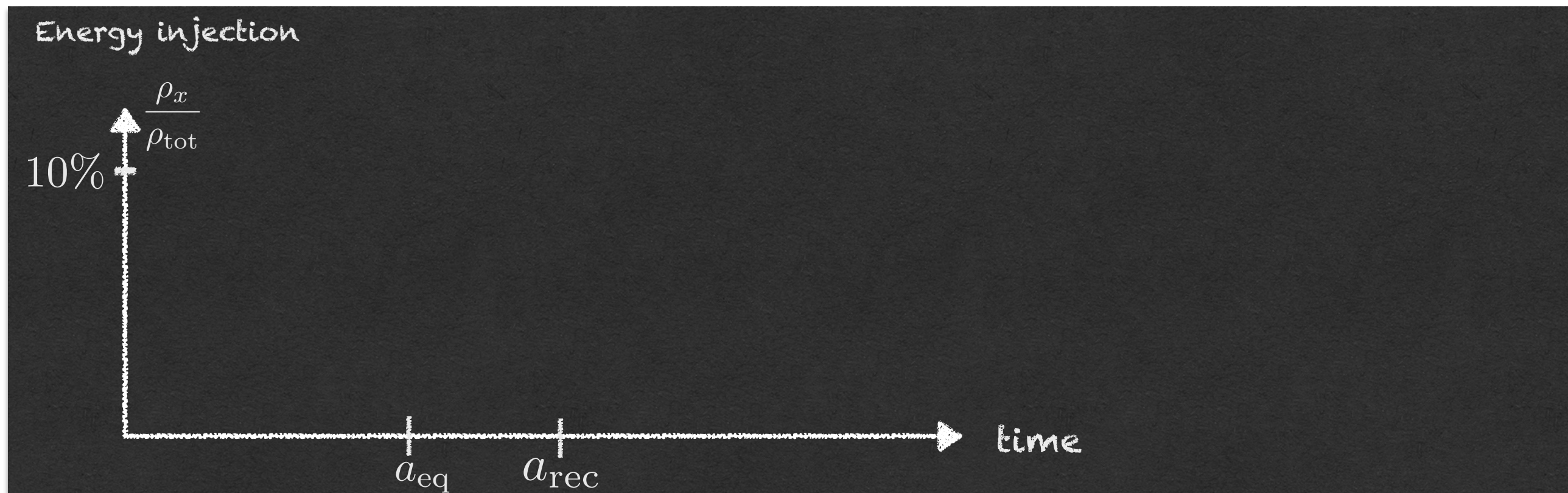
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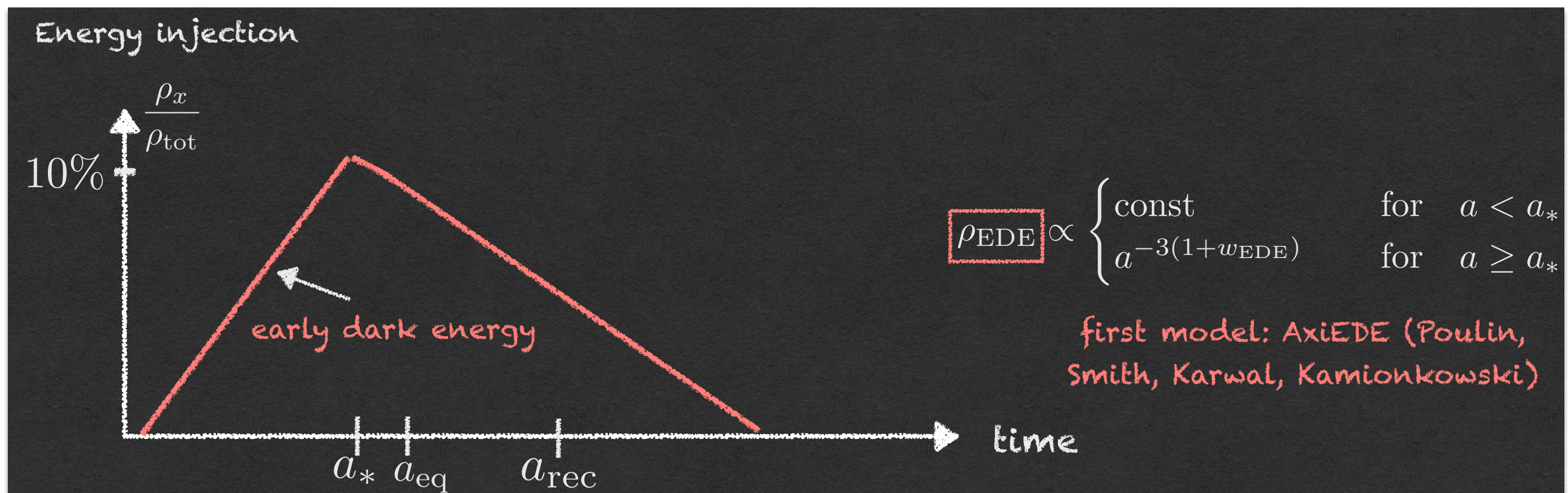
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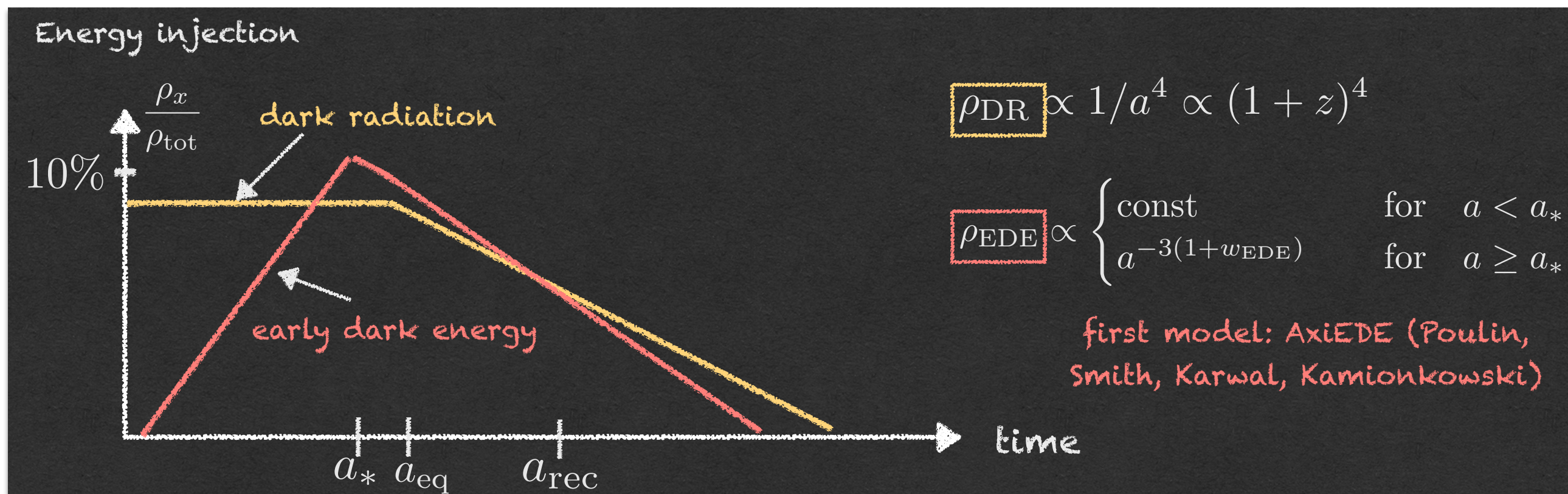
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► **DR:** Provides **continuous** injection before matter–radiation equality — Hot NEDE.

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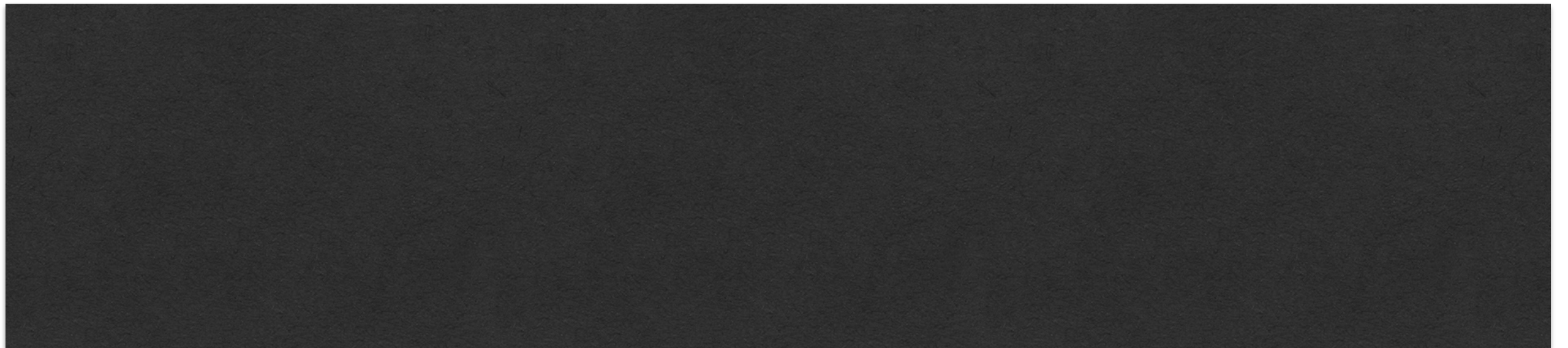
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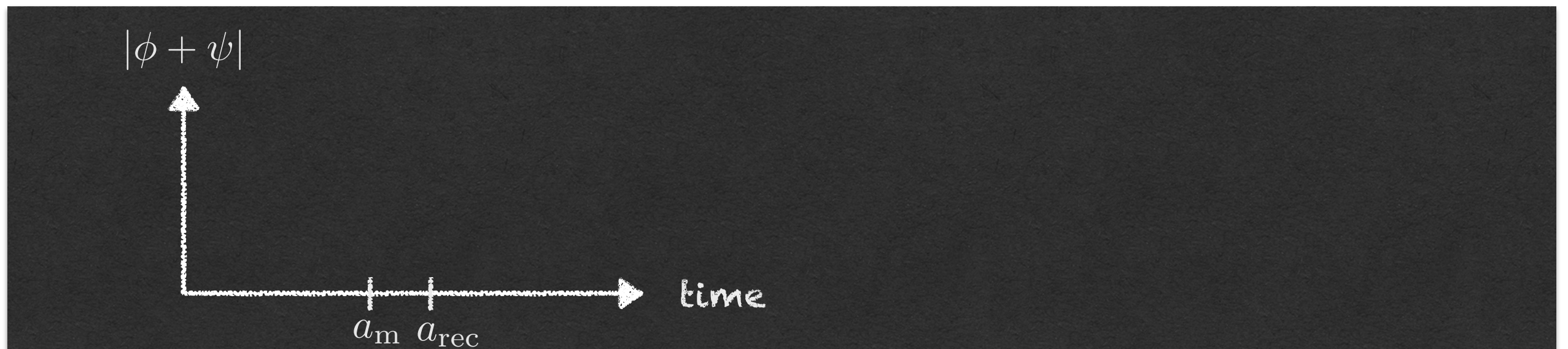
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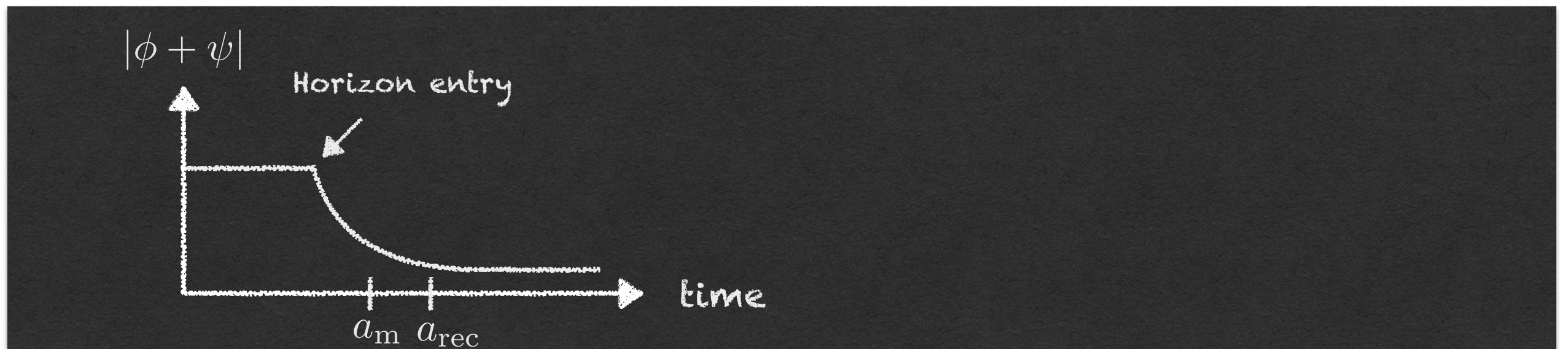
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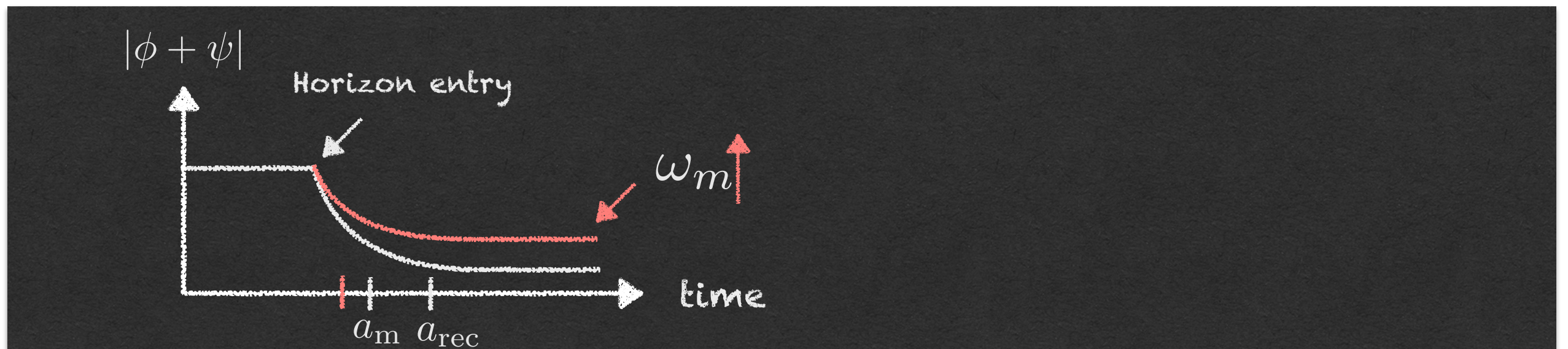
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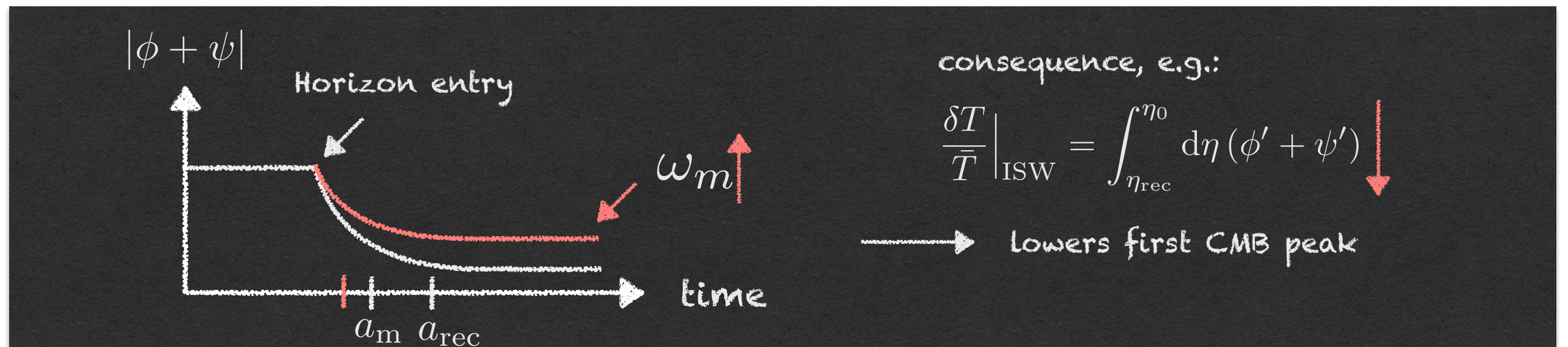
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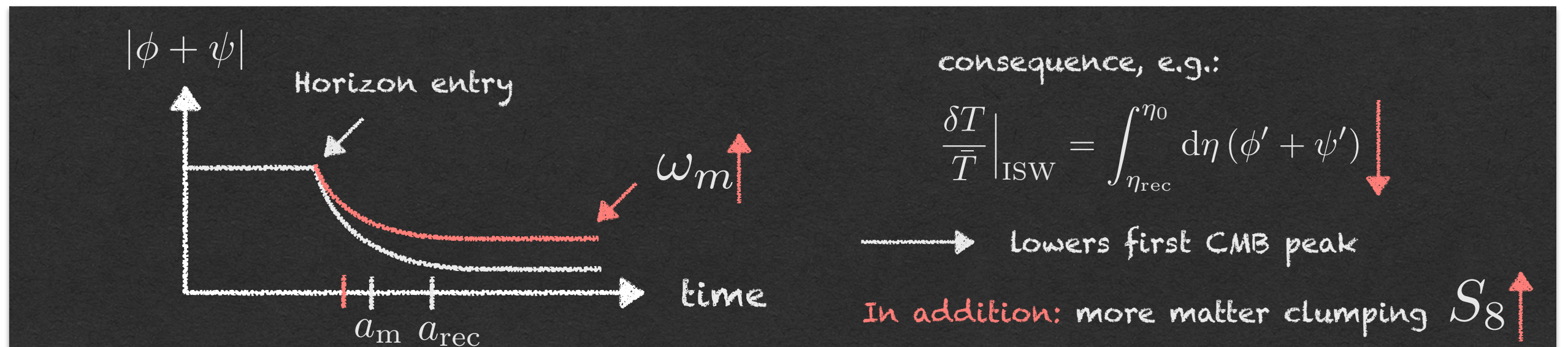
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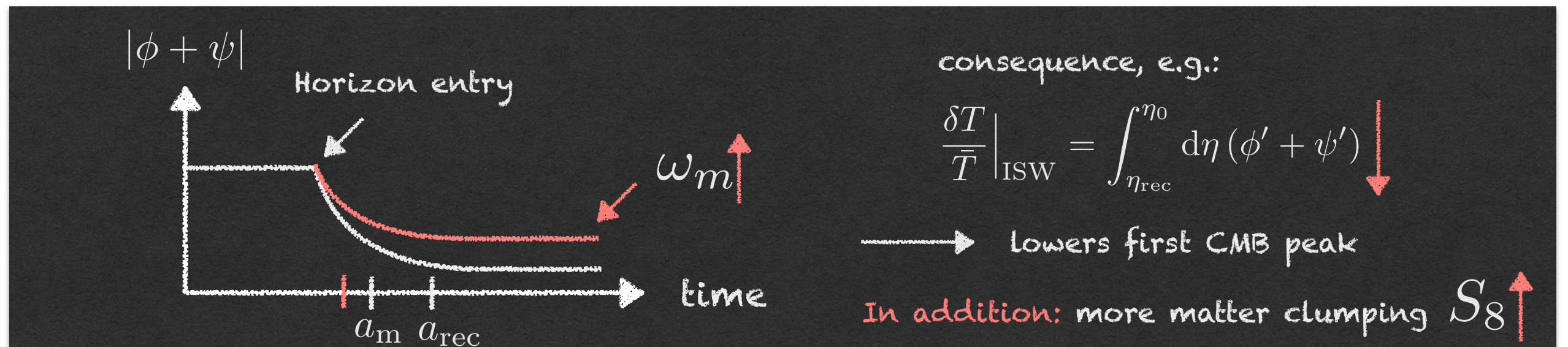
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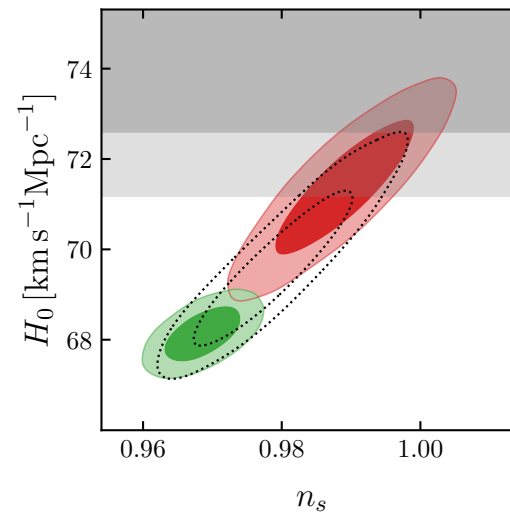


- Any early-time model needs to come with compensation mechanism, **examples:**
 - (i) delay matter domination, (ii) dark sector acoustic oscillations, (iii) fuzzy dark matter
- **Important lesson:** The detailed compensation mechanism requires a specific model!

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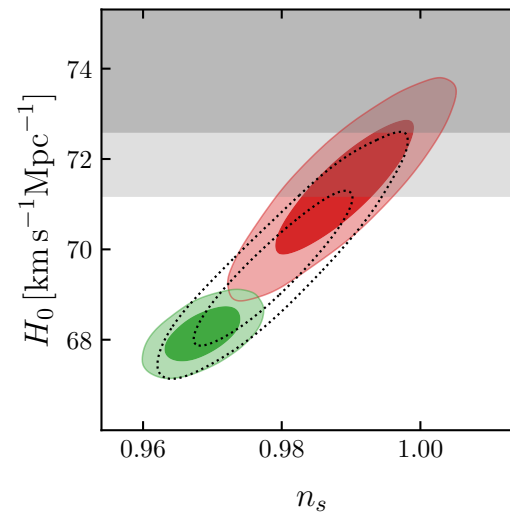


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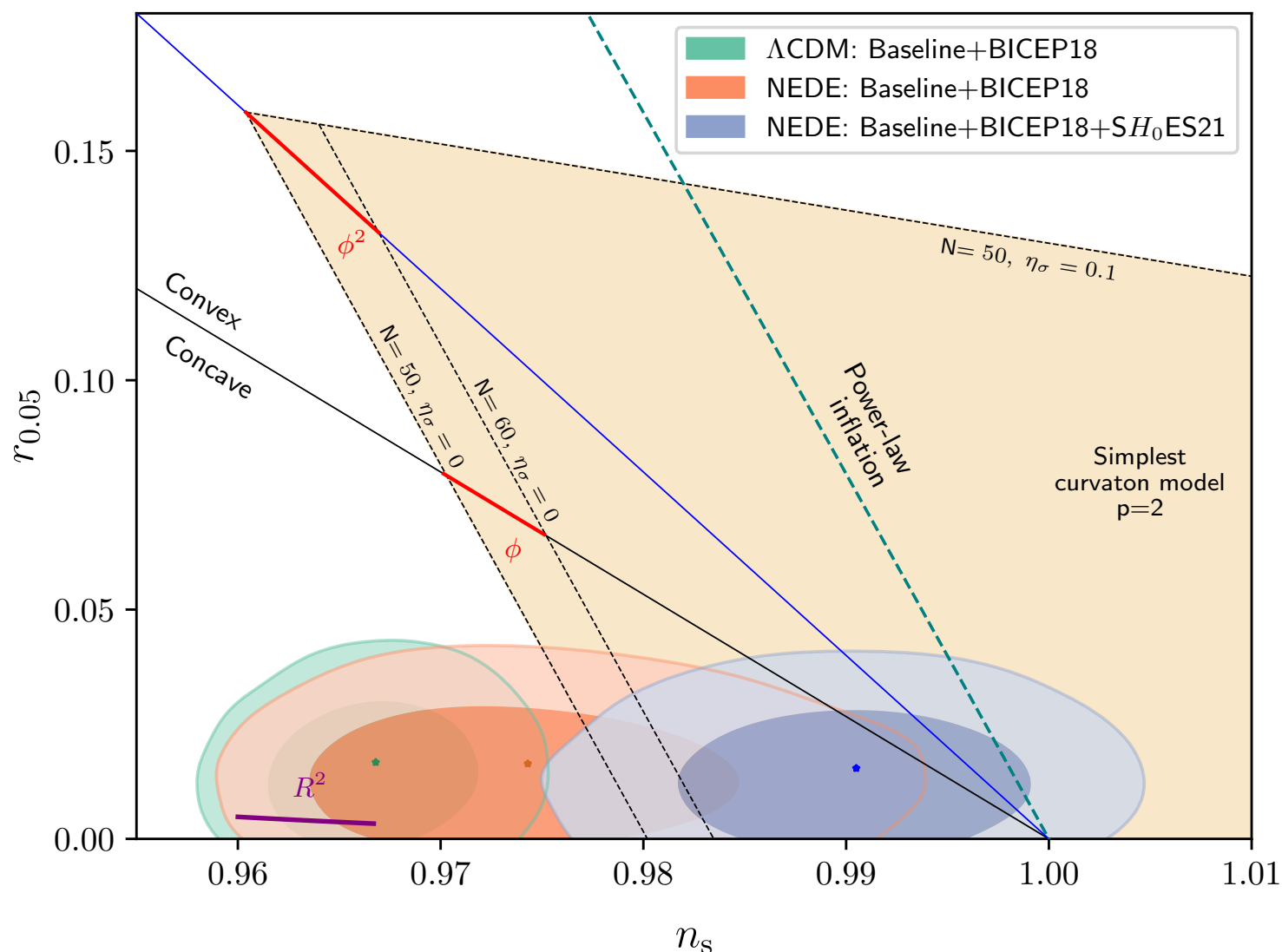
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- ◆ Could bring back to life simple models of inflation, e.g.:
 - quadratic potential + curvaton
 - power-law inflation (exp. potentials)
- ◆ **For now:** Keep in mind LCDM dependence of primordial constraints.

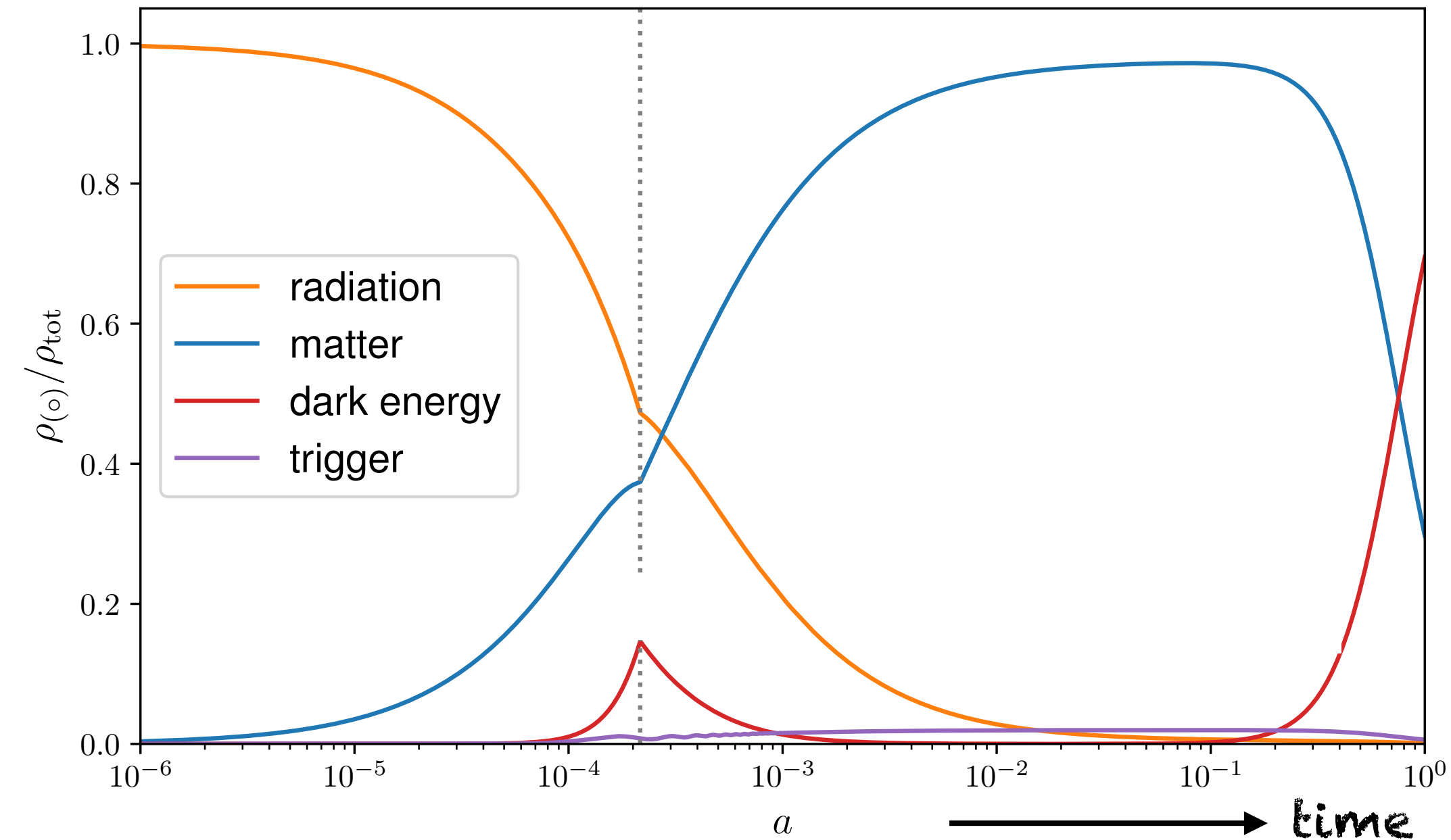
Part A: Early Dark Energy

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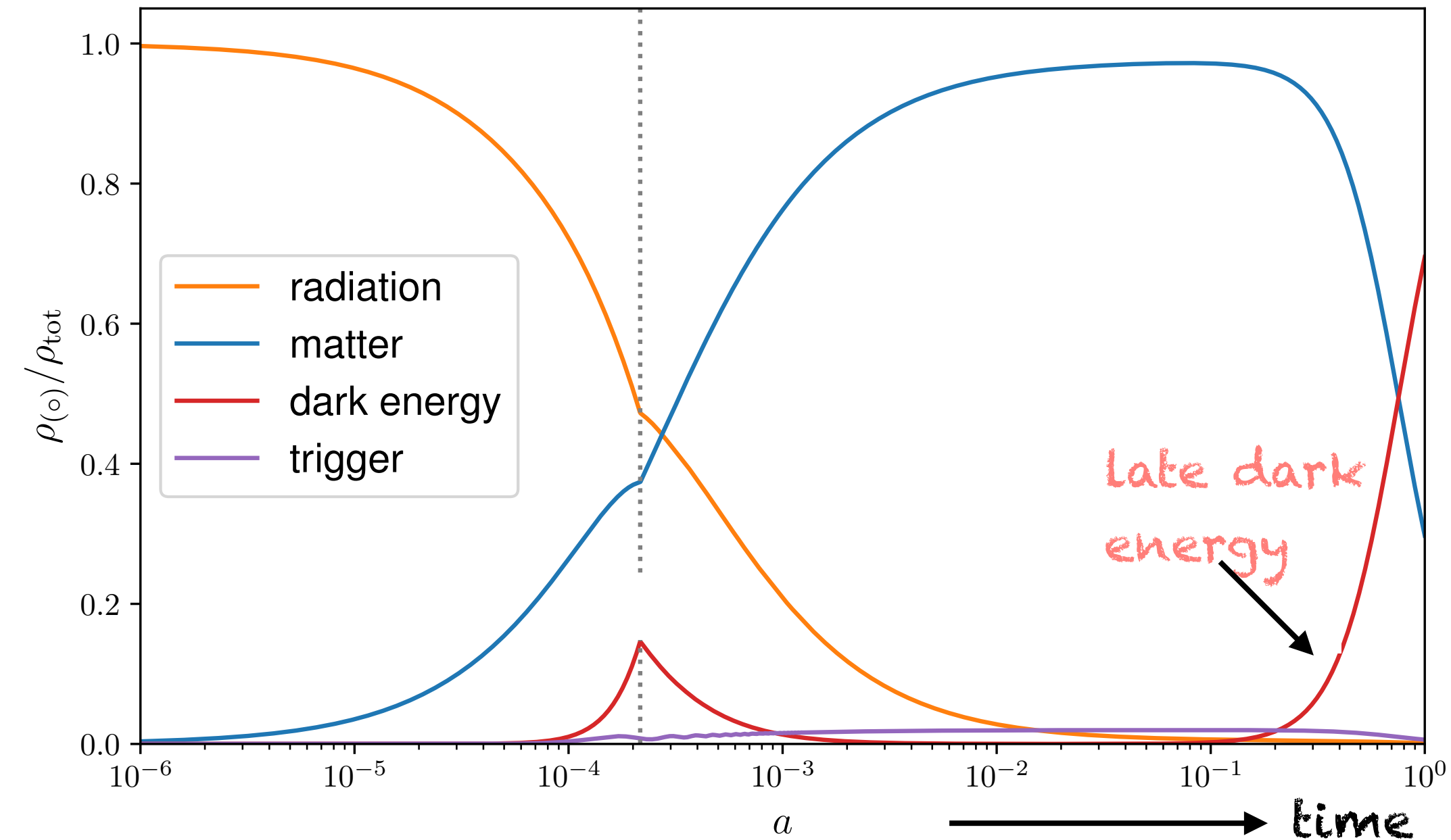
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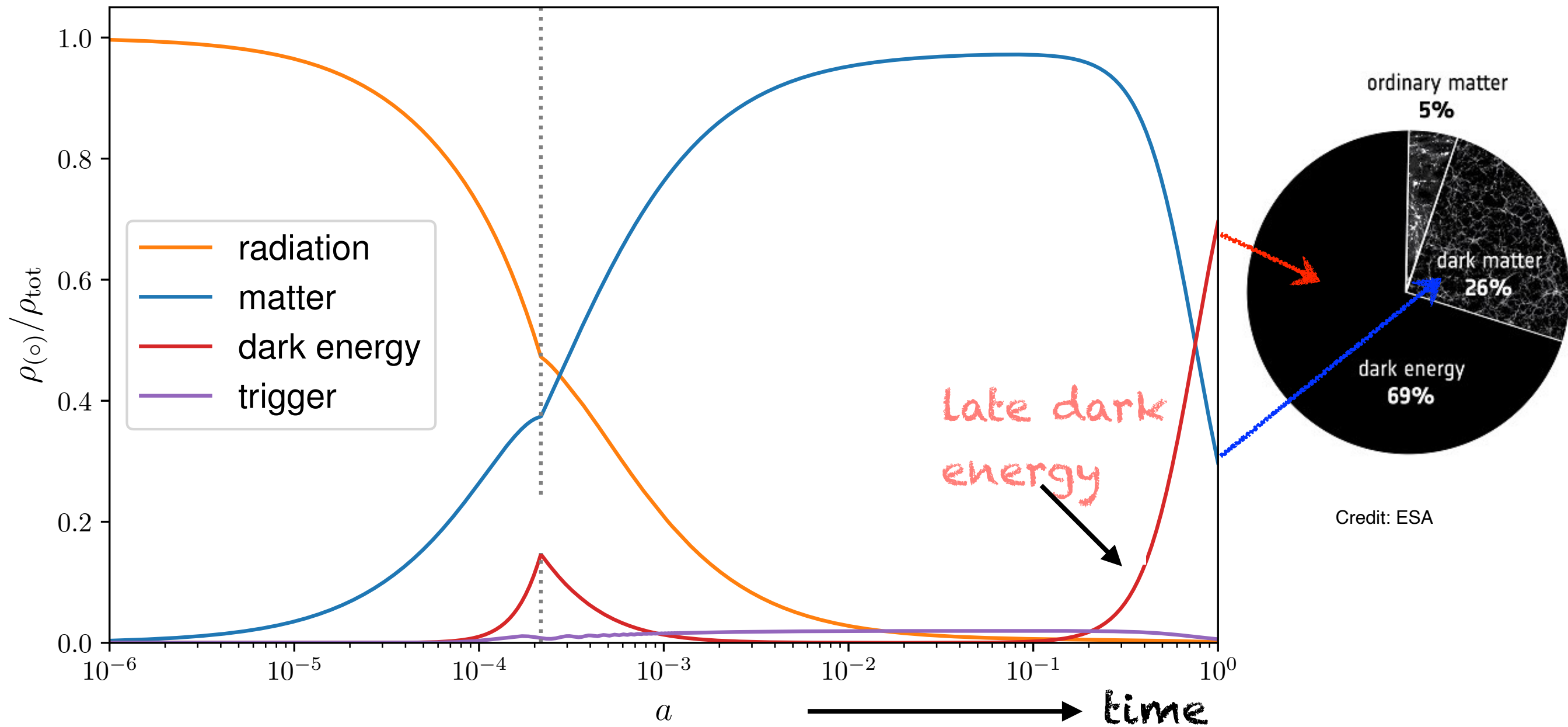
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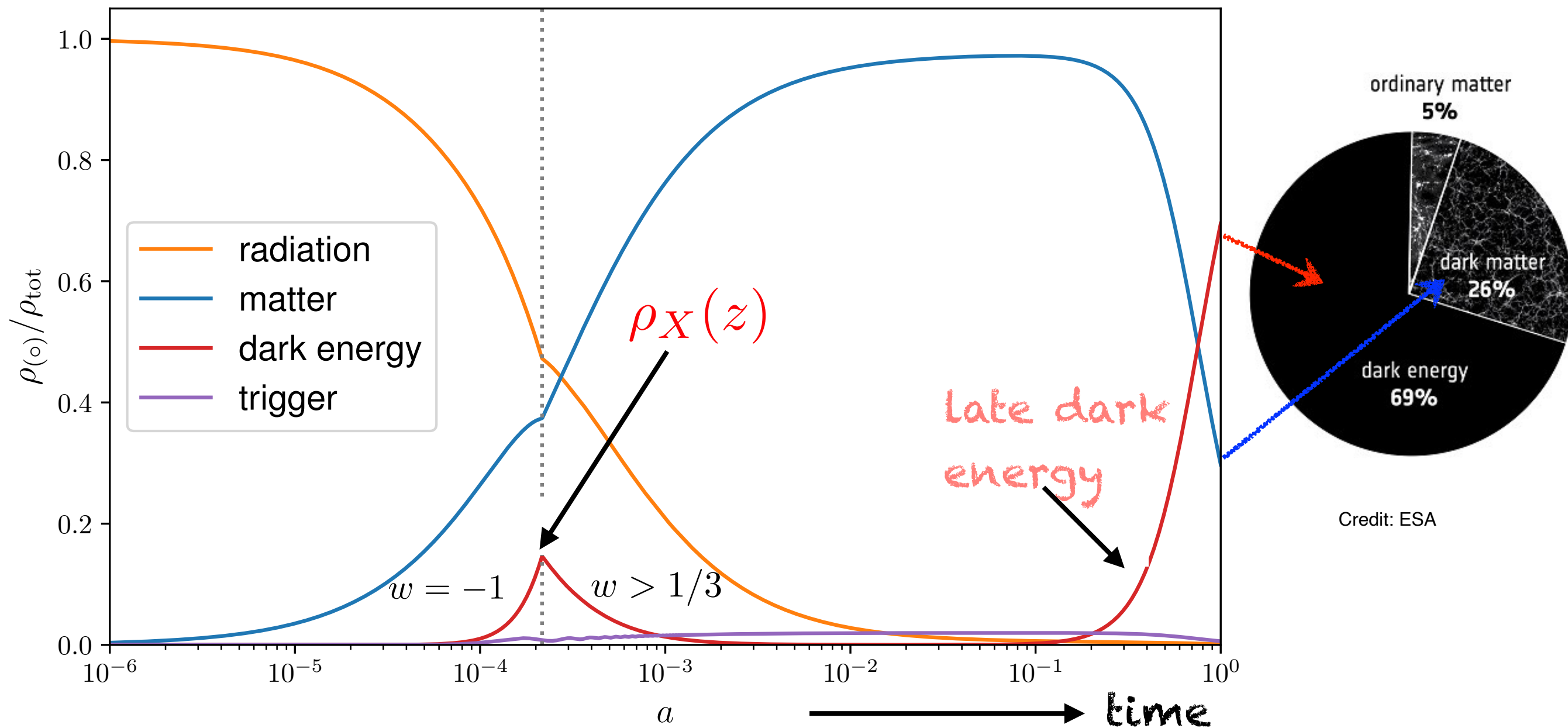
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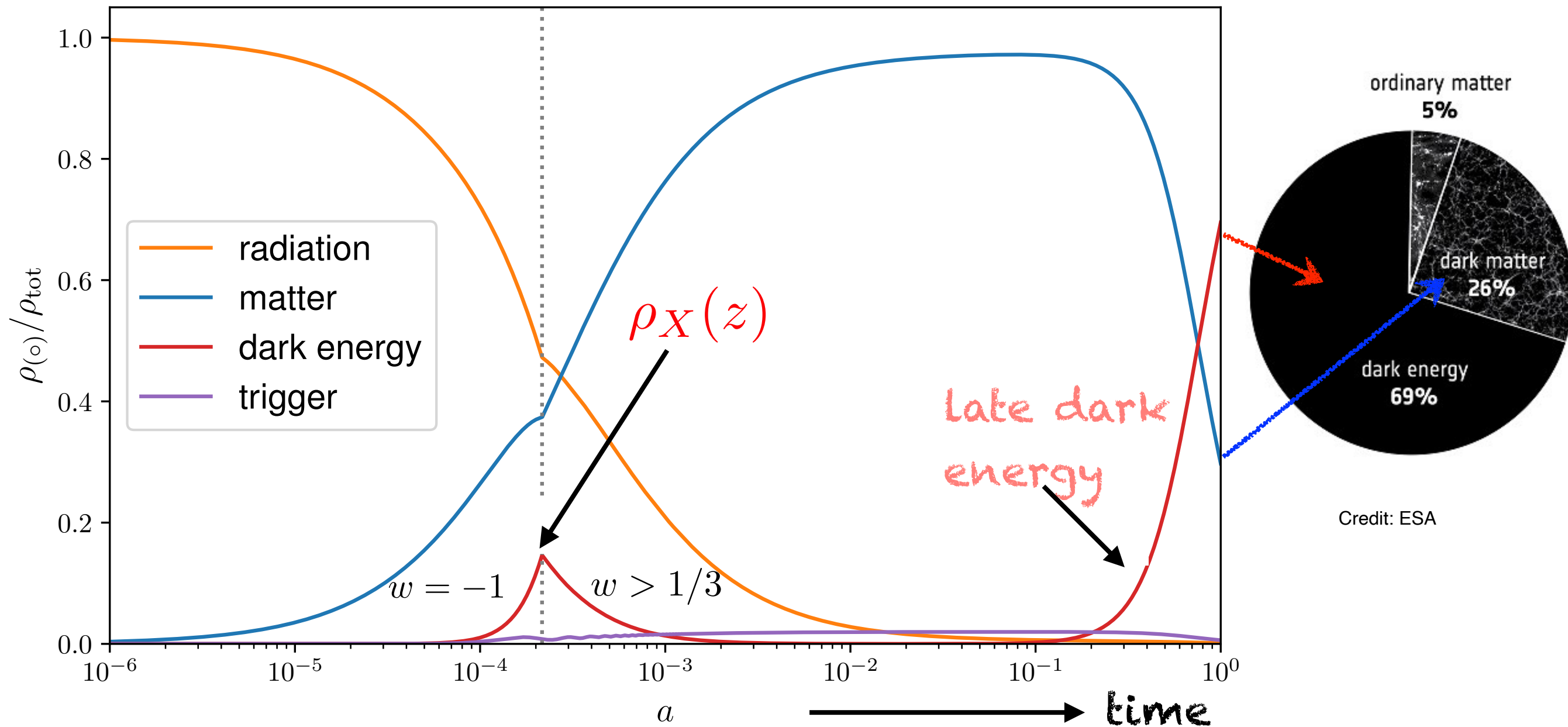
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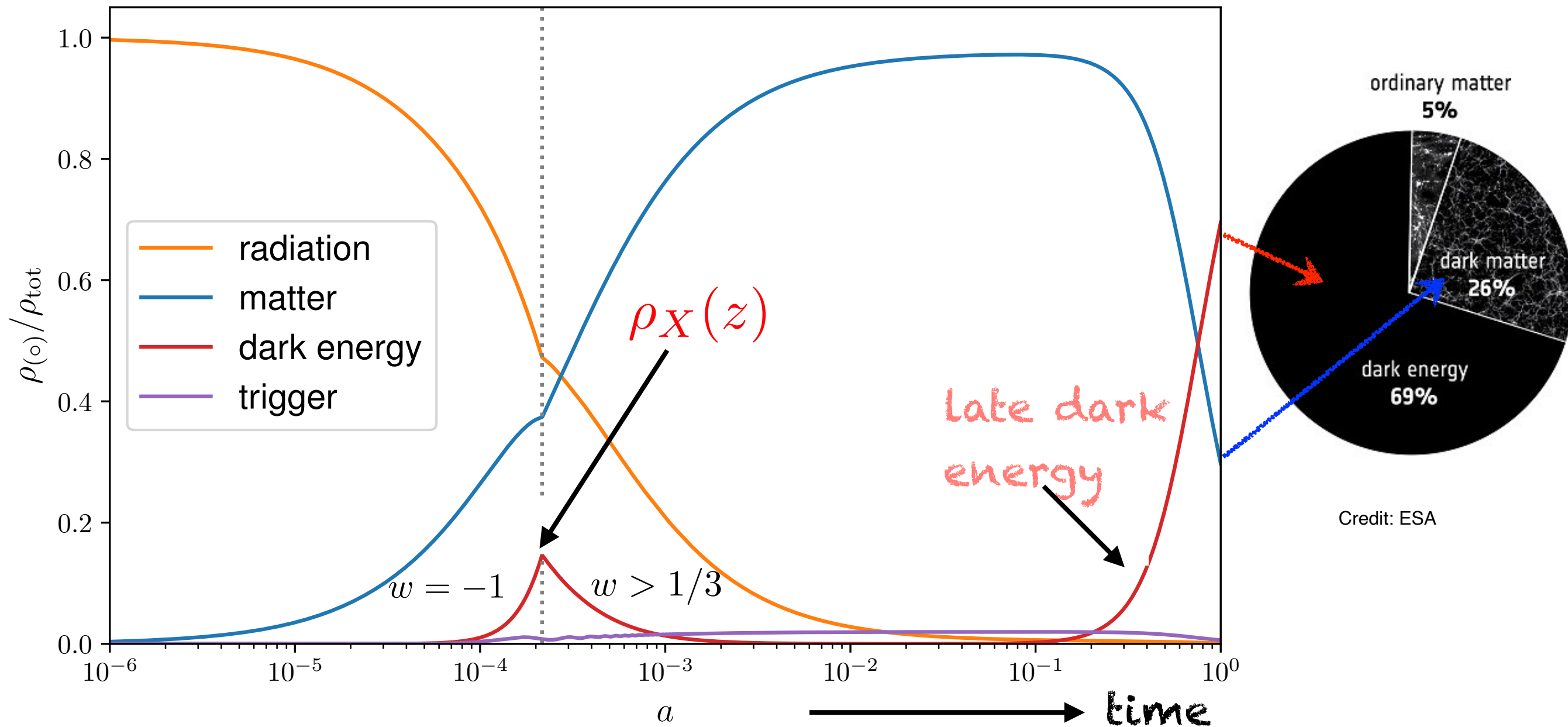
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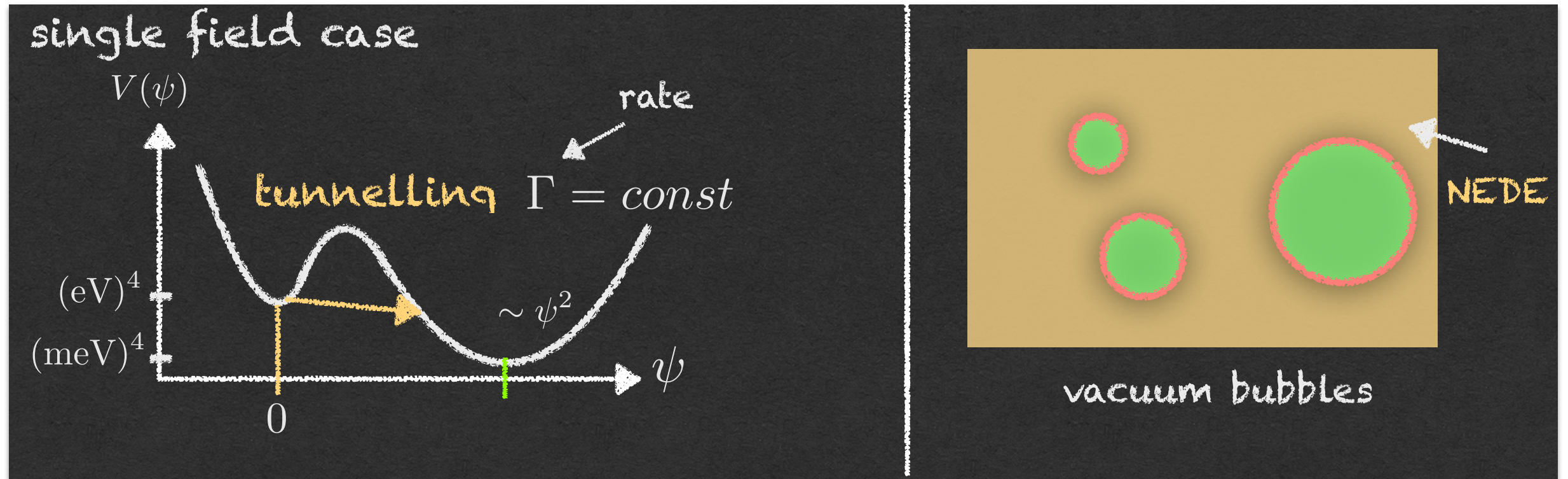


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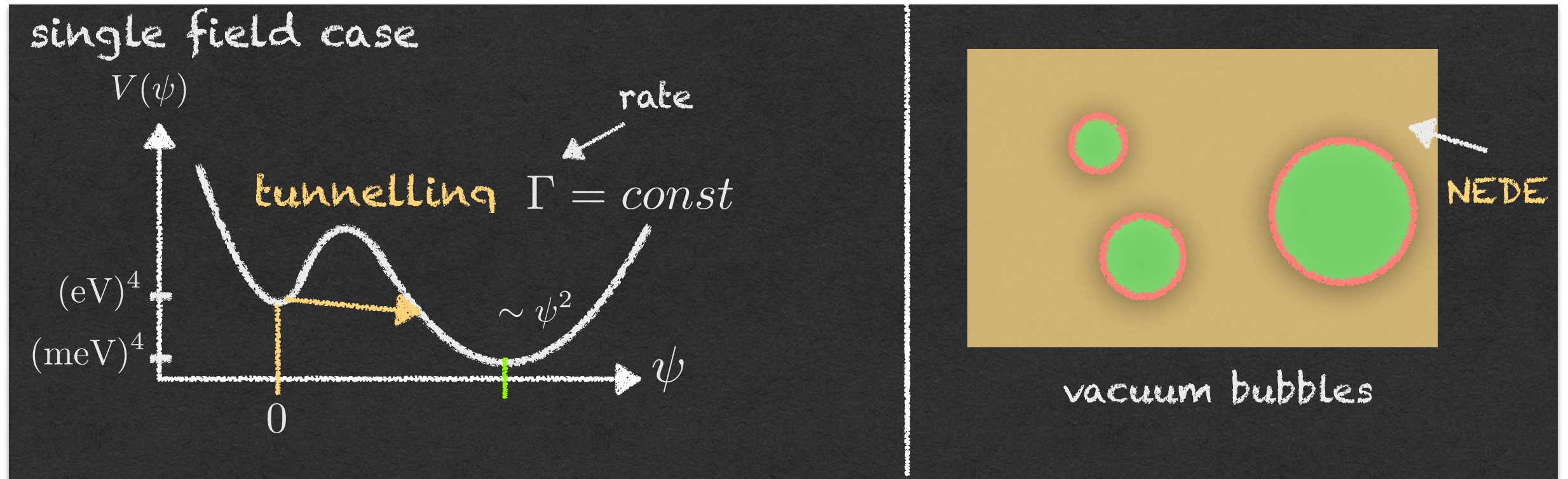
► Insight with Martin S. Sloth 2019: Looks like a **vacuum phase transition!**

Vacuum phase transition

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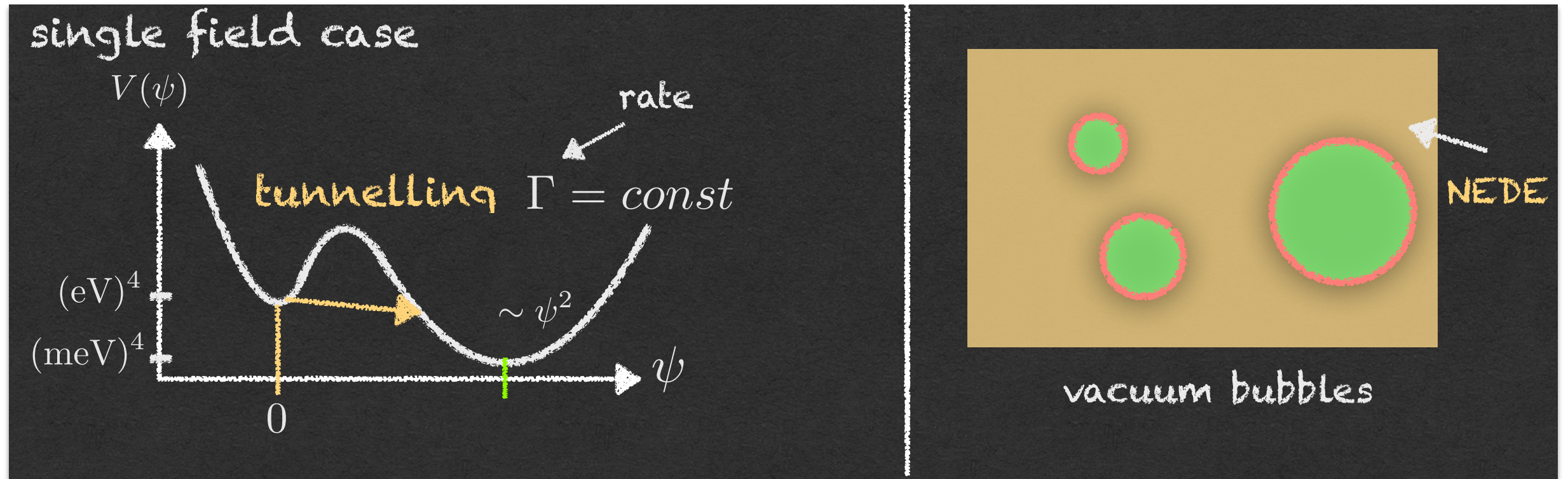


Vacuum phase transition



- **Hubble tension:** EDE provided by (decaying) false vacuum energy / latent heat.

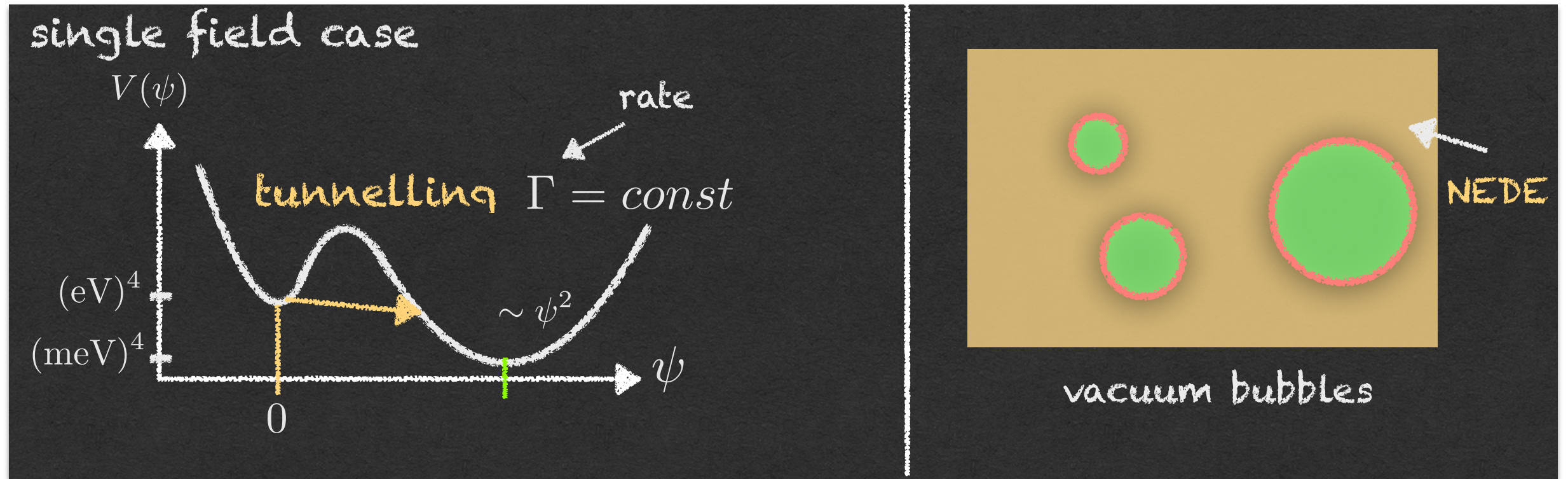
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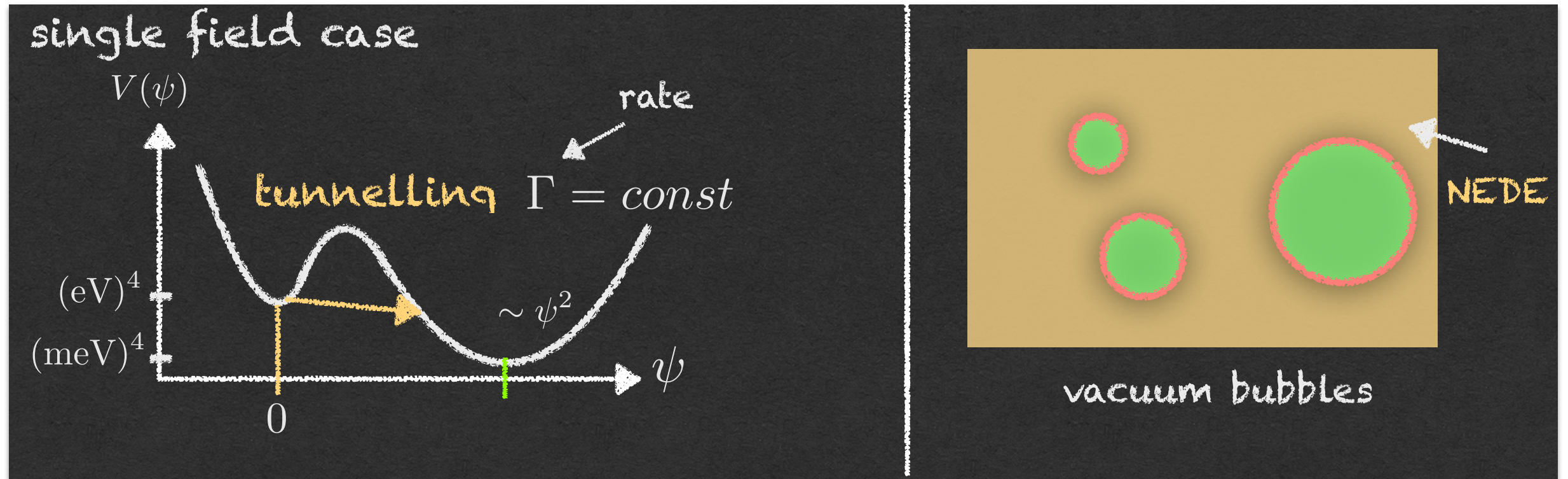


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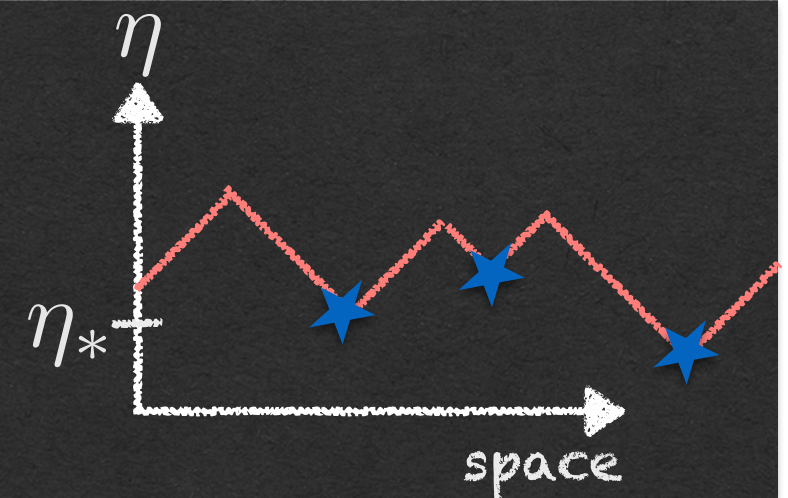
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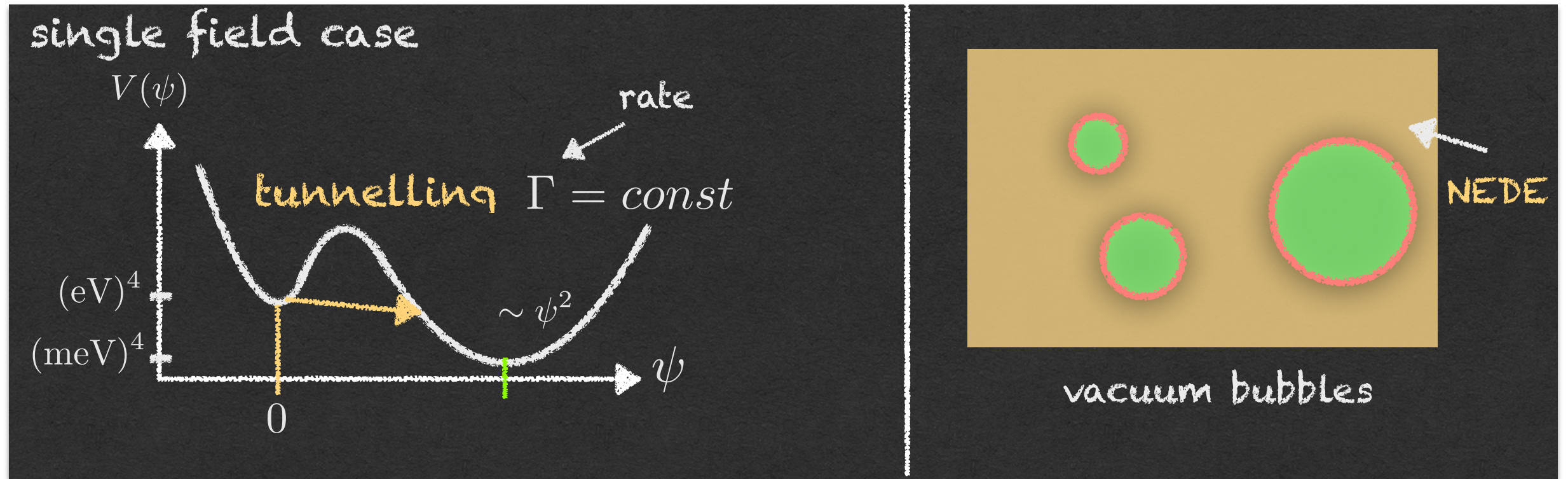
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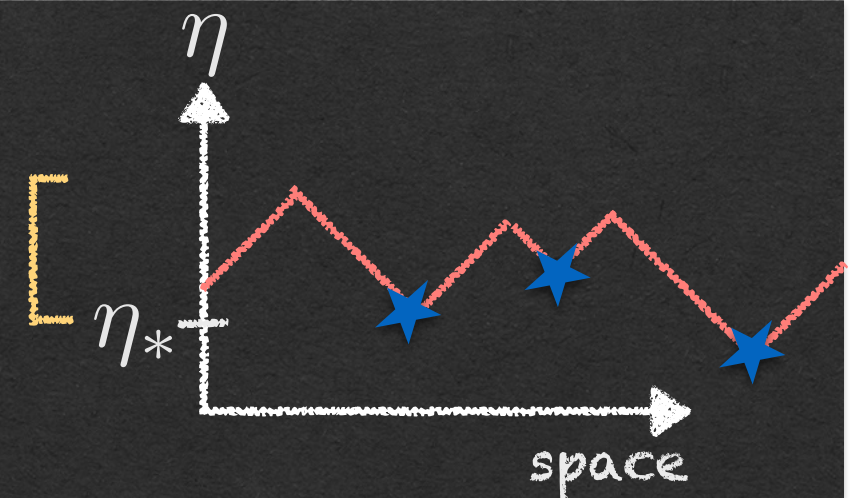


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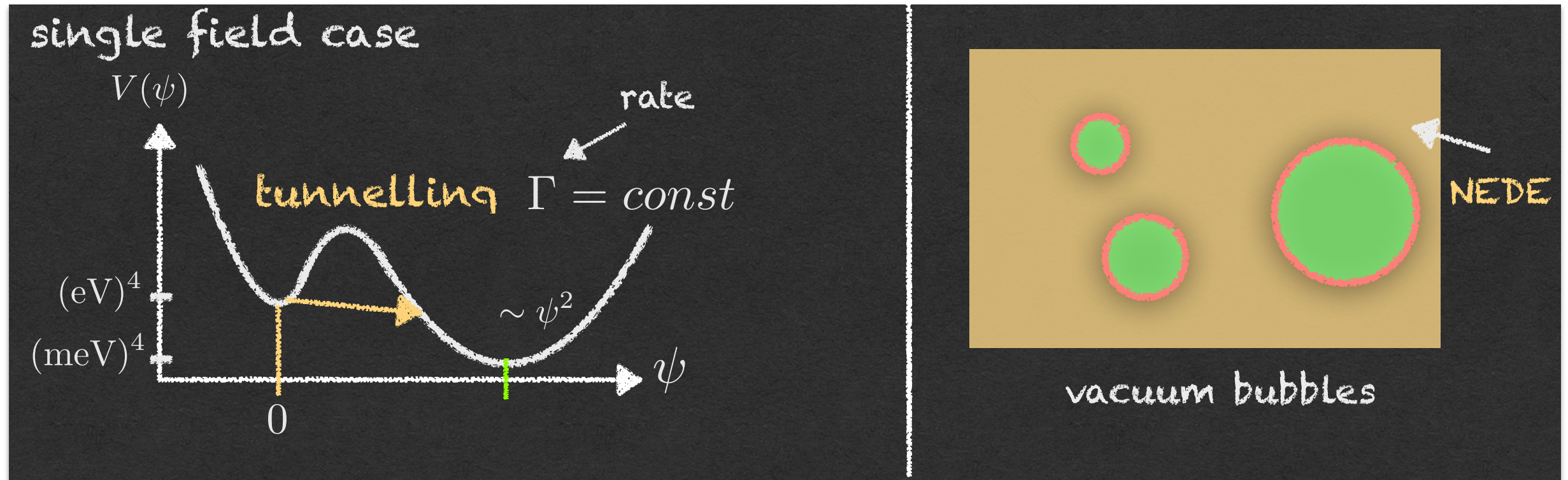
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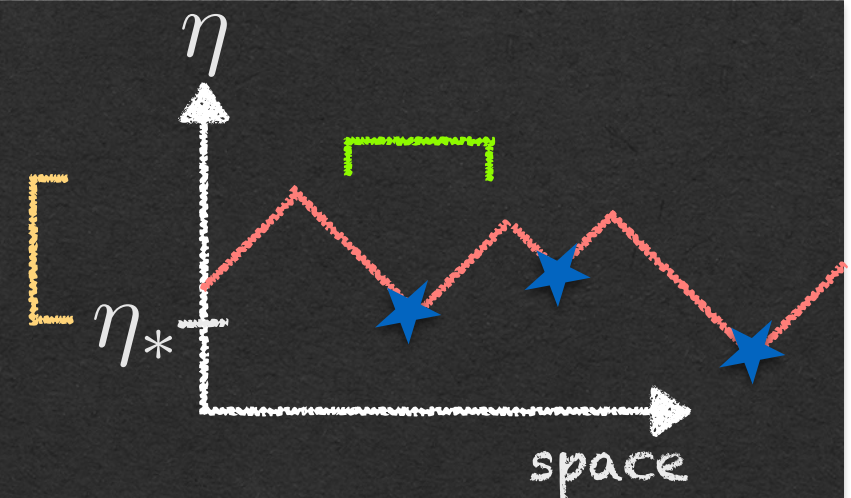
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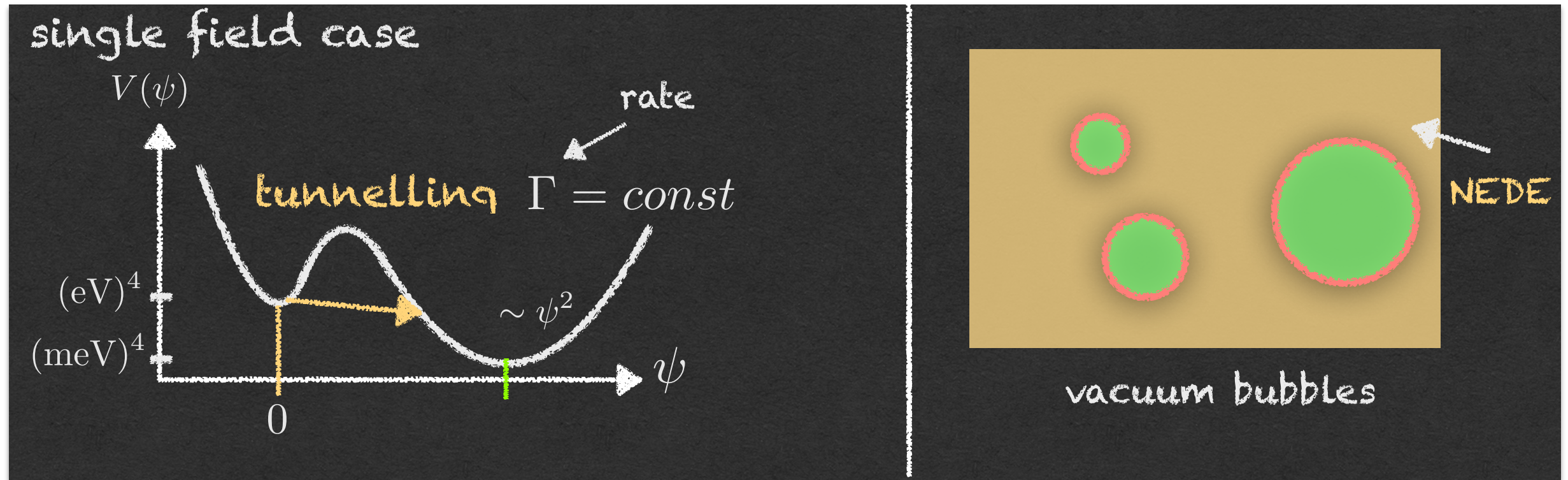
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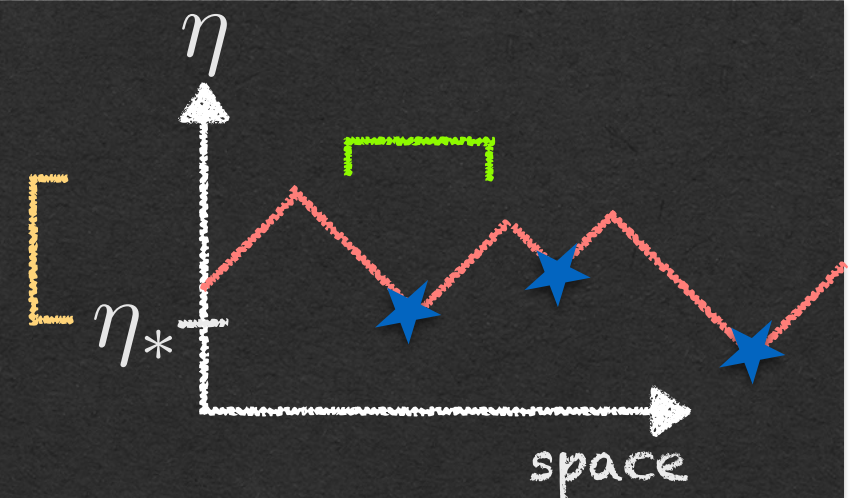
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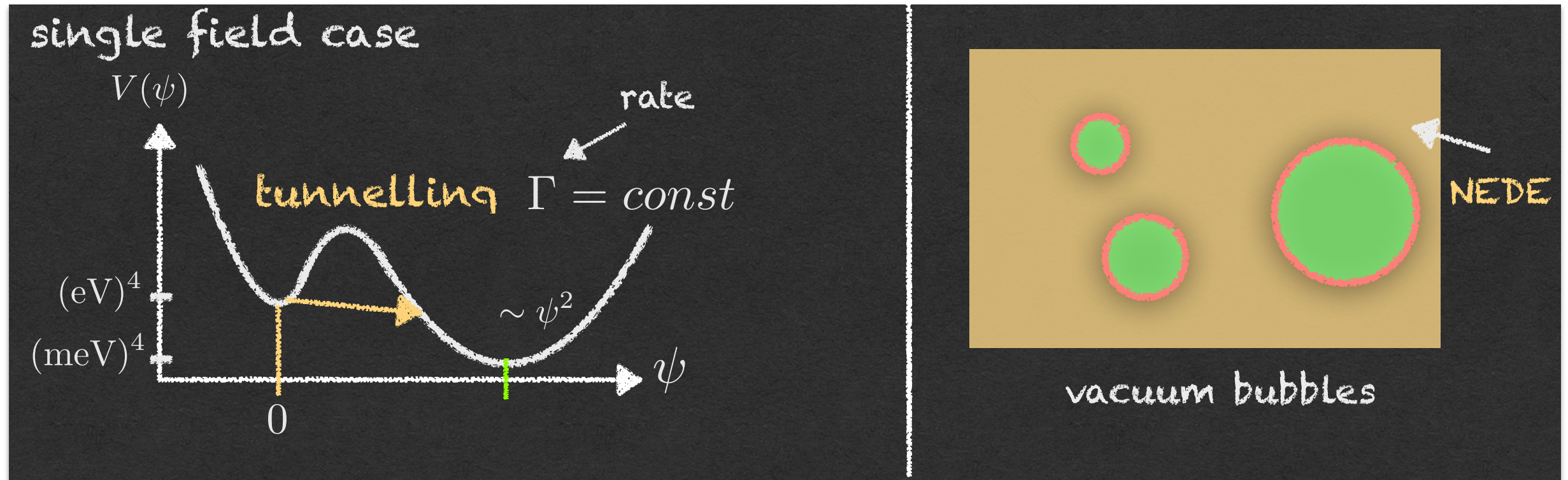
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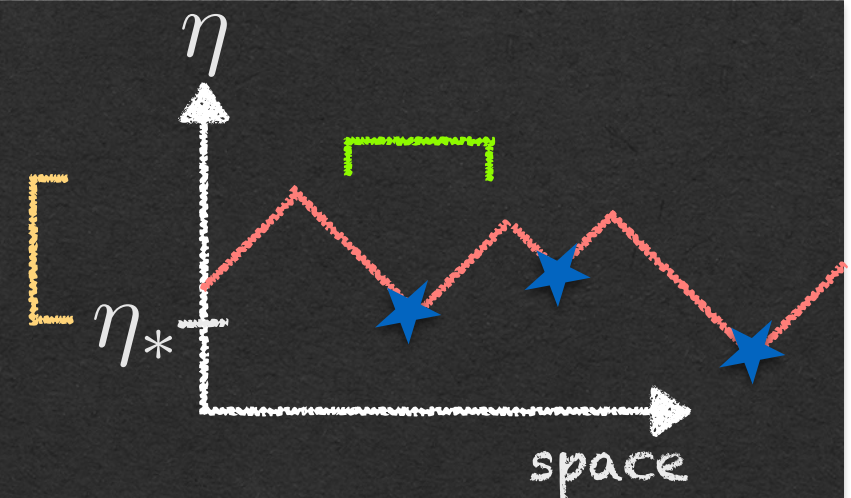
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► **Idea:** Make tunneling rate time dependent: Two models **Cold** and **Hot NEDE**.

Cold New Early Dark Energy

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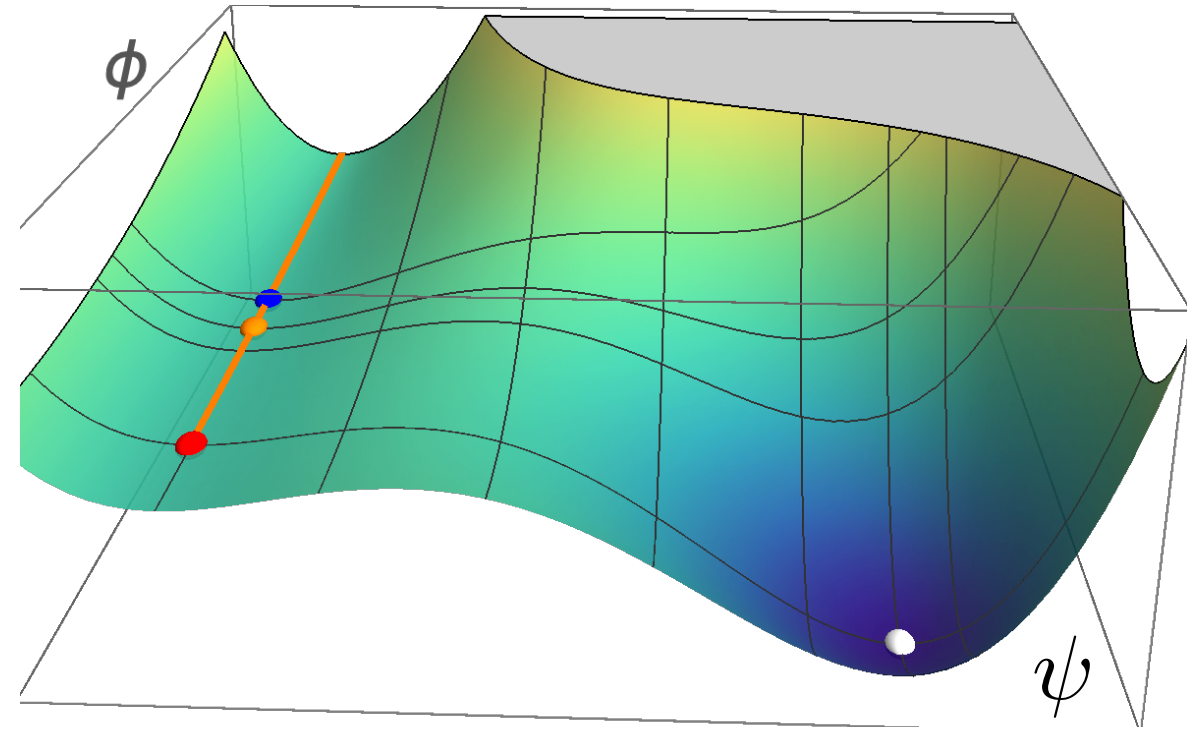
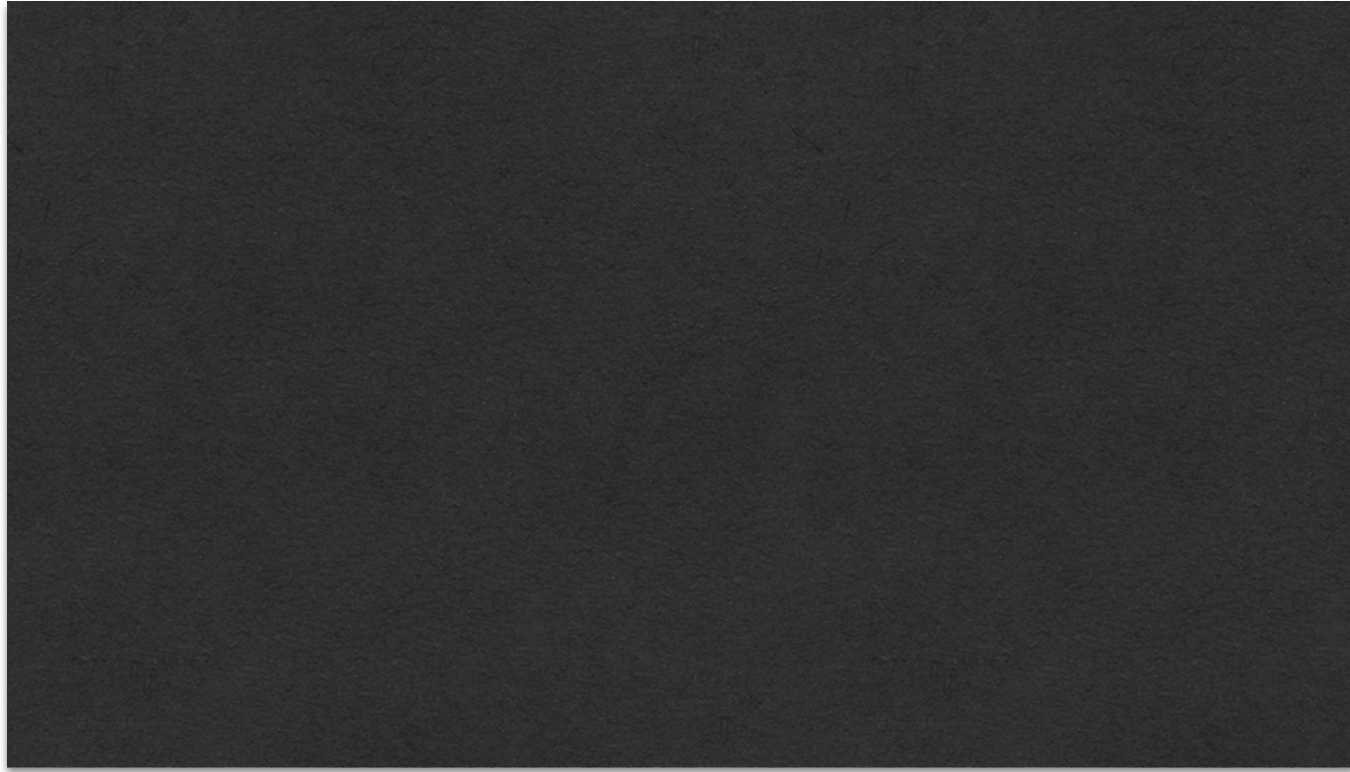
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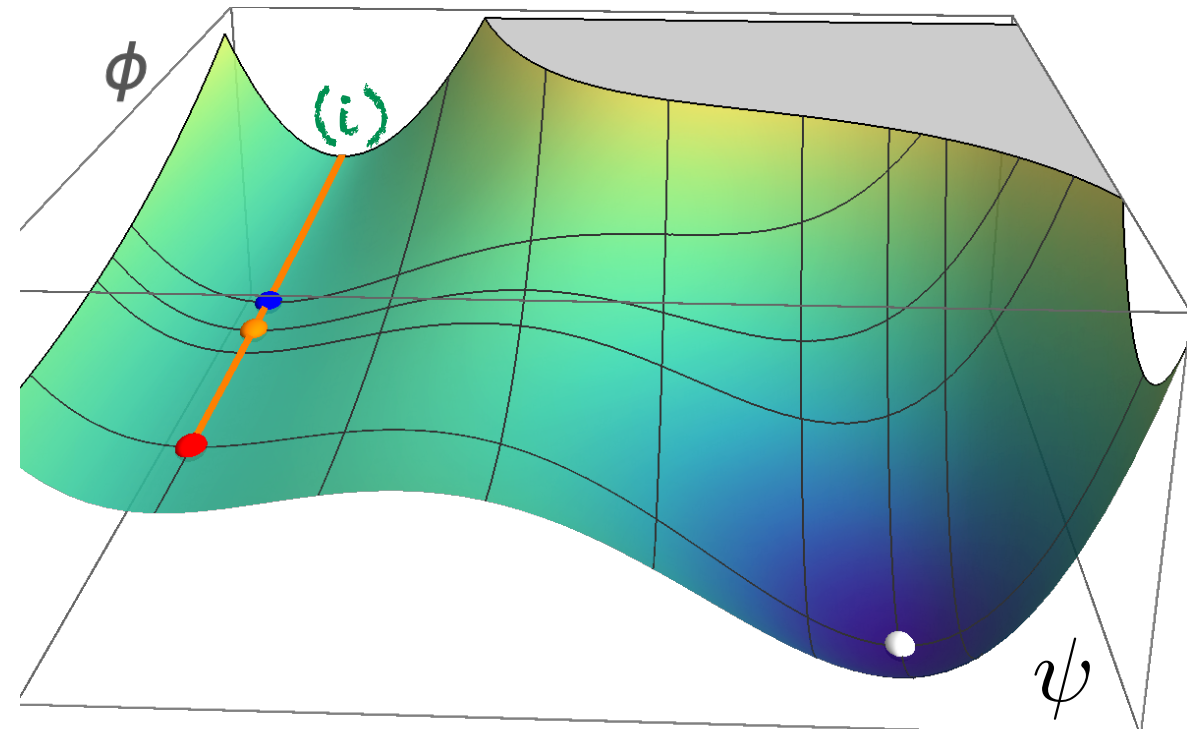
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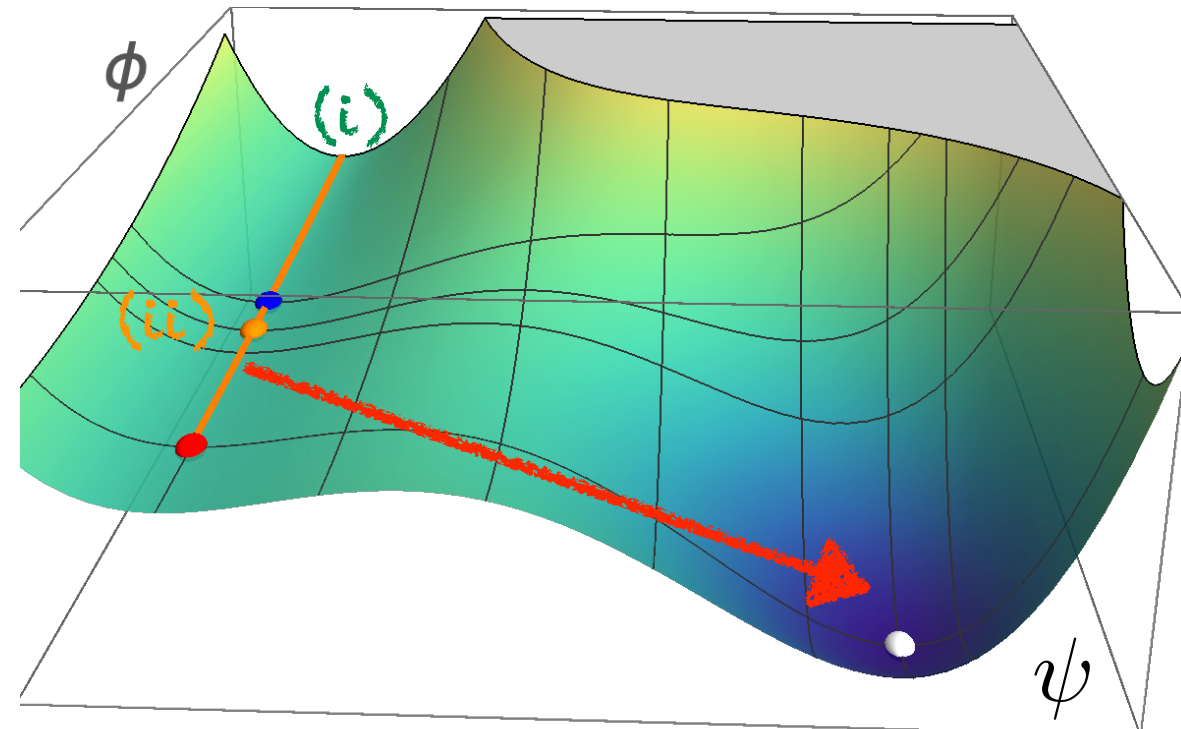
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recent bound:
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[G.Elor++, 2311.16222]



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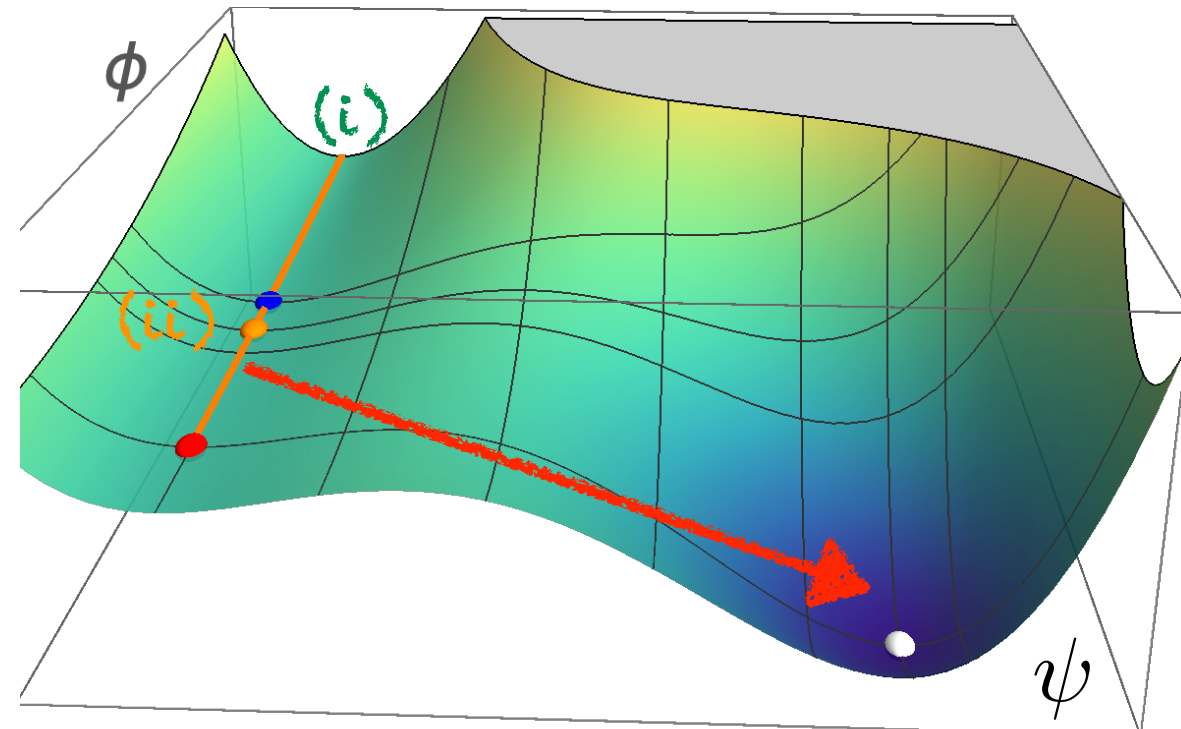
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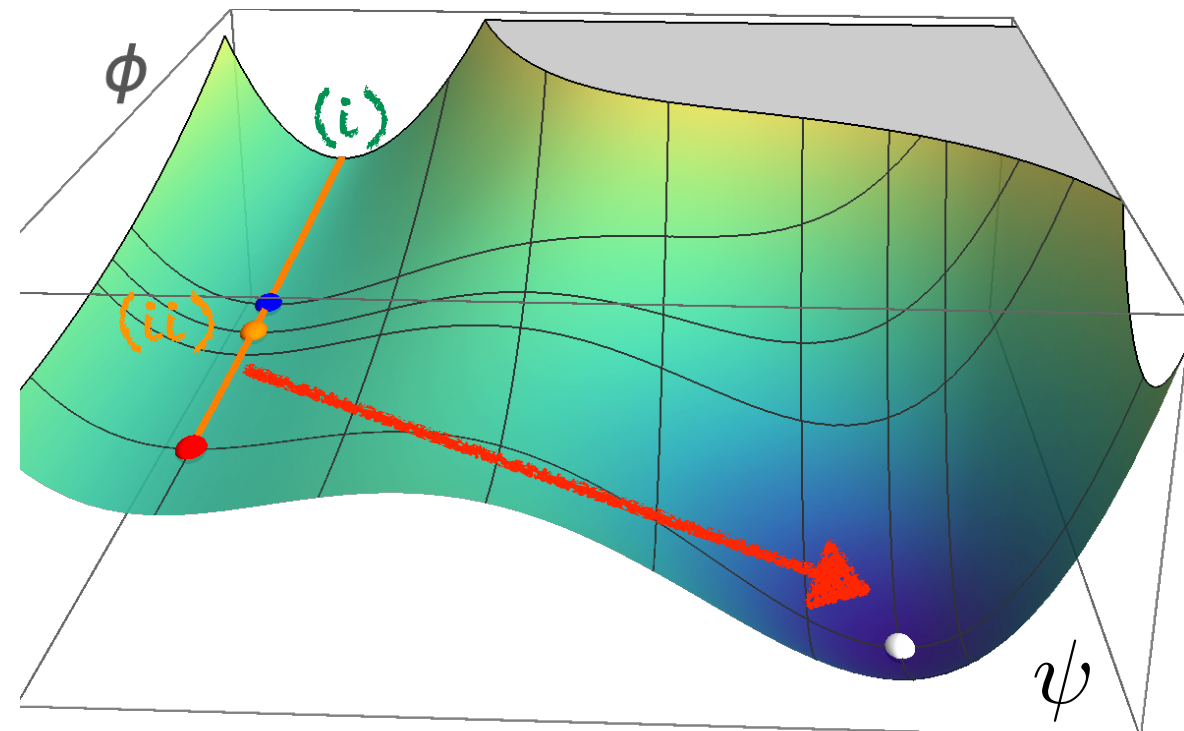
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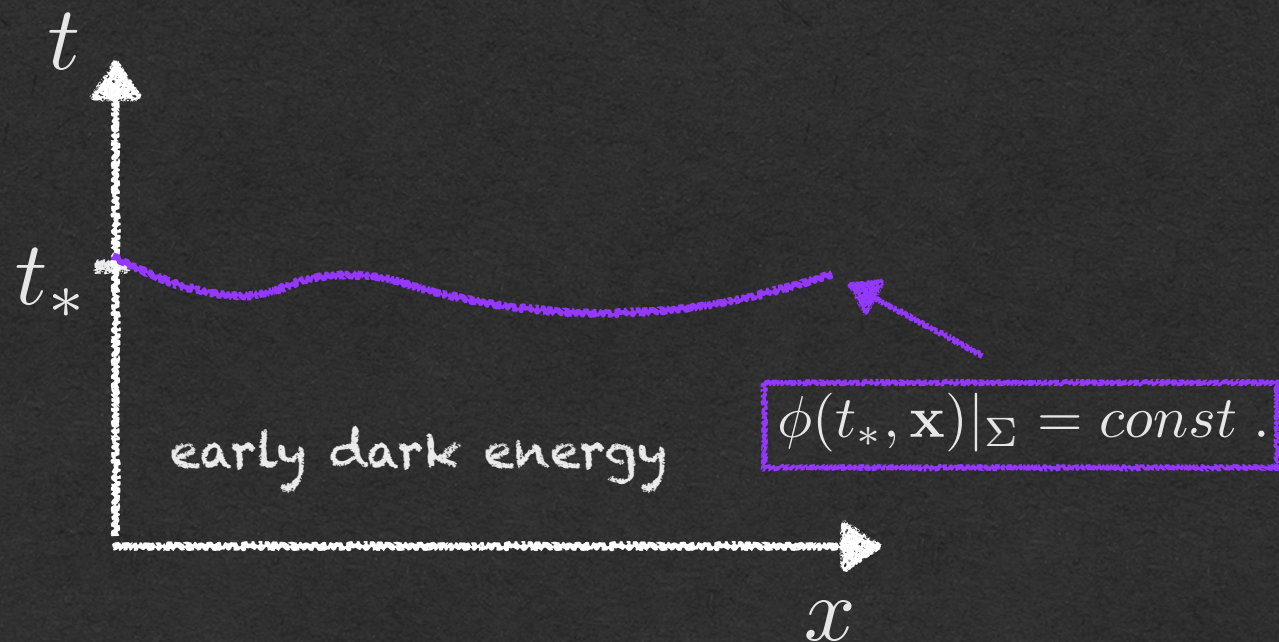
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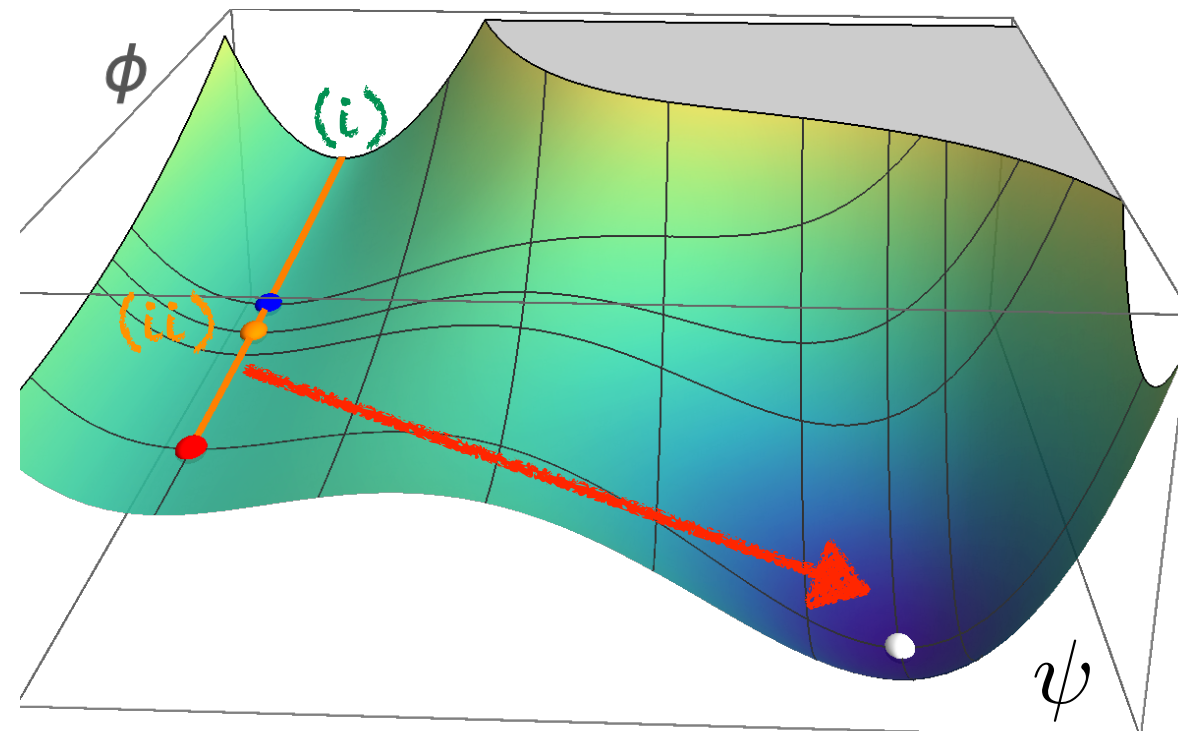
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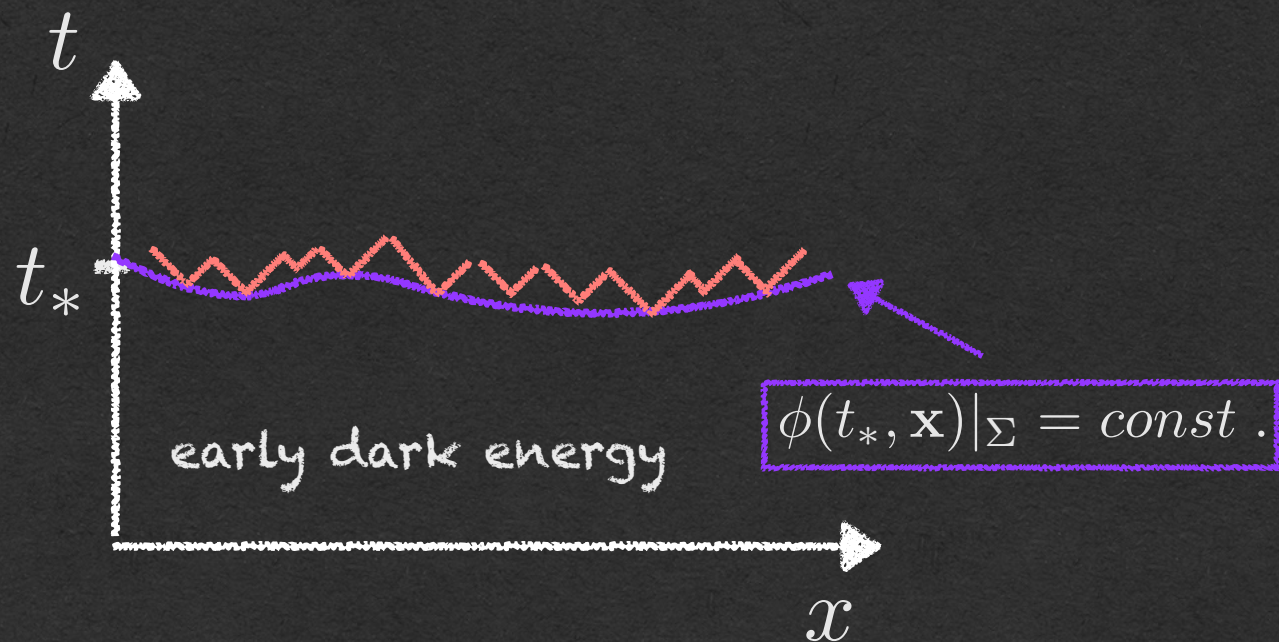
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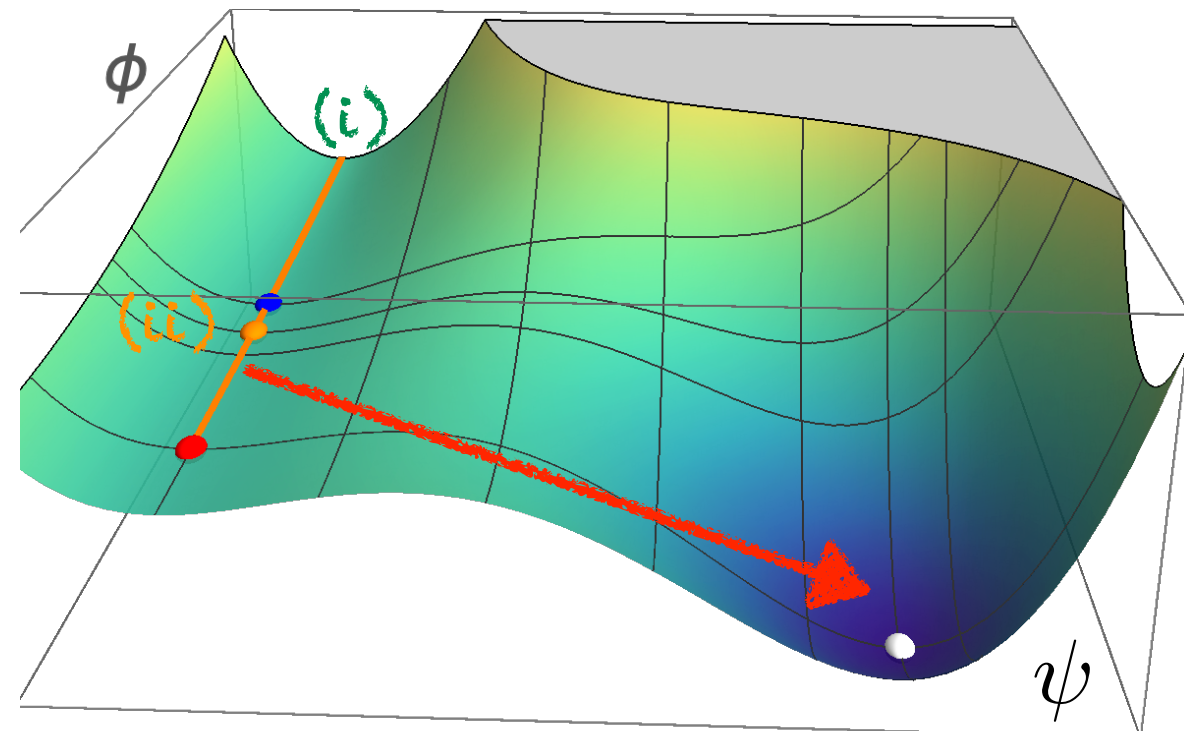
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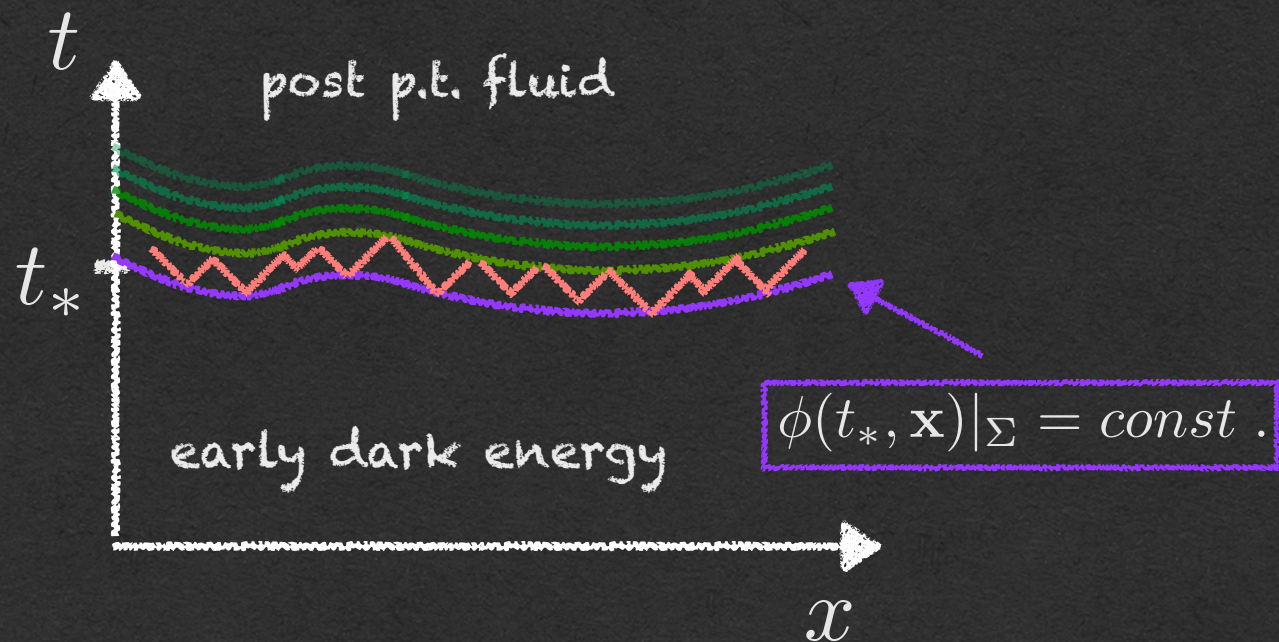
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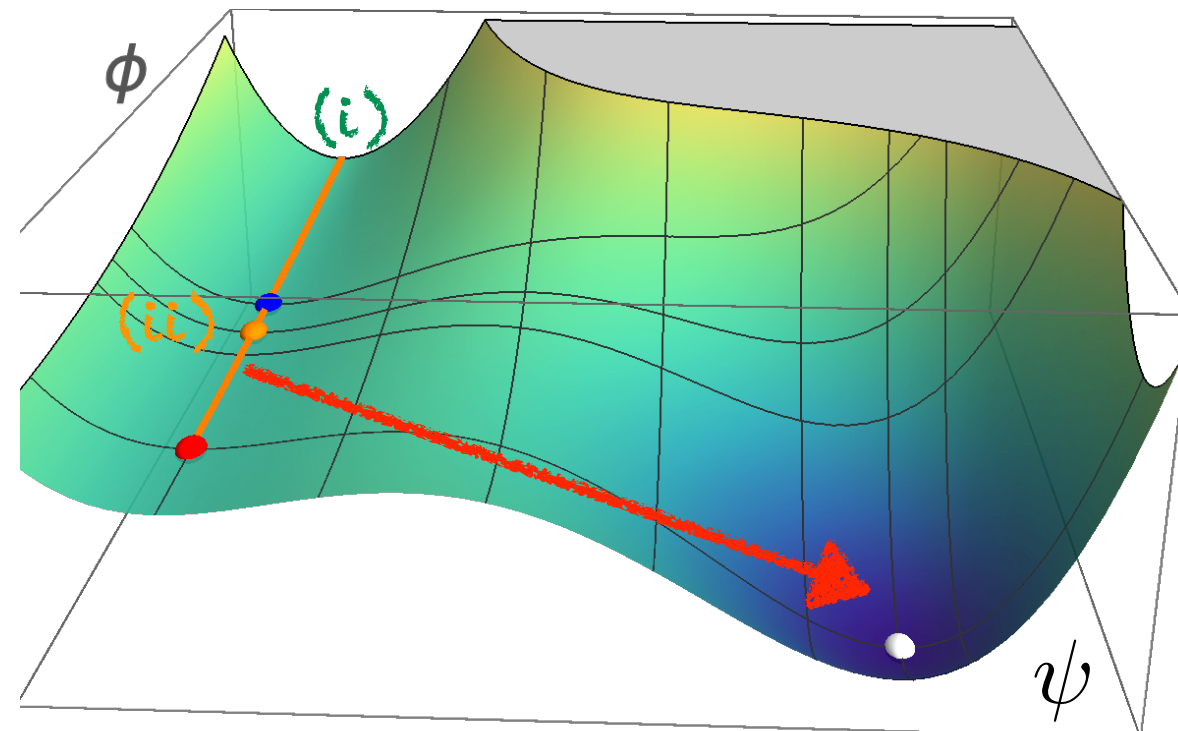
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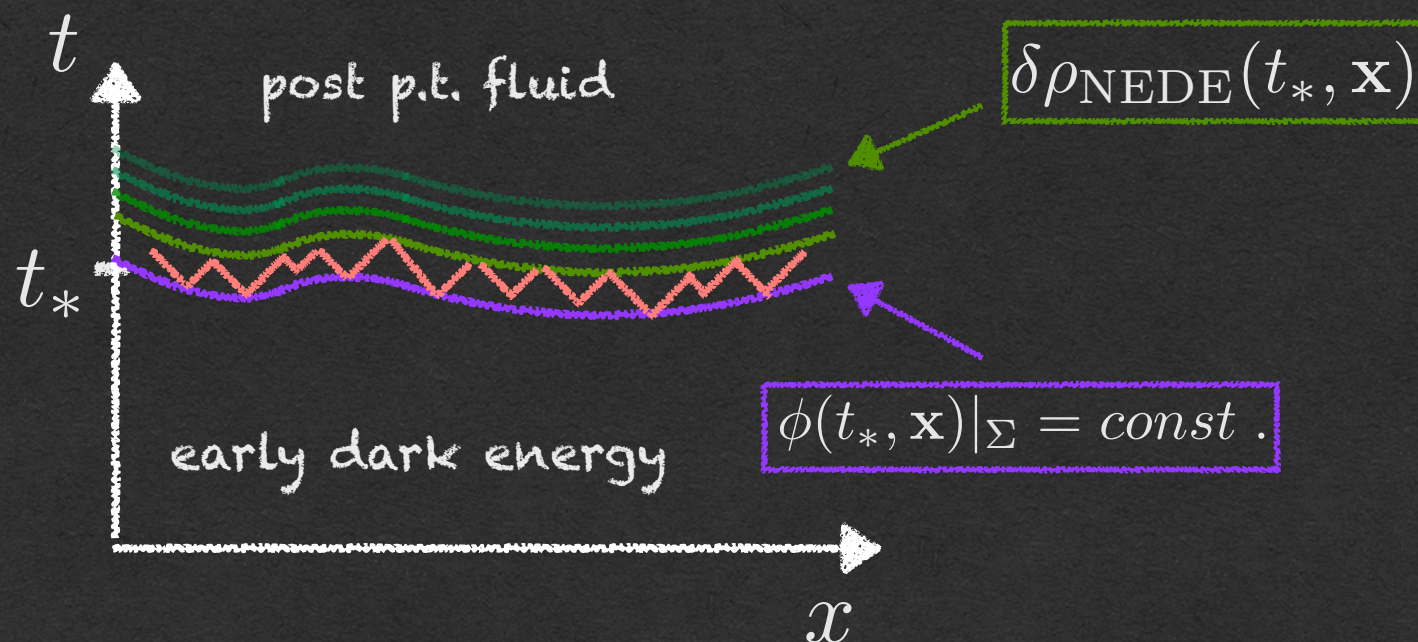
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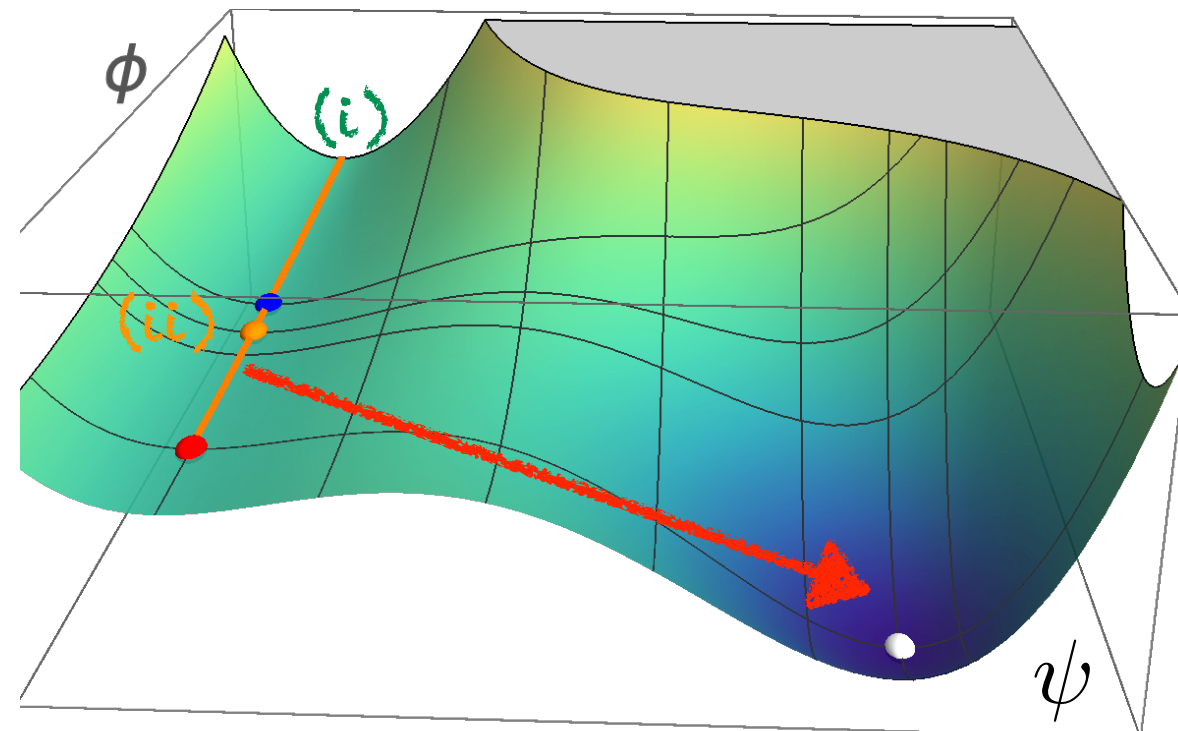
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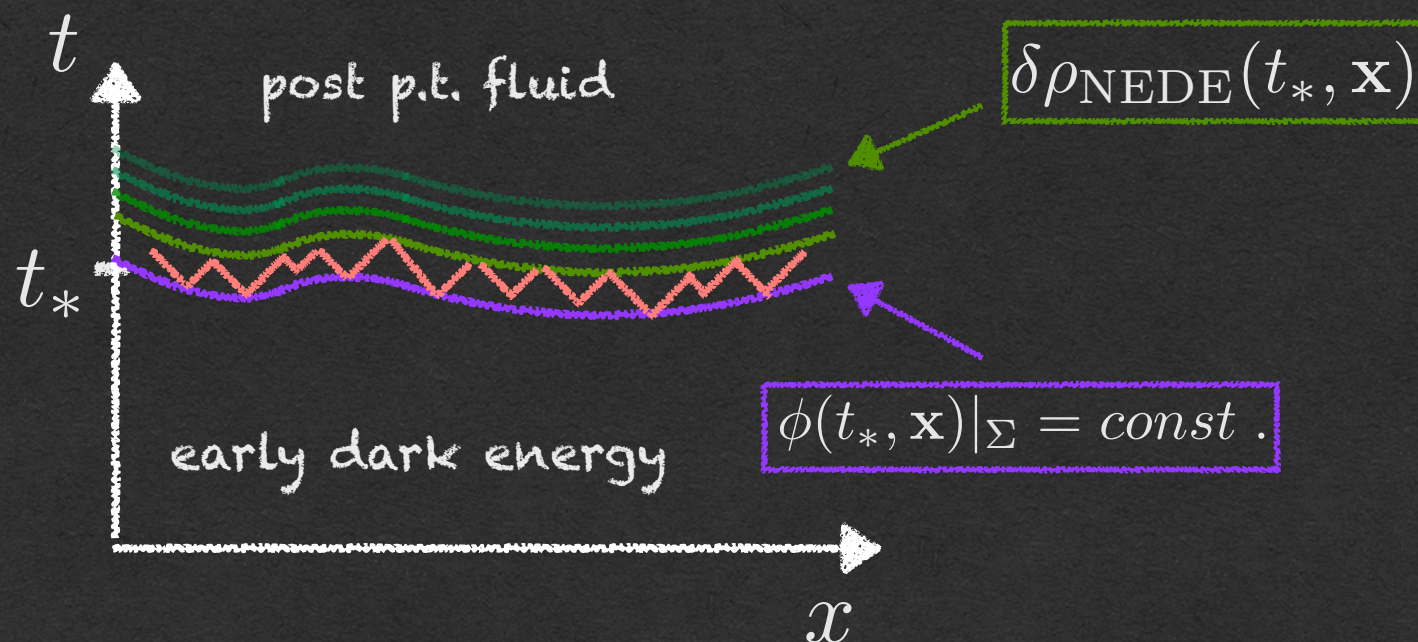
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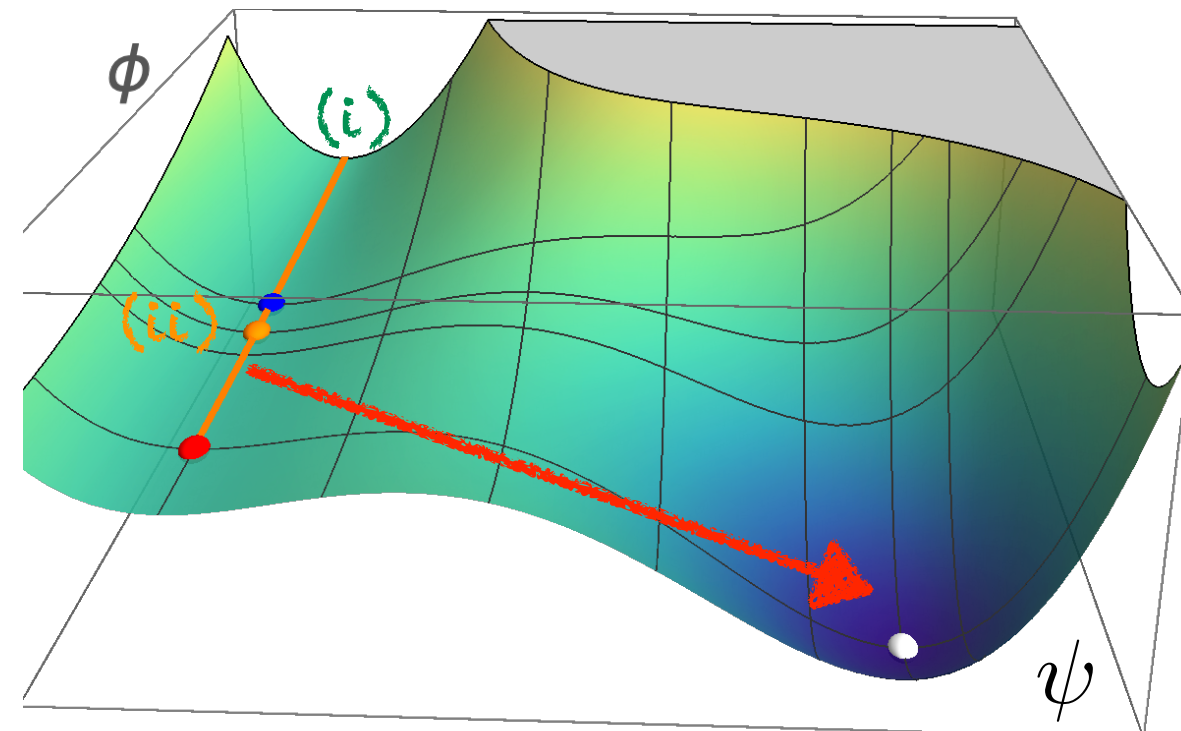
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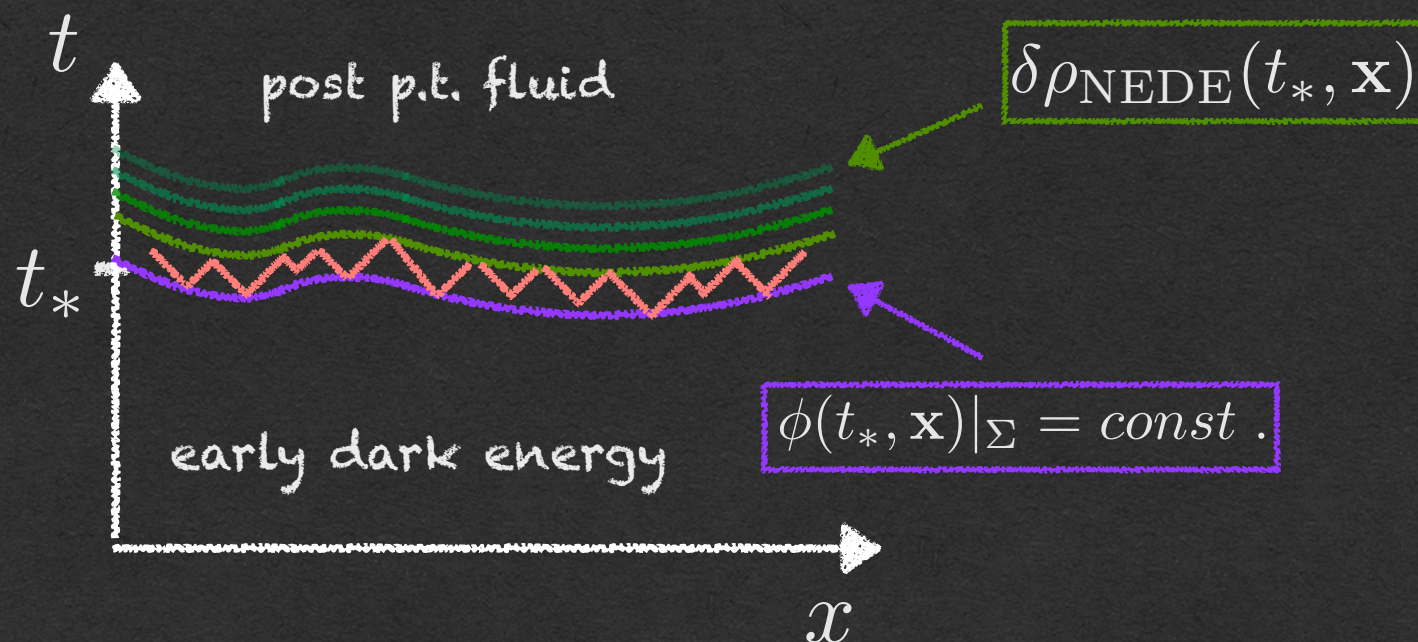
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A testable scenario

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The H_0 Olympics: A fair ranking of proposed models

Schöneberg et al., 2022

Nils Schöneberg,^a Guillermo Franco Abellán,^b Andrea Pérez Sánchez,^a Samuel J. Witte,^c Vivian Poulin,^b and Julien Lesgourgues^a

Model	ΔN_{param}	M_B	Gaussian Tension	Q_{DMAP} Tension		$\Delta\chi^2$	ΔAIC		Finalist
ΛCDM	0	-19.416 ± 0.012	4.4σ	4.5σ	\times	0.00	0.00	\times	\times
Majoron	3	-19.380 ± 0.027	3.0σ	2.9σ	✓	-13.74	-7.74	✓	✓ ②
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axiEDE

NEDE

[Planck 2018 + BAO + Pantheon (+ SH0ES)]

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EDE-type models reduce tension to ~ 2 sigma.

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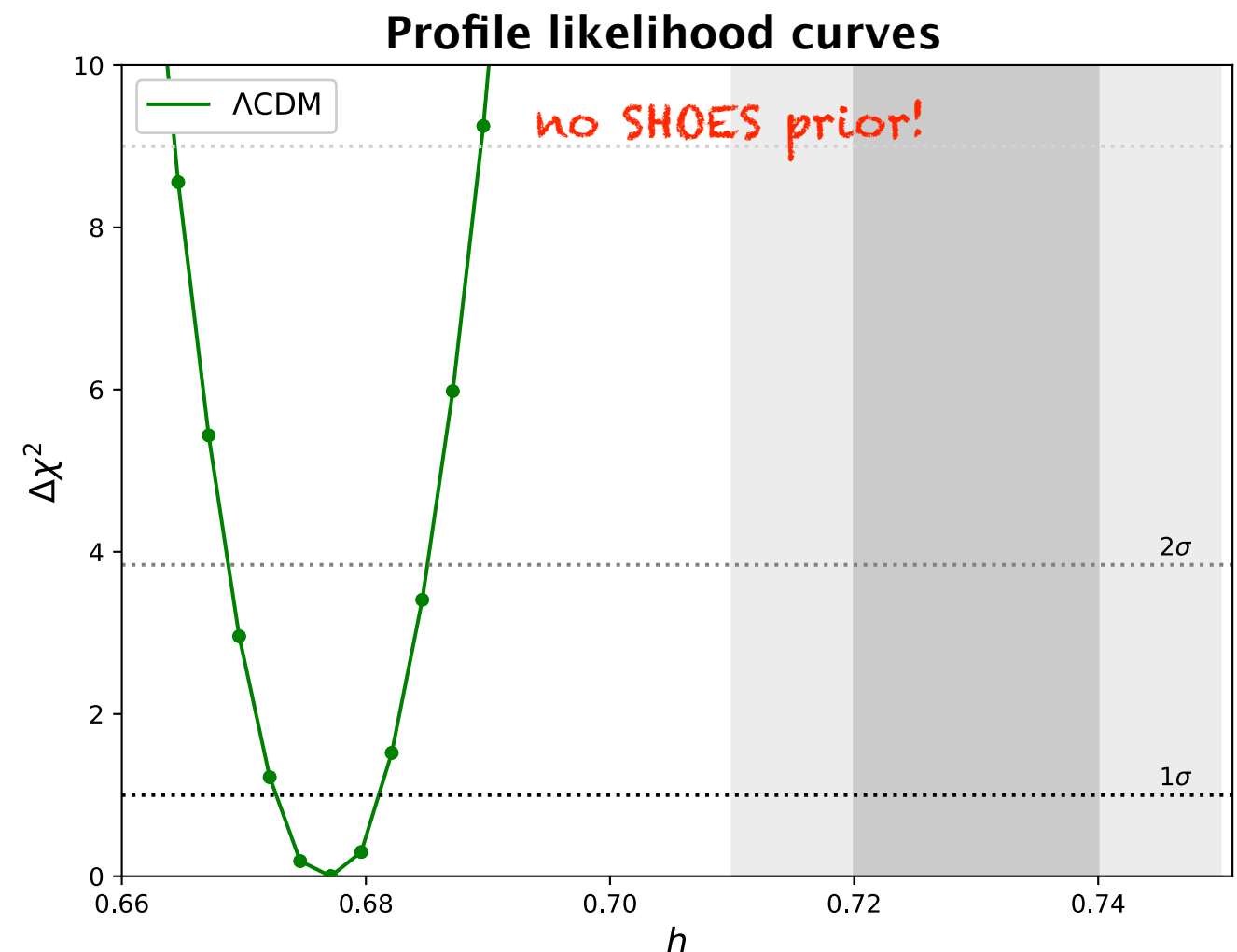
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Update 2024

(with A. Chatrchyan, V. Poulin, and M. S. Sloth)



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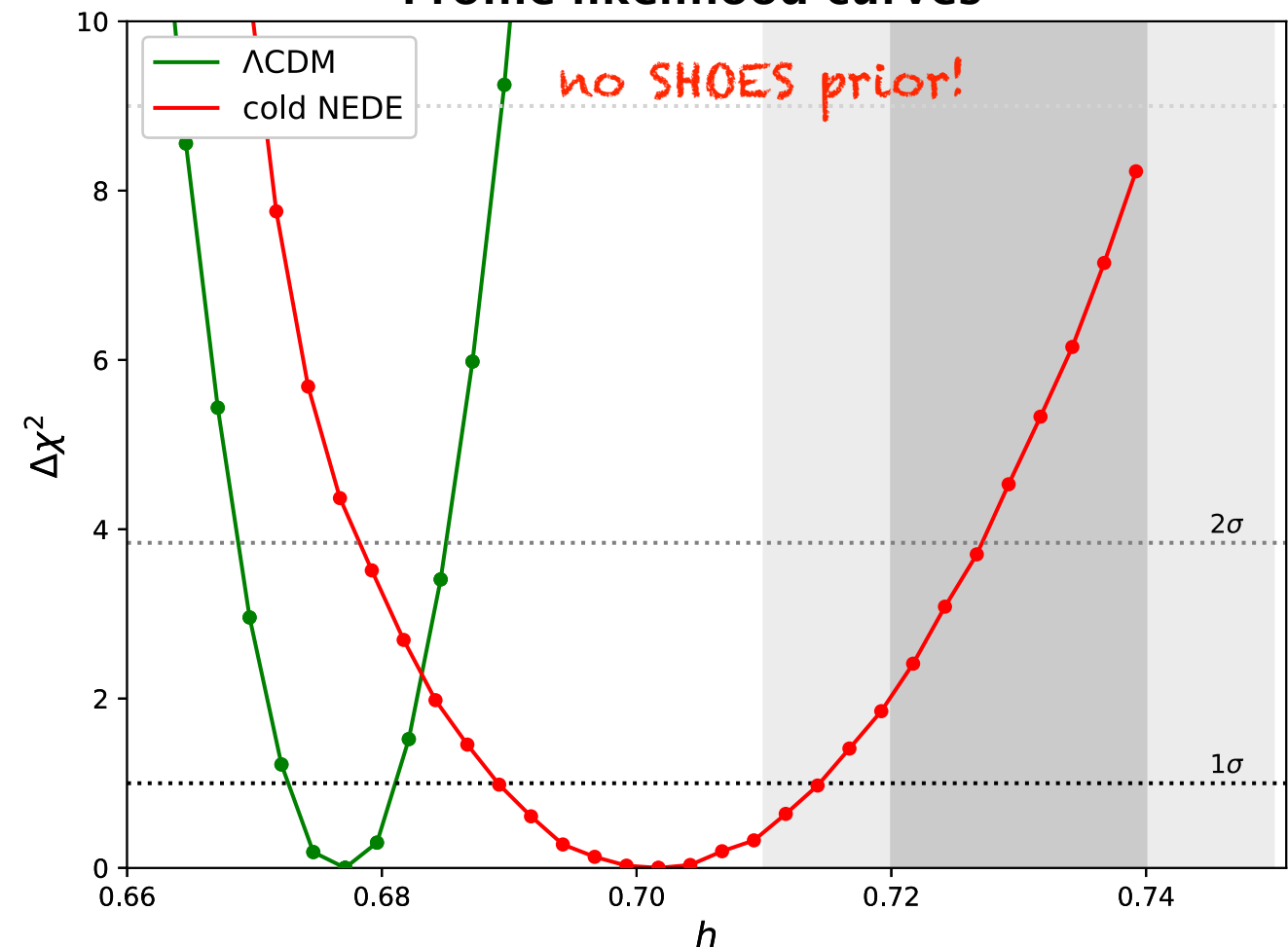
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Schöneberg et al., 2022

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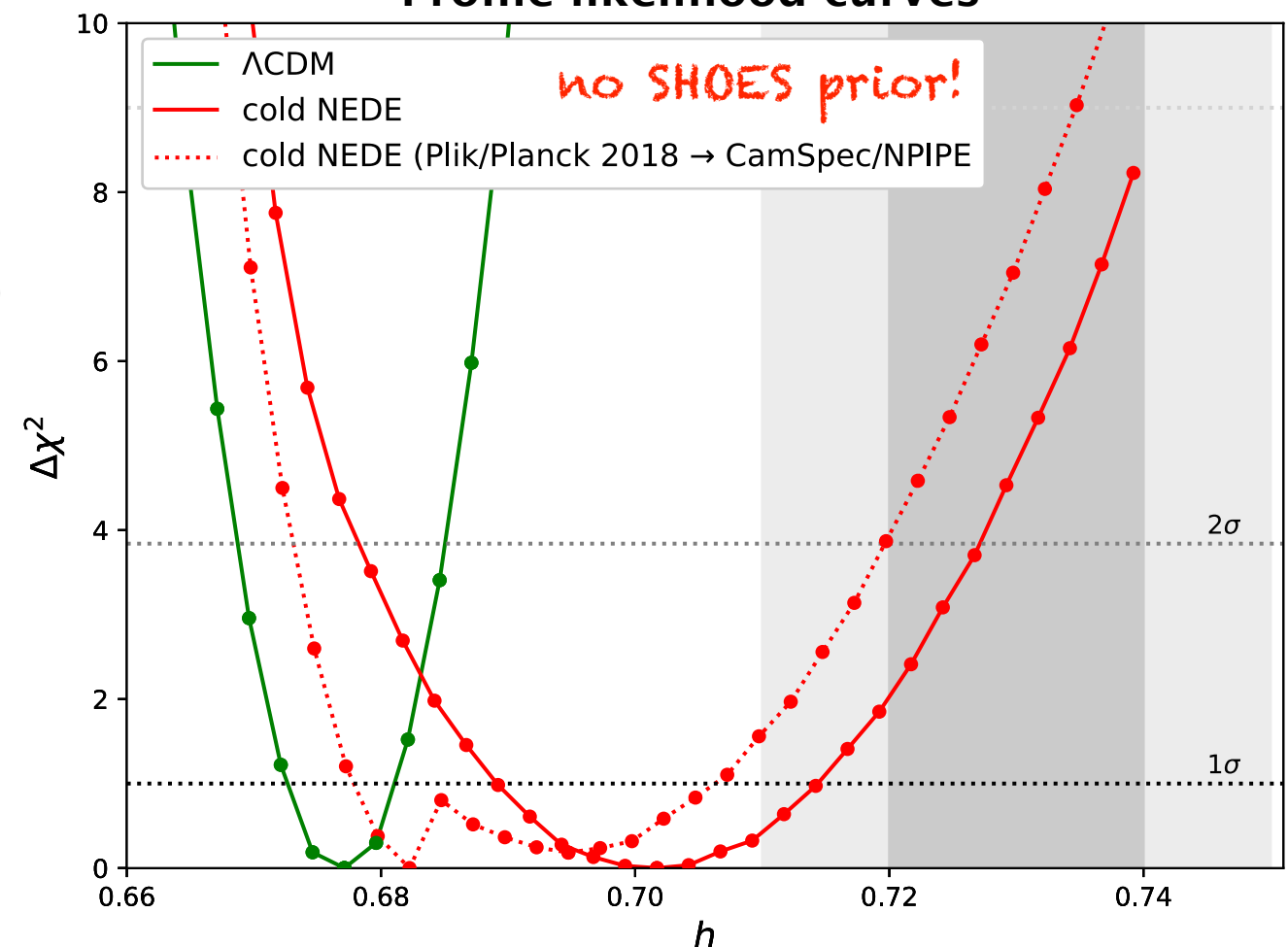
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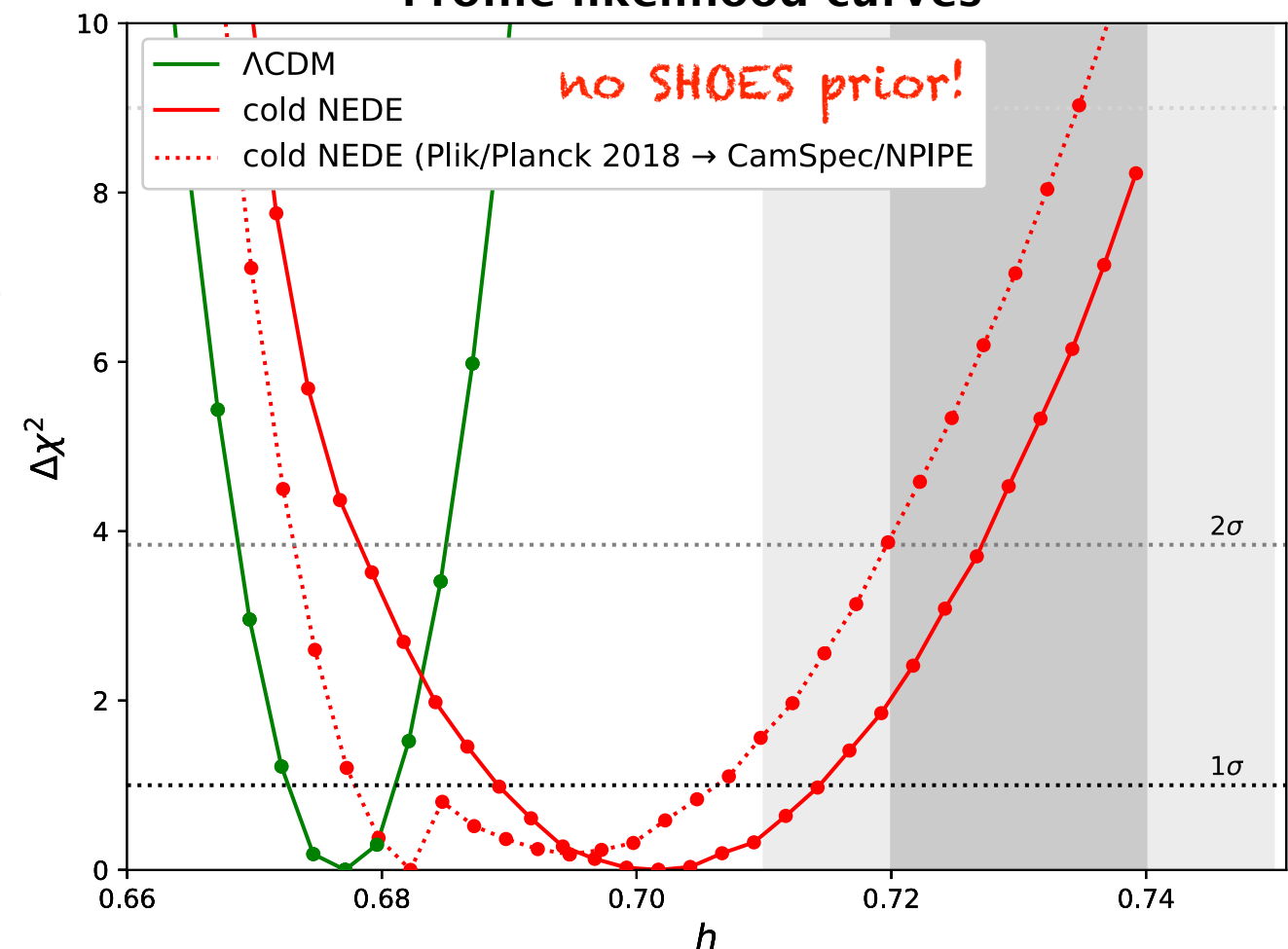
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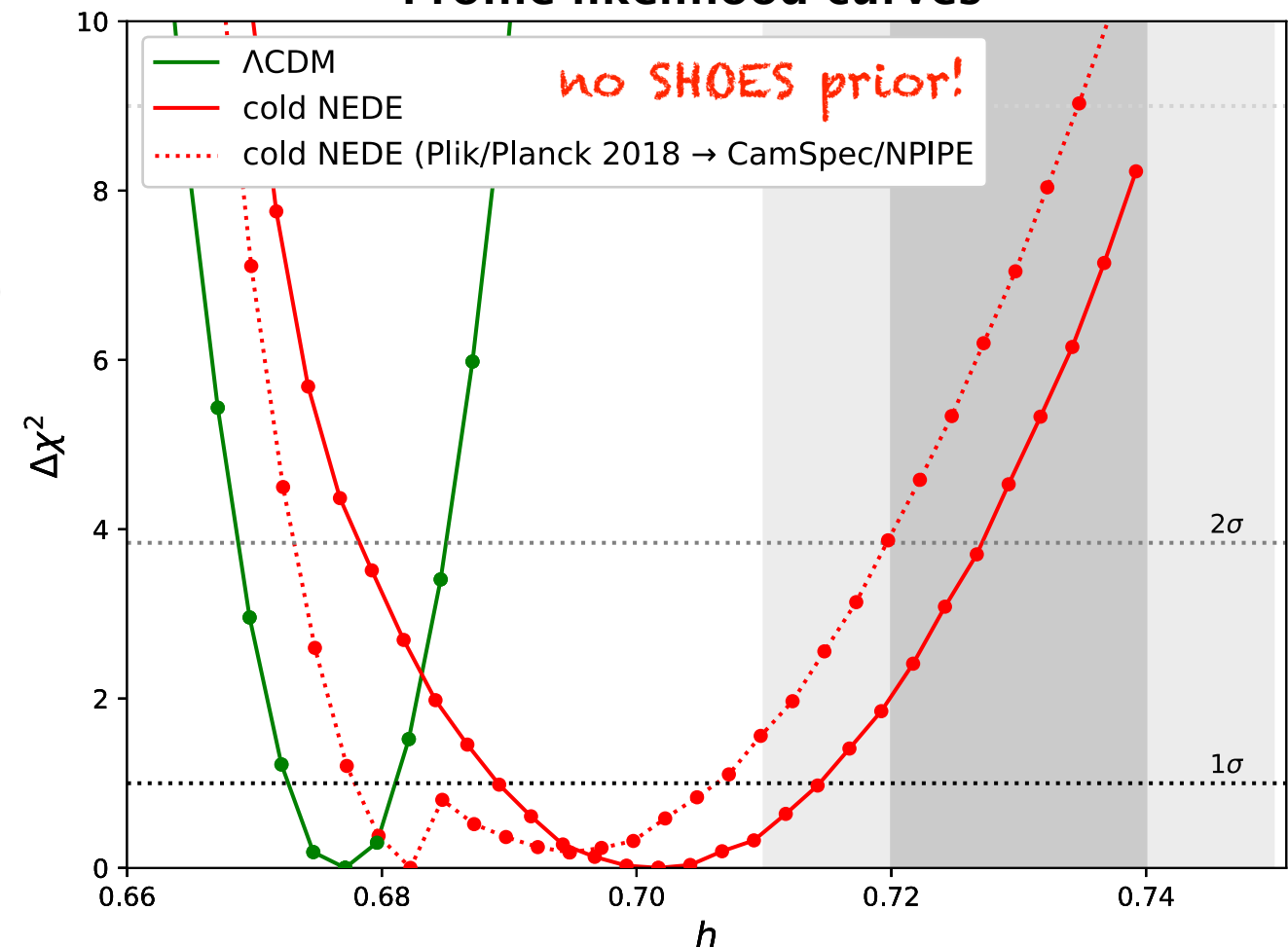
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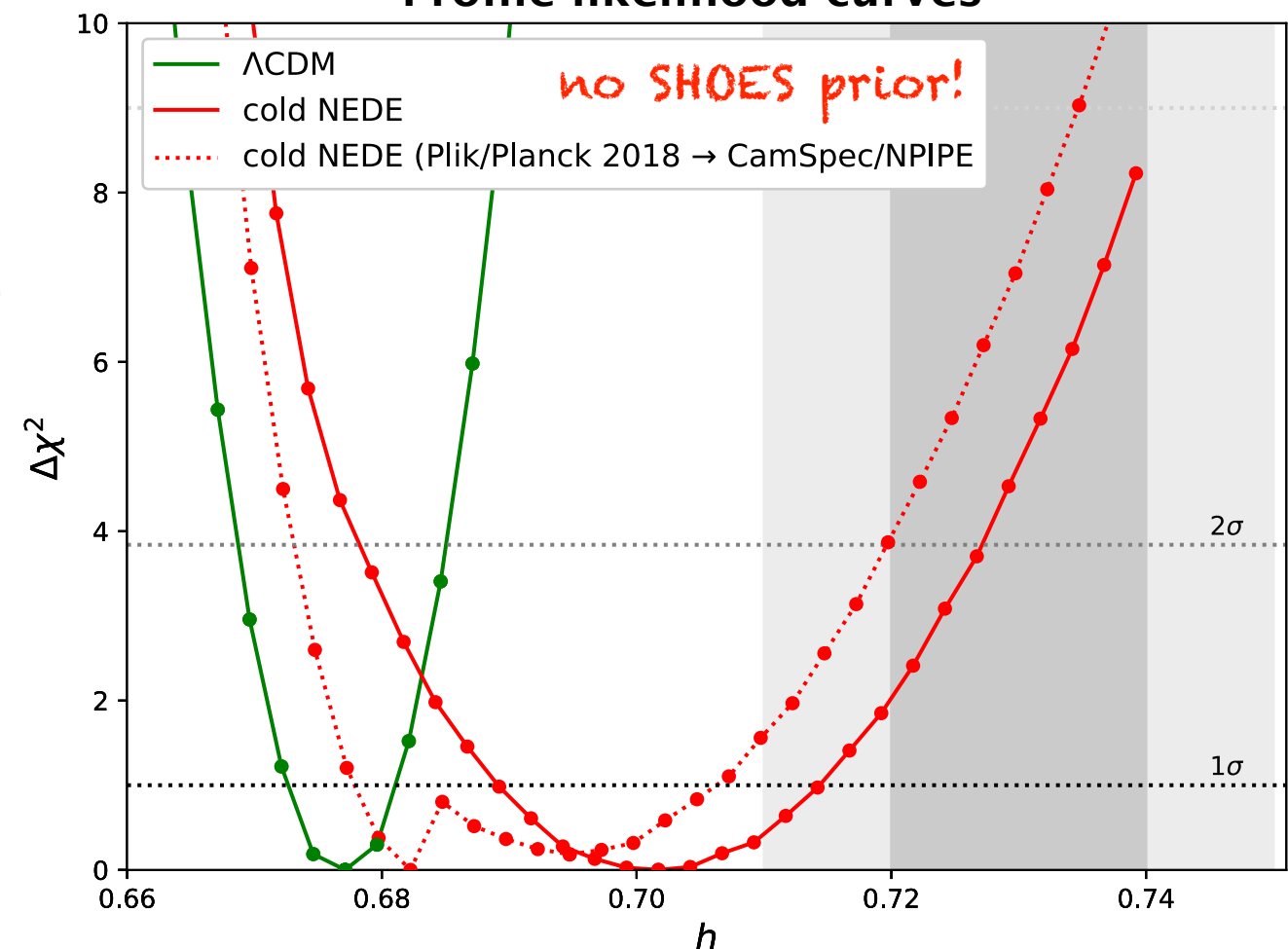
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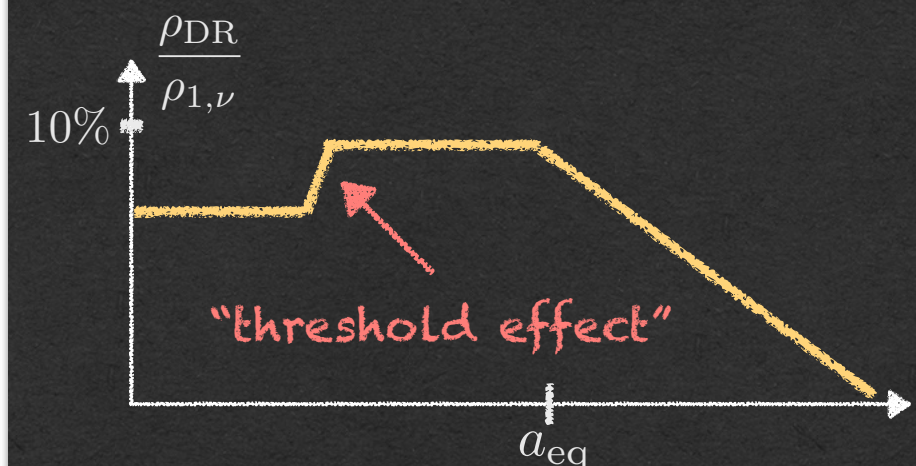
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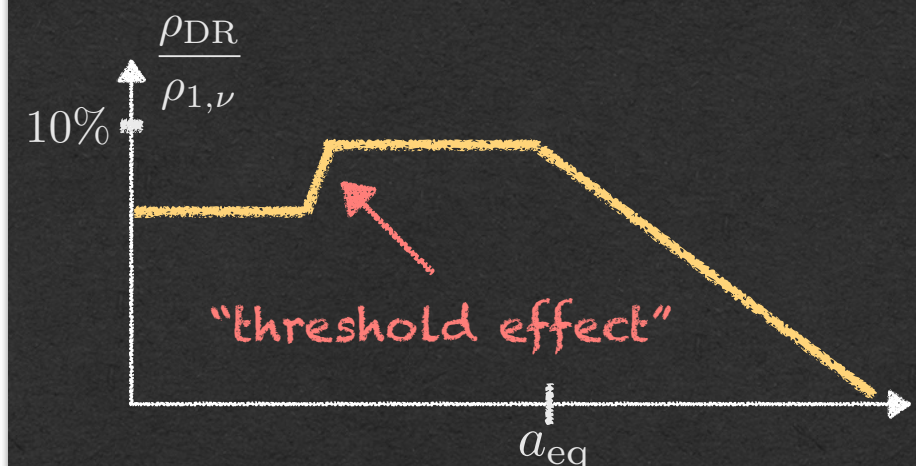
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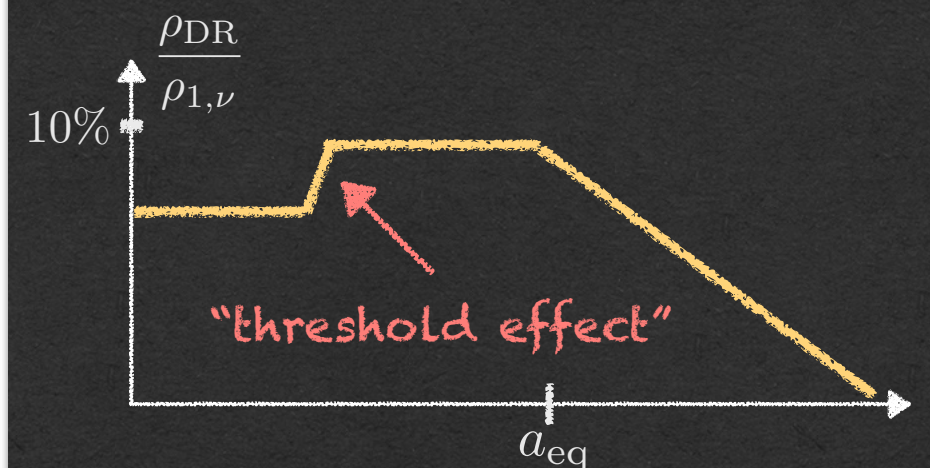
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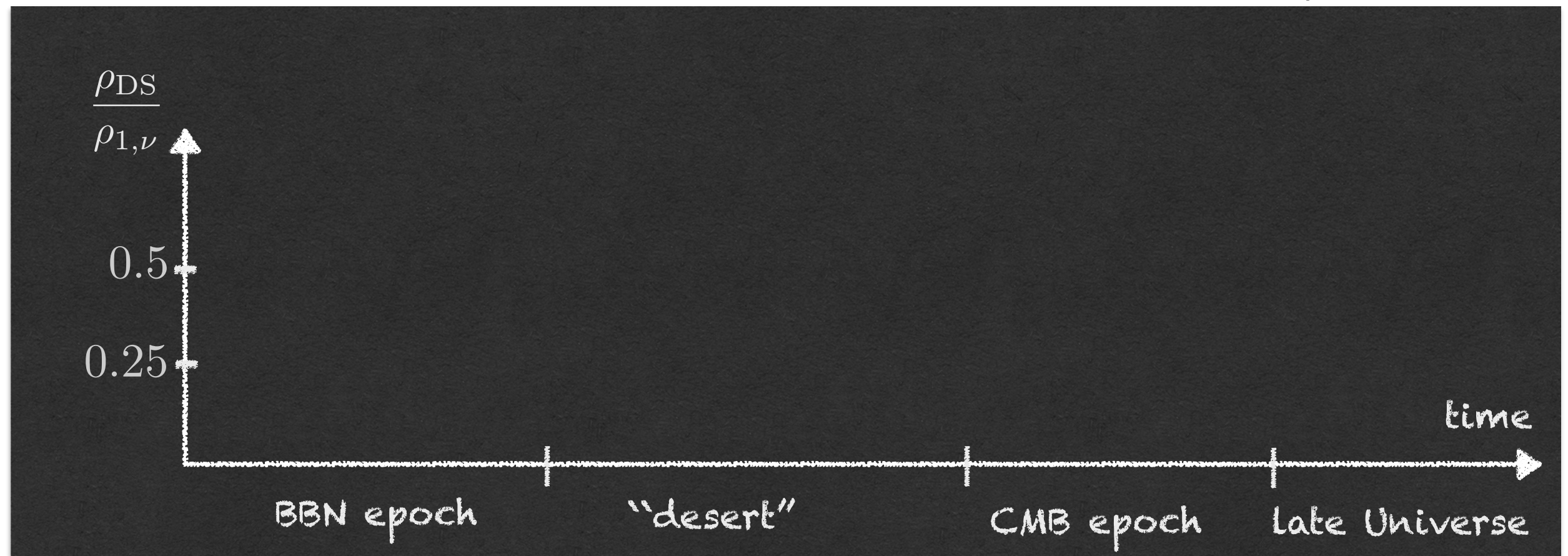
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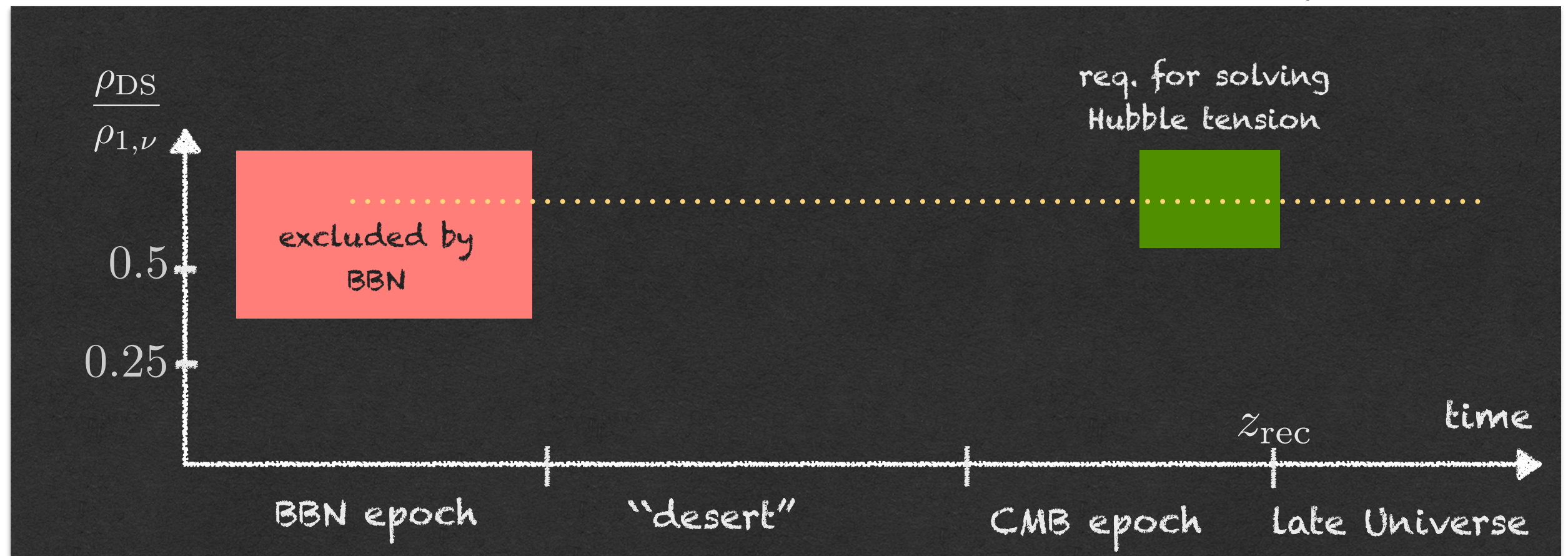
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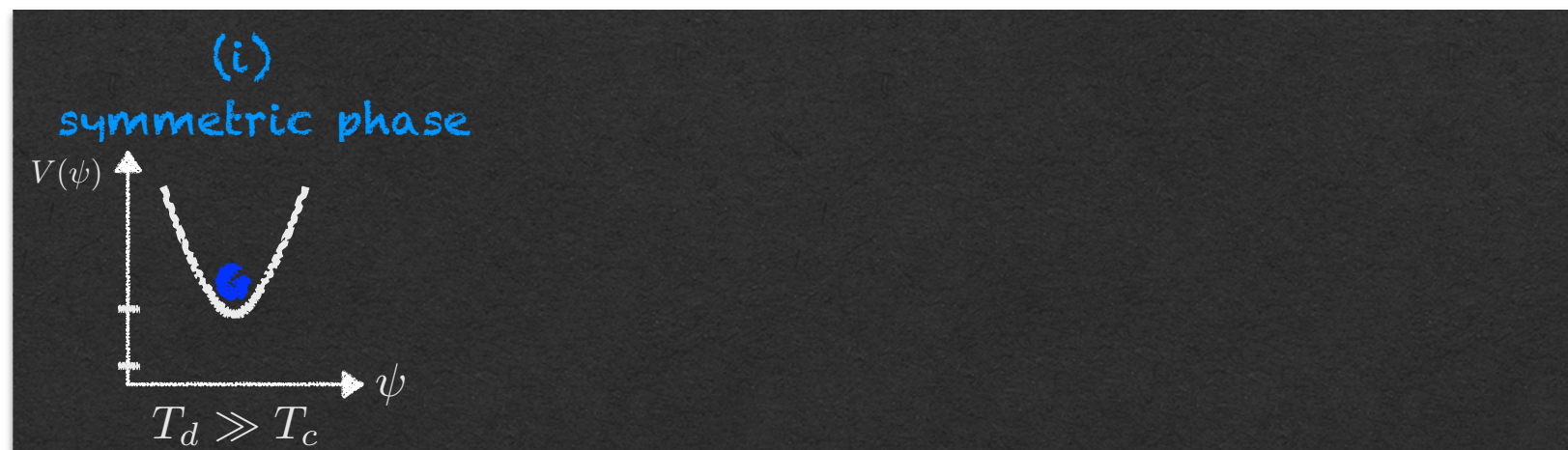
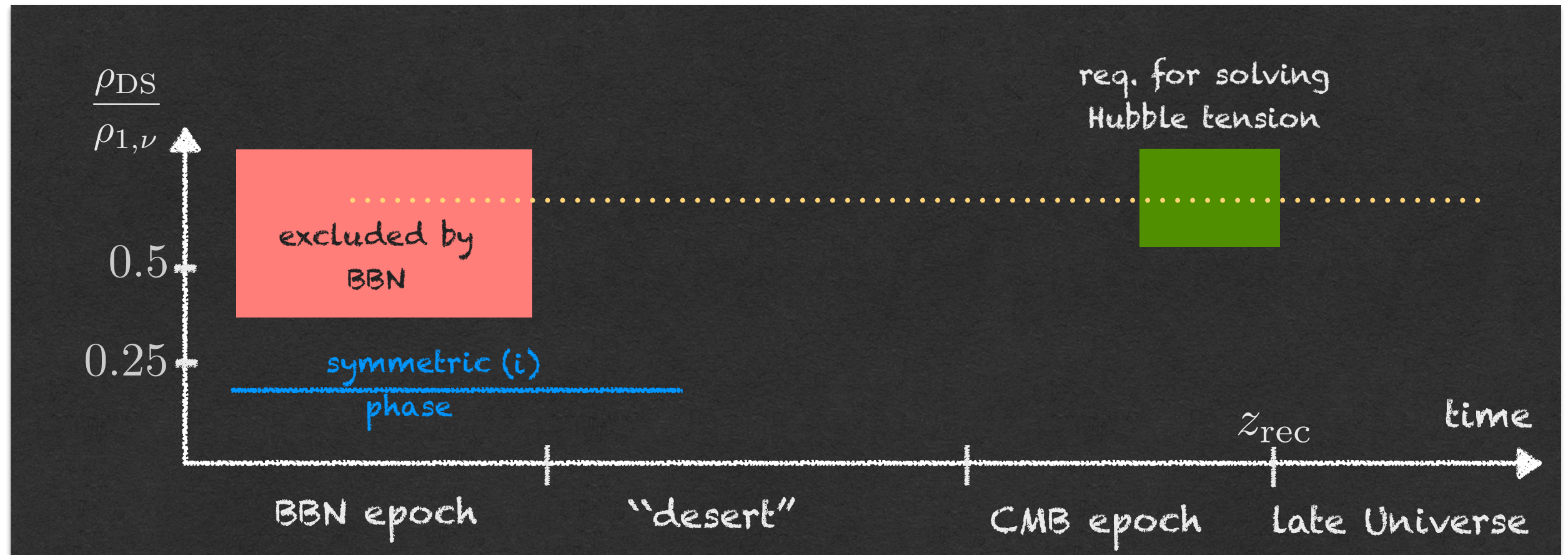
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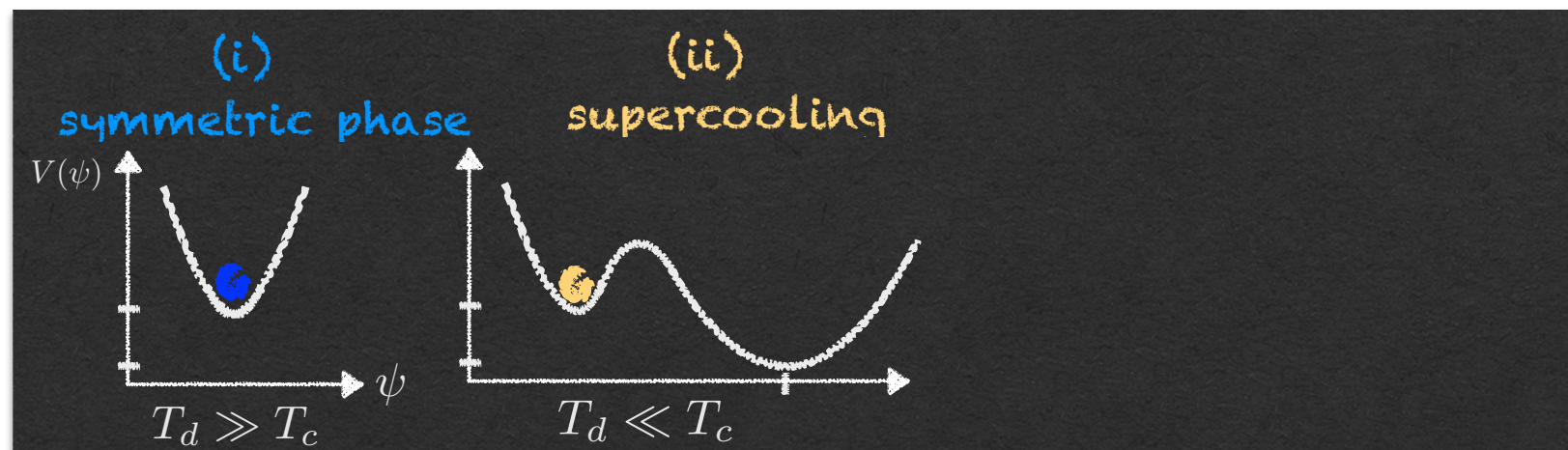
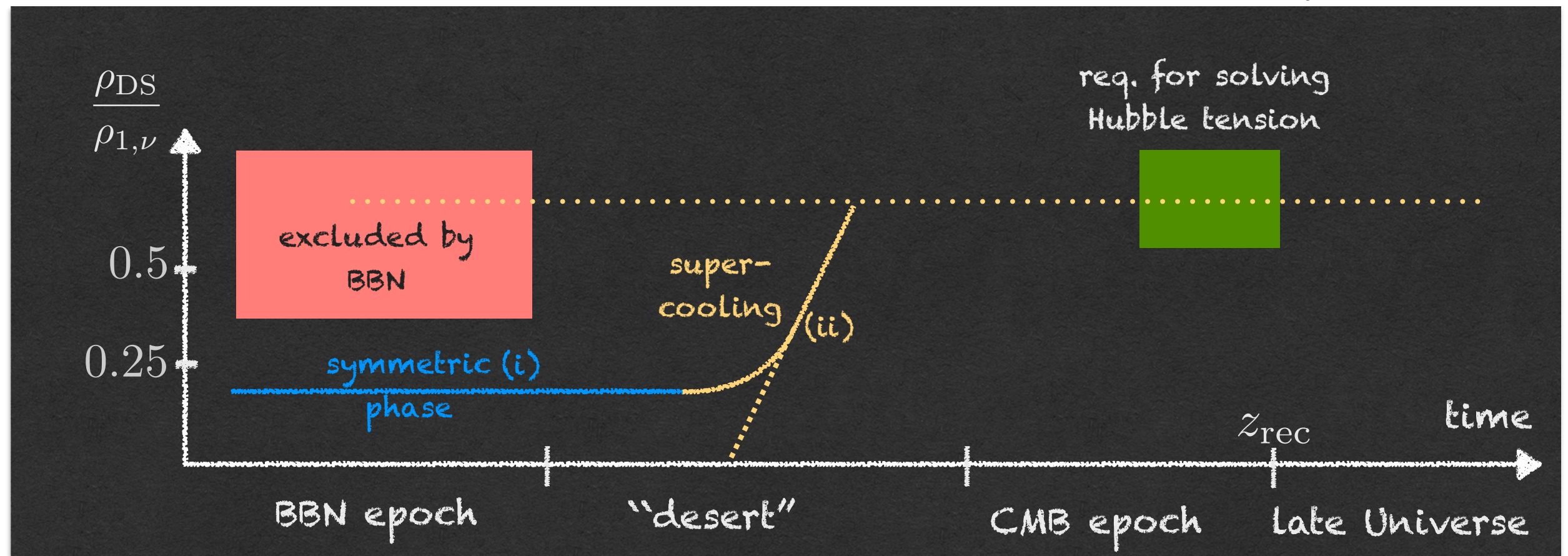
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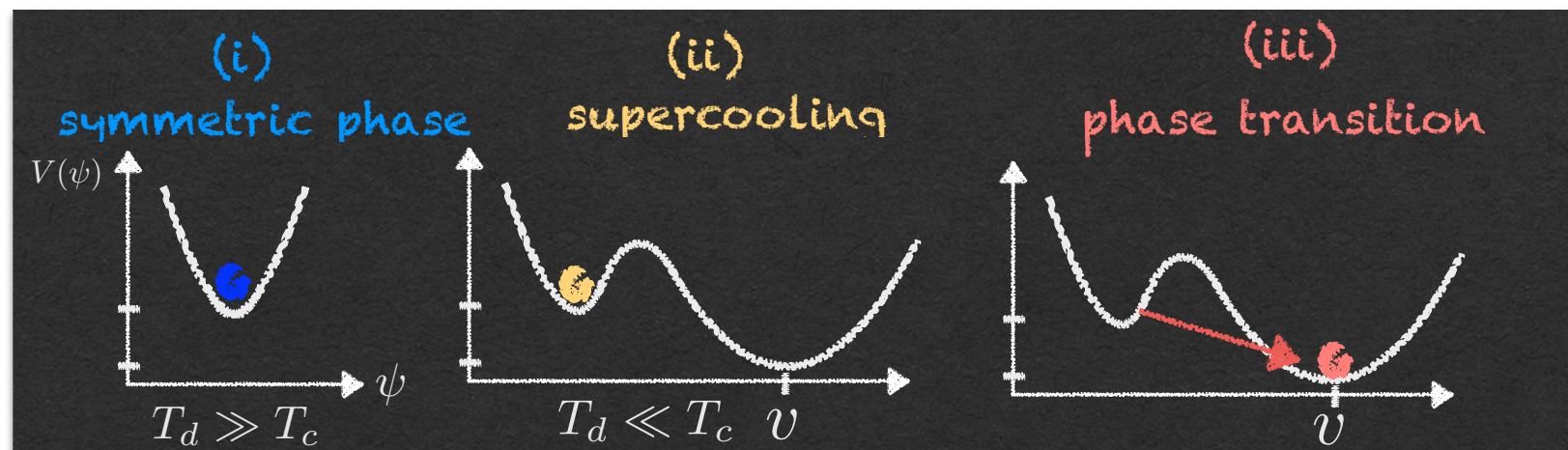
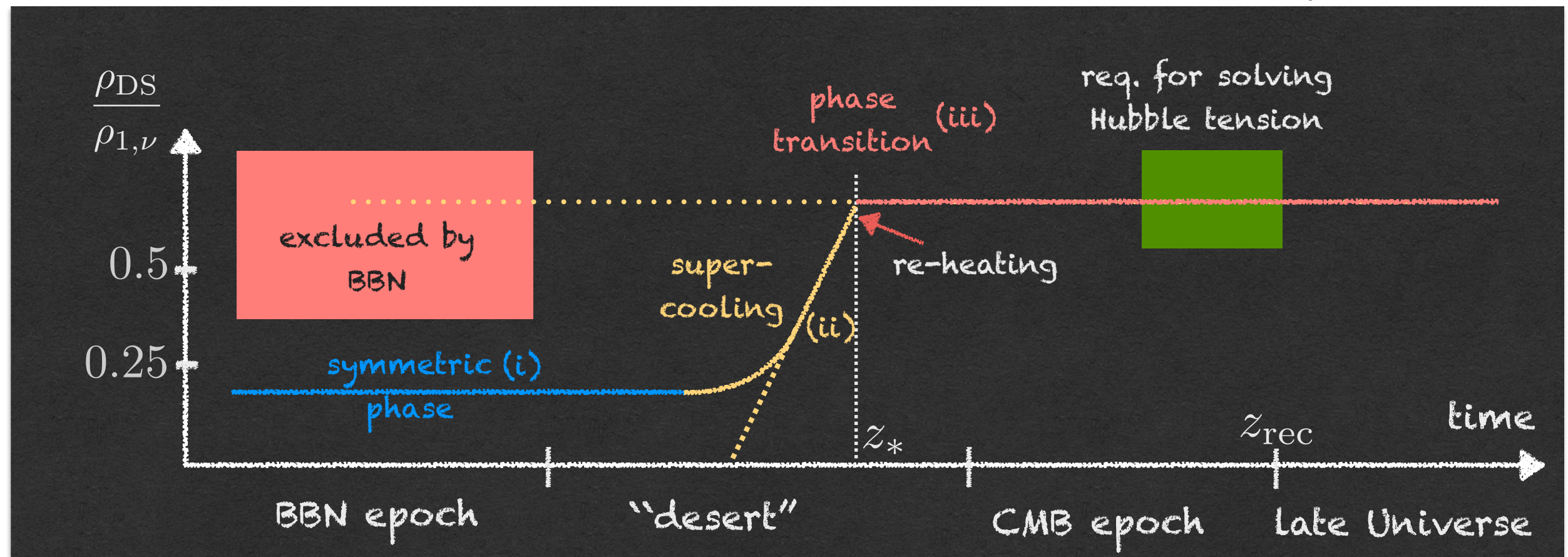
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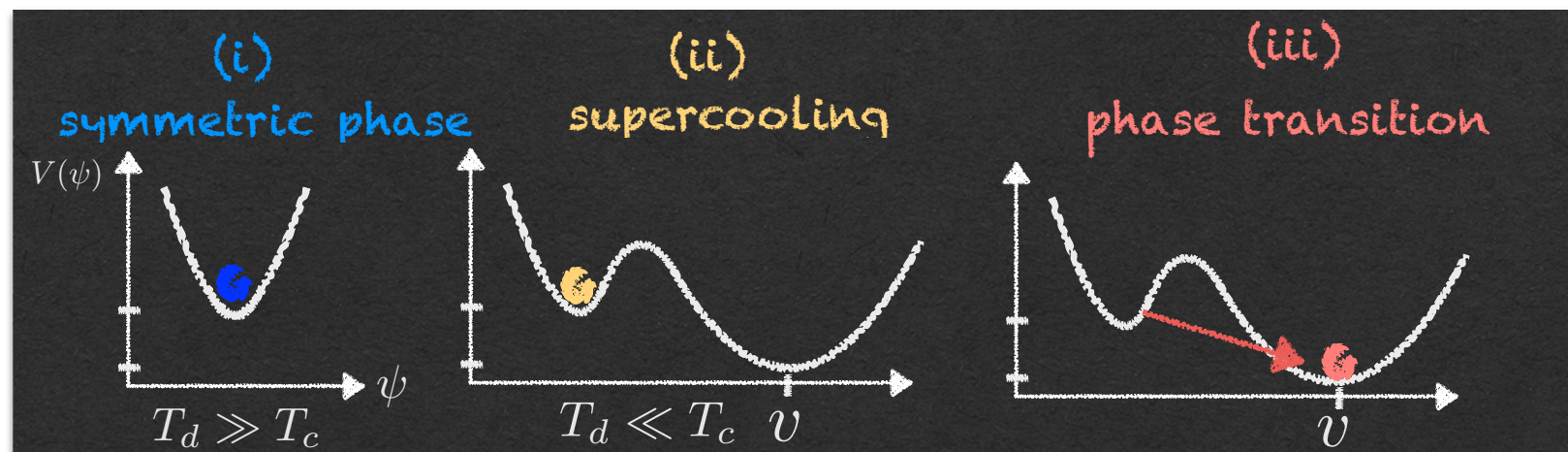
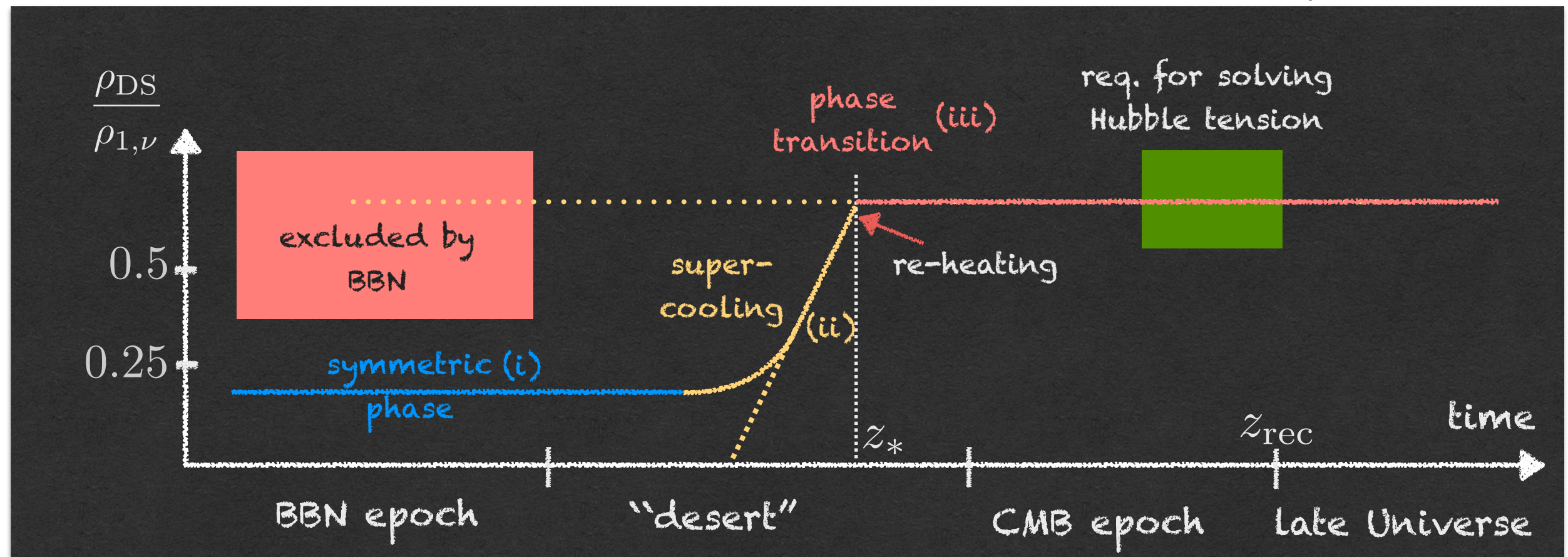
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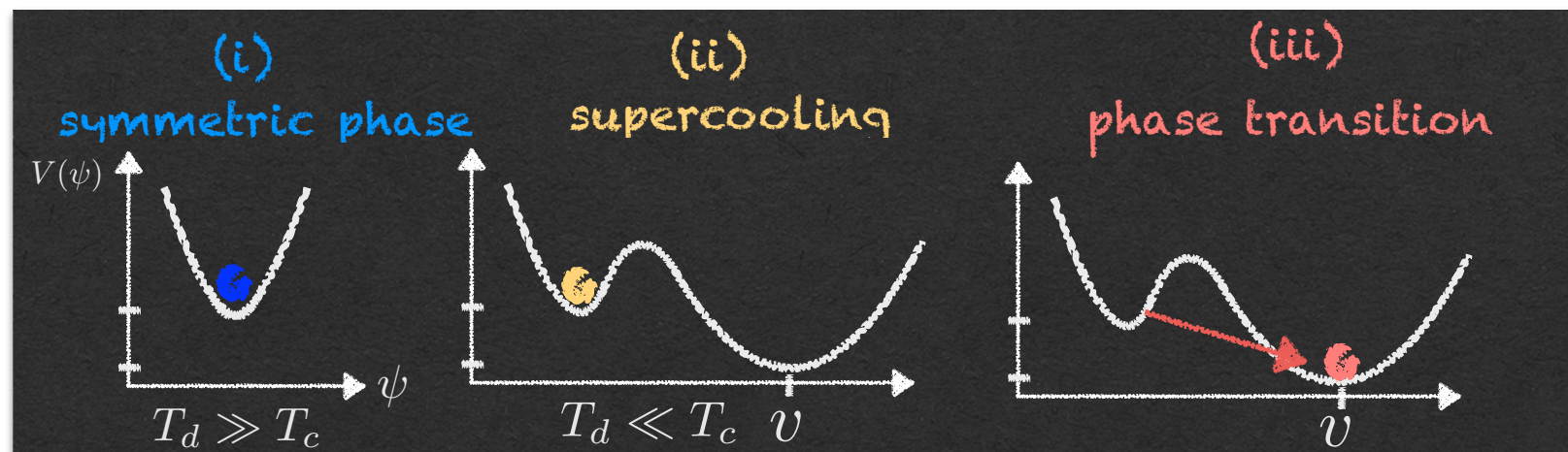
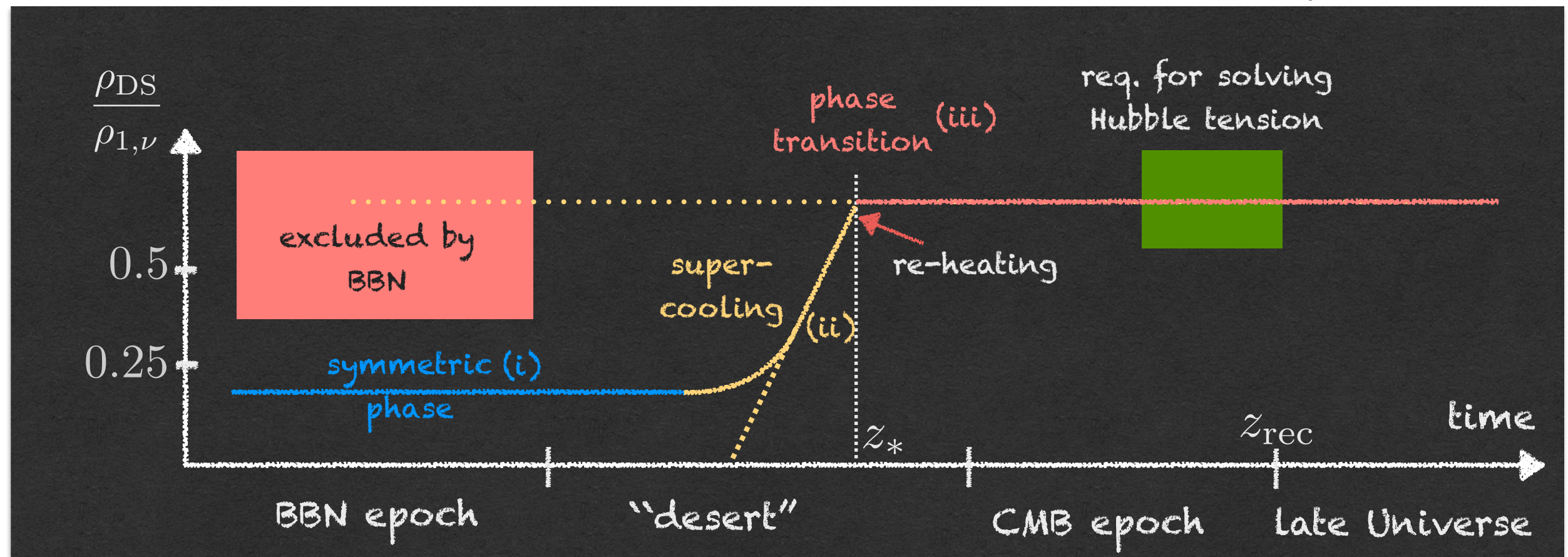
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Central simplifying assumptions:

- (i) $p(t) \propto \exp(\beta(t - t_*))$ with $\beta \gg H_* = 1/(2t_*)$ ↓
bubbles remain small
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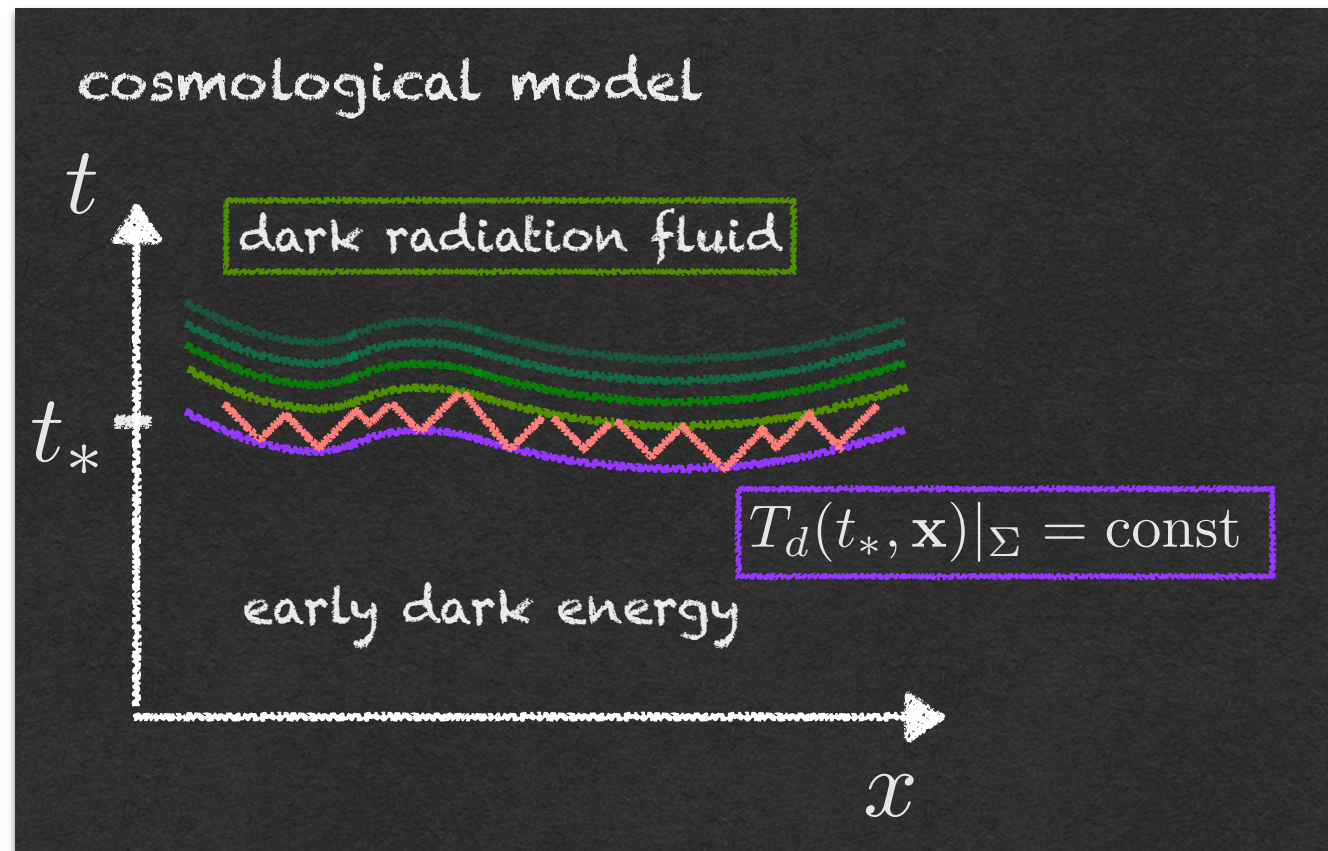
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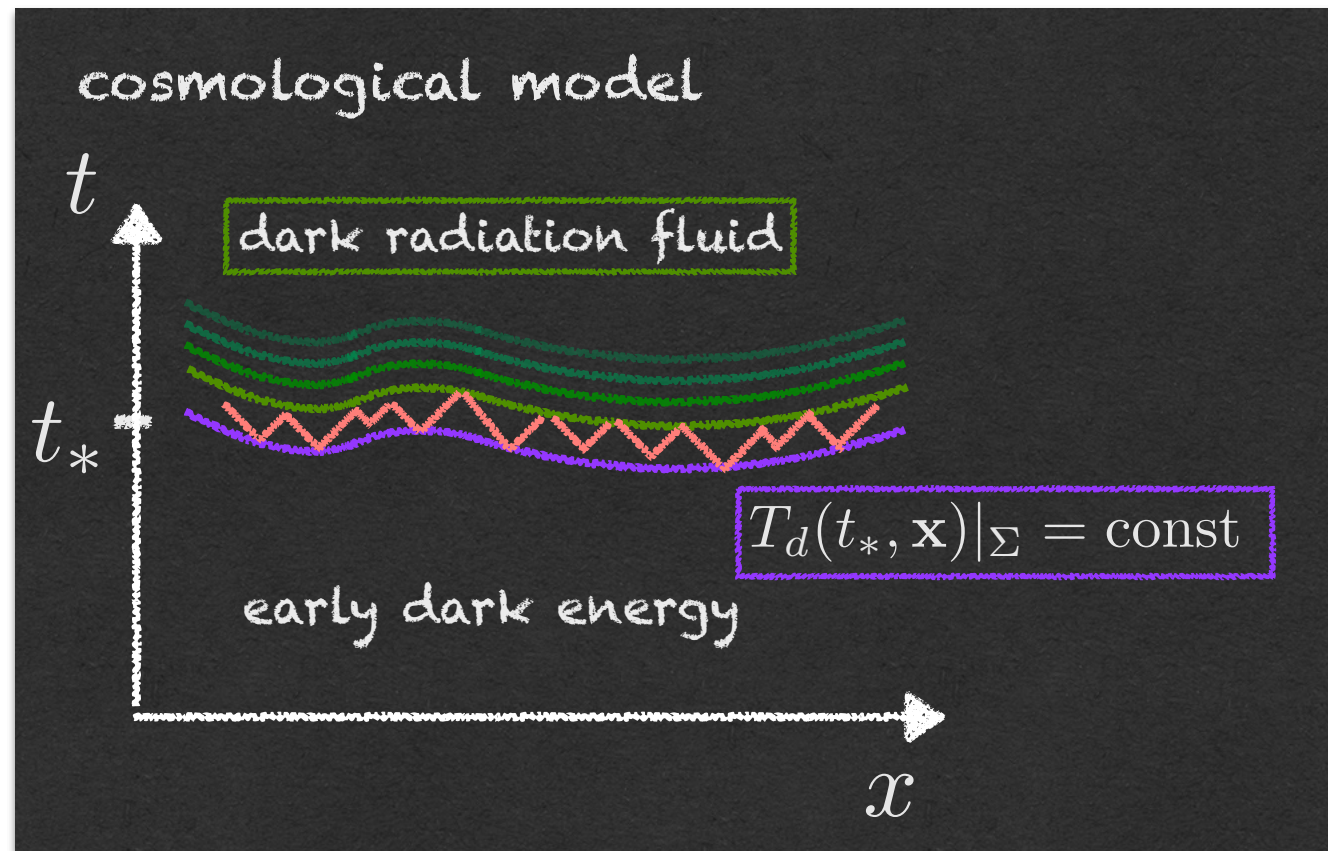
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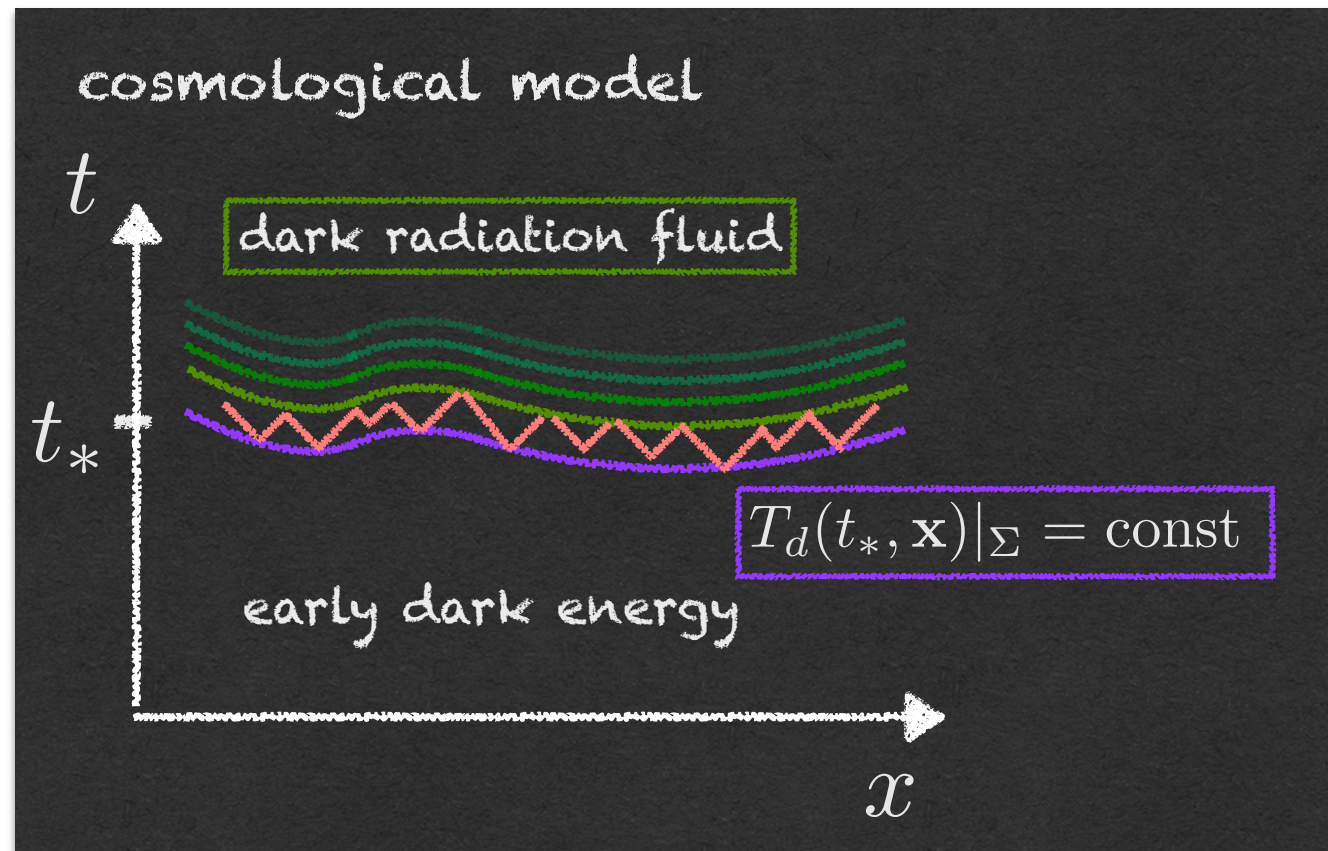
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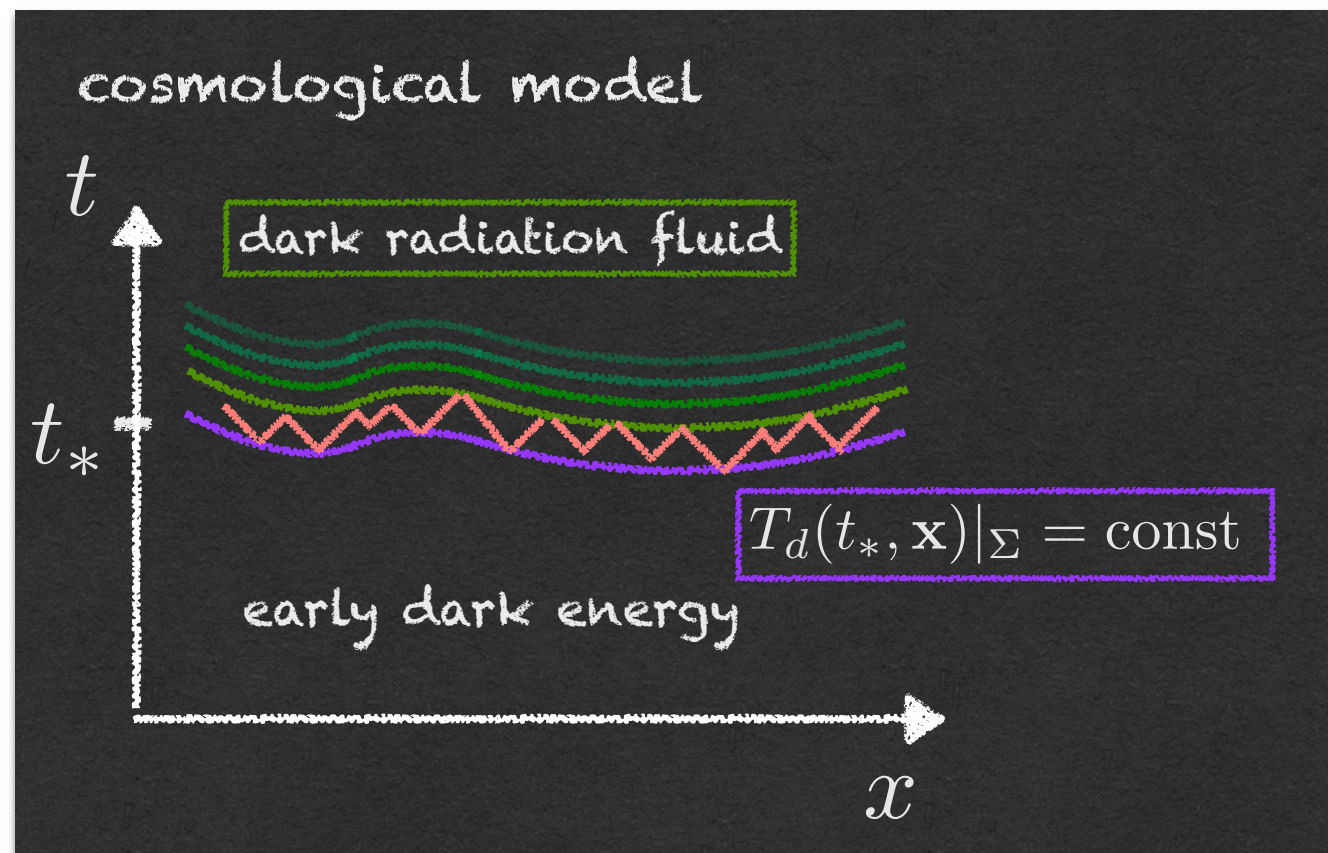
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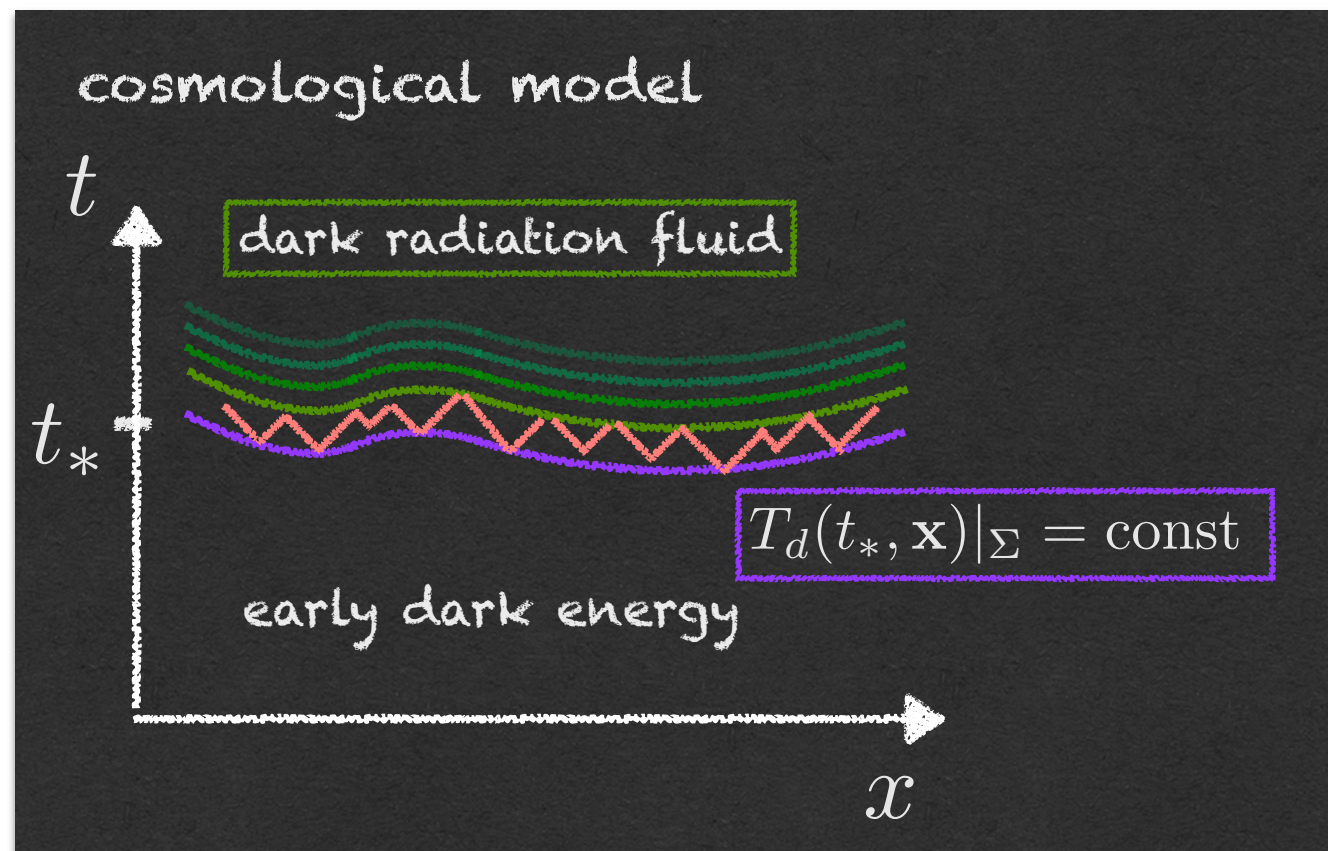
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2. Adiabatic fluctuations in dark sector temperature seed perturbations in post-p.t. fluid.

→ relevant for CMB + LSS

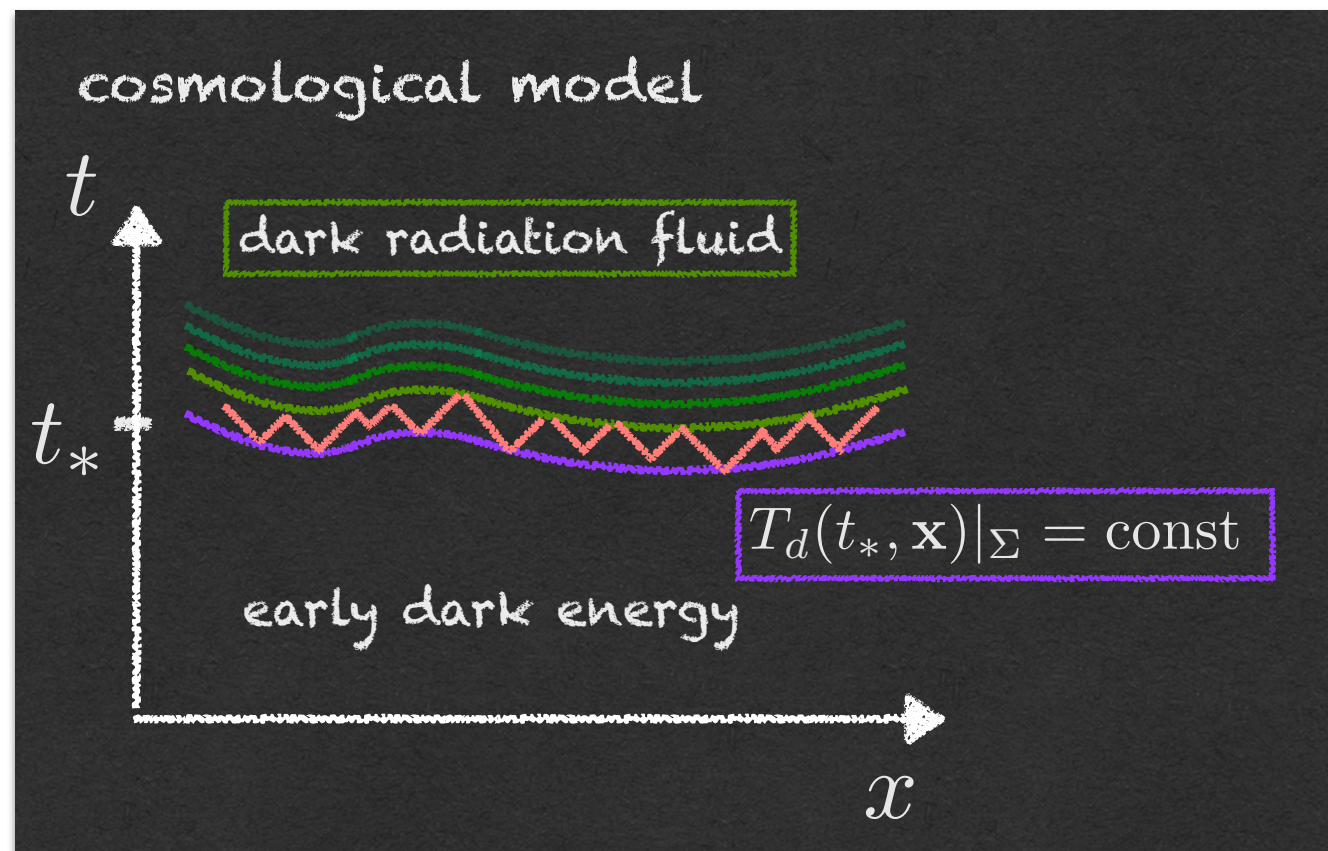
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► First results (triggered stepped SIDR):

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◆ **Residual tension: 2.8 sigma**

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Cosmological Perturbation Theory

recent bound:
 $\beta > 320 H_*$
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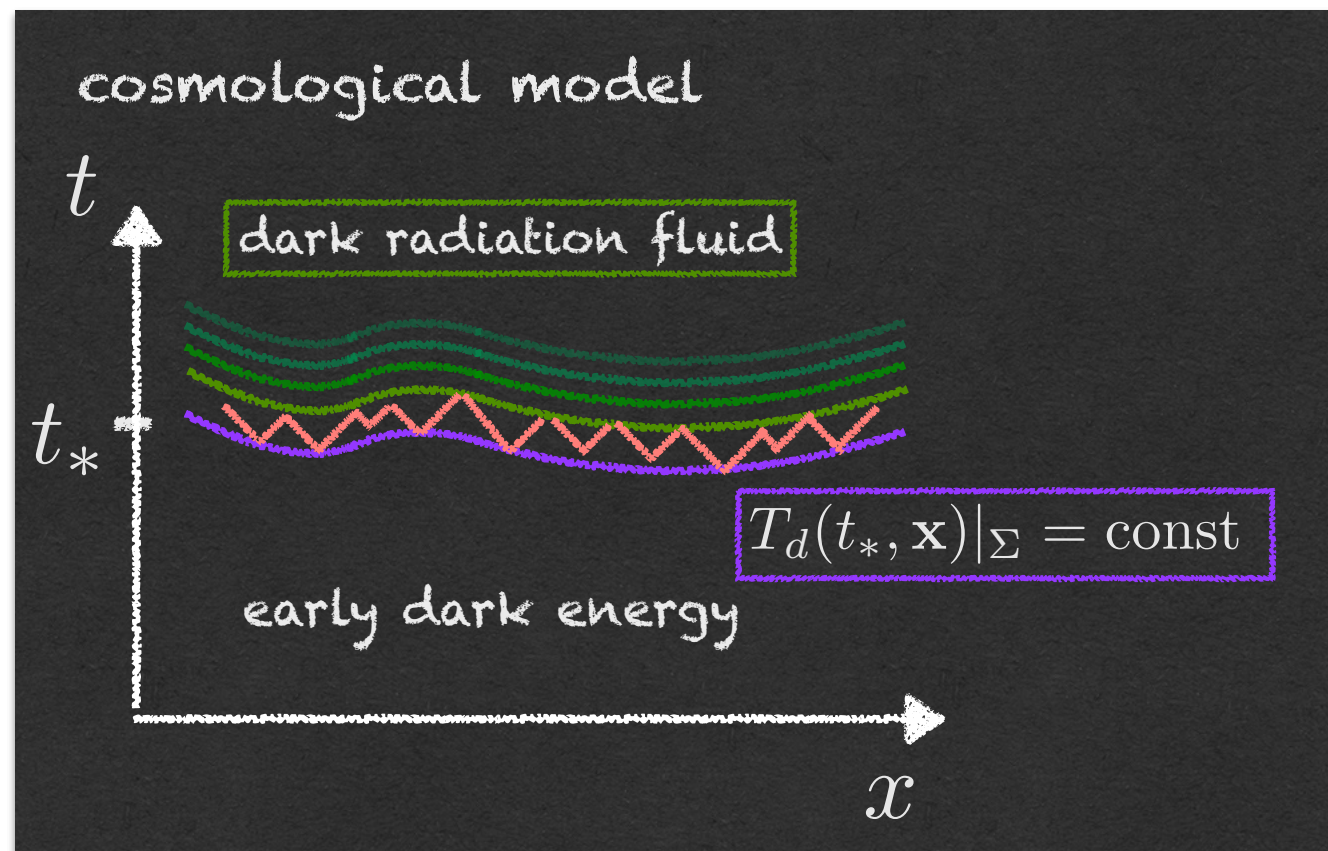
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► Work in progress [...] Stay tuned!

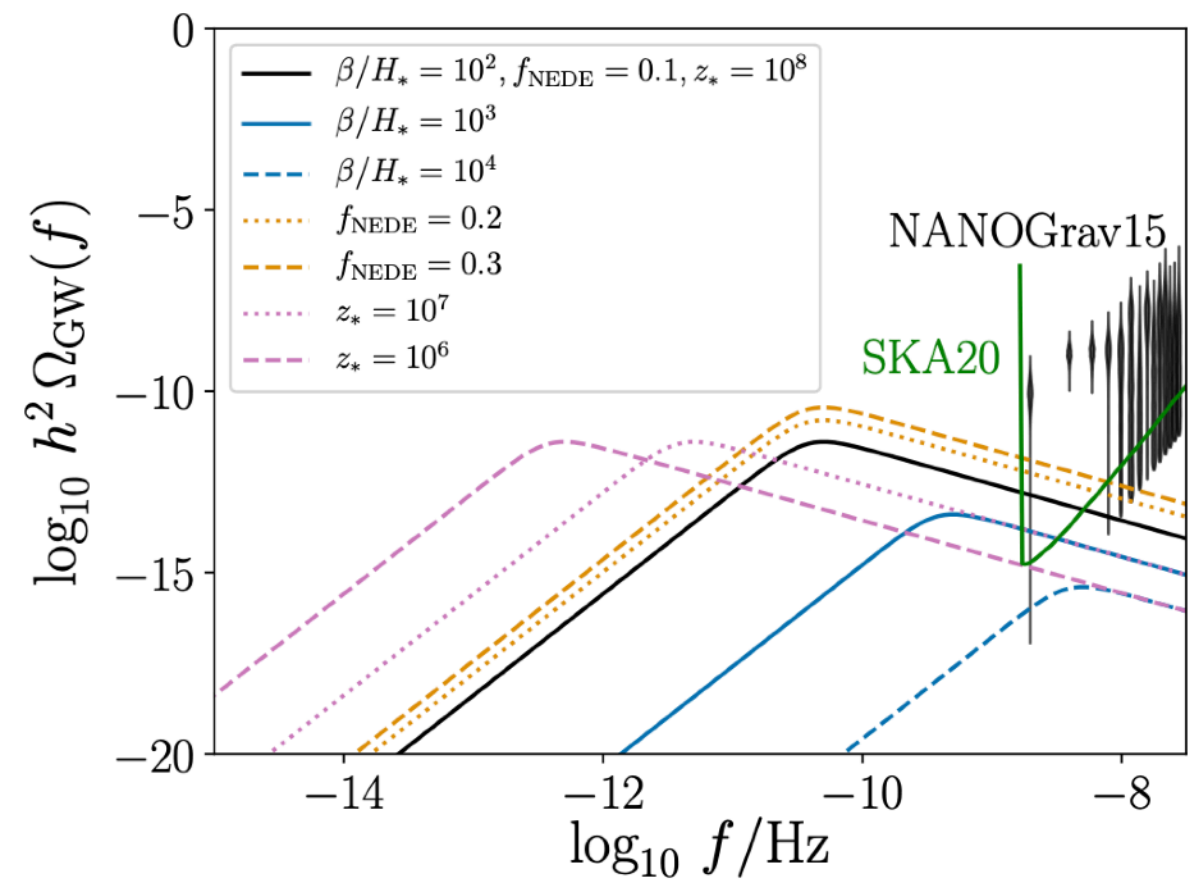
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 - ◆ Overlap with PTAs for phase transitions that occur close to BBN and create large bubbles.

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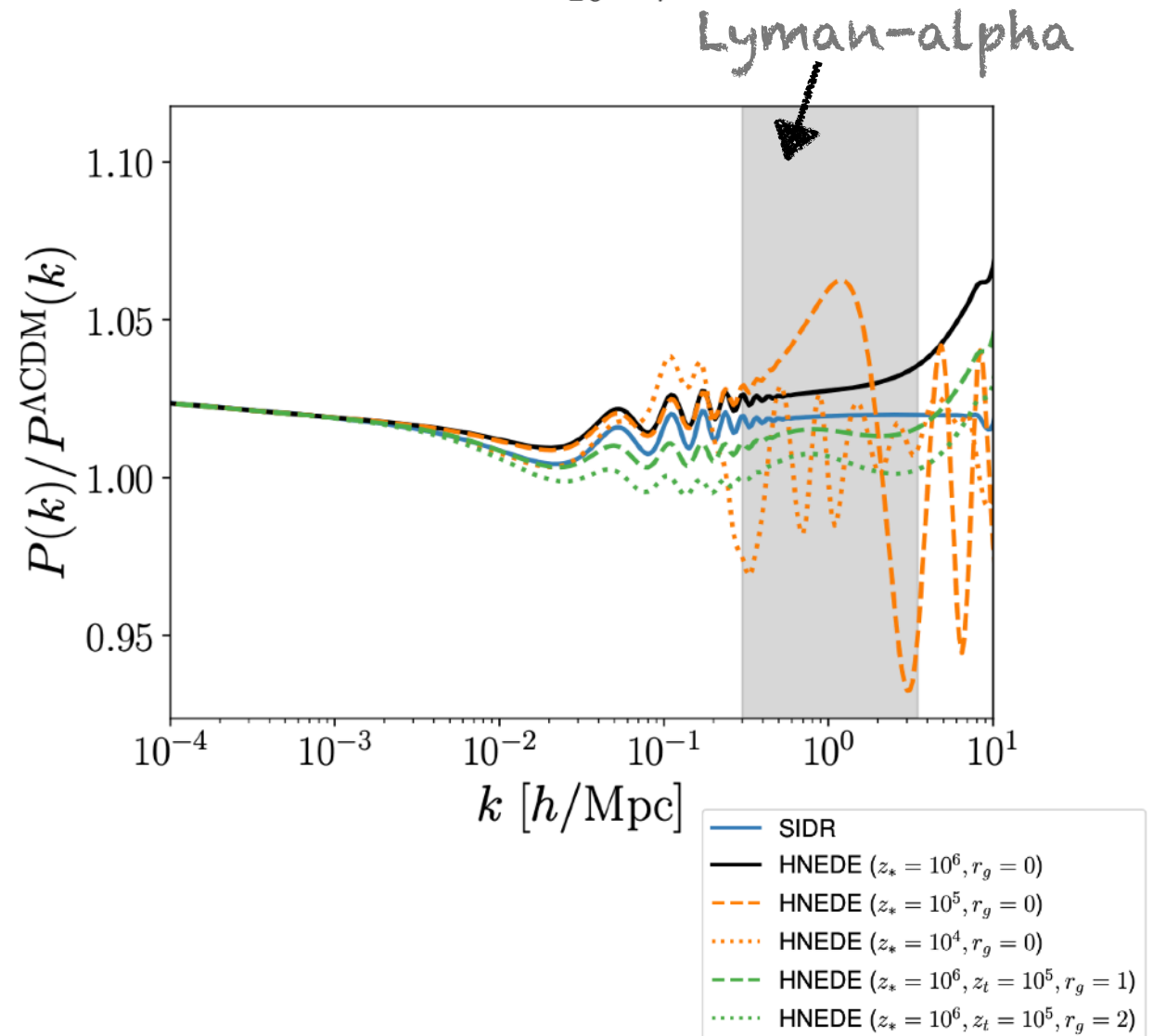
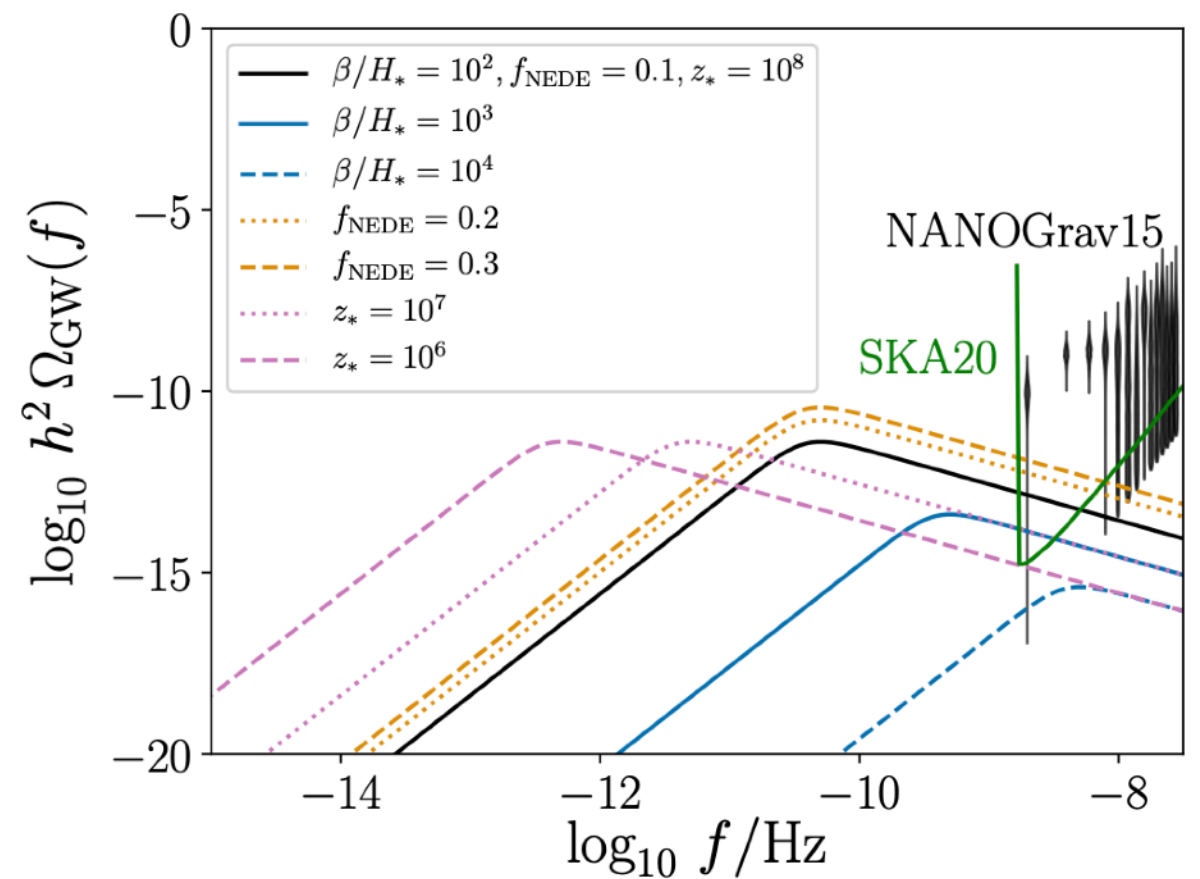
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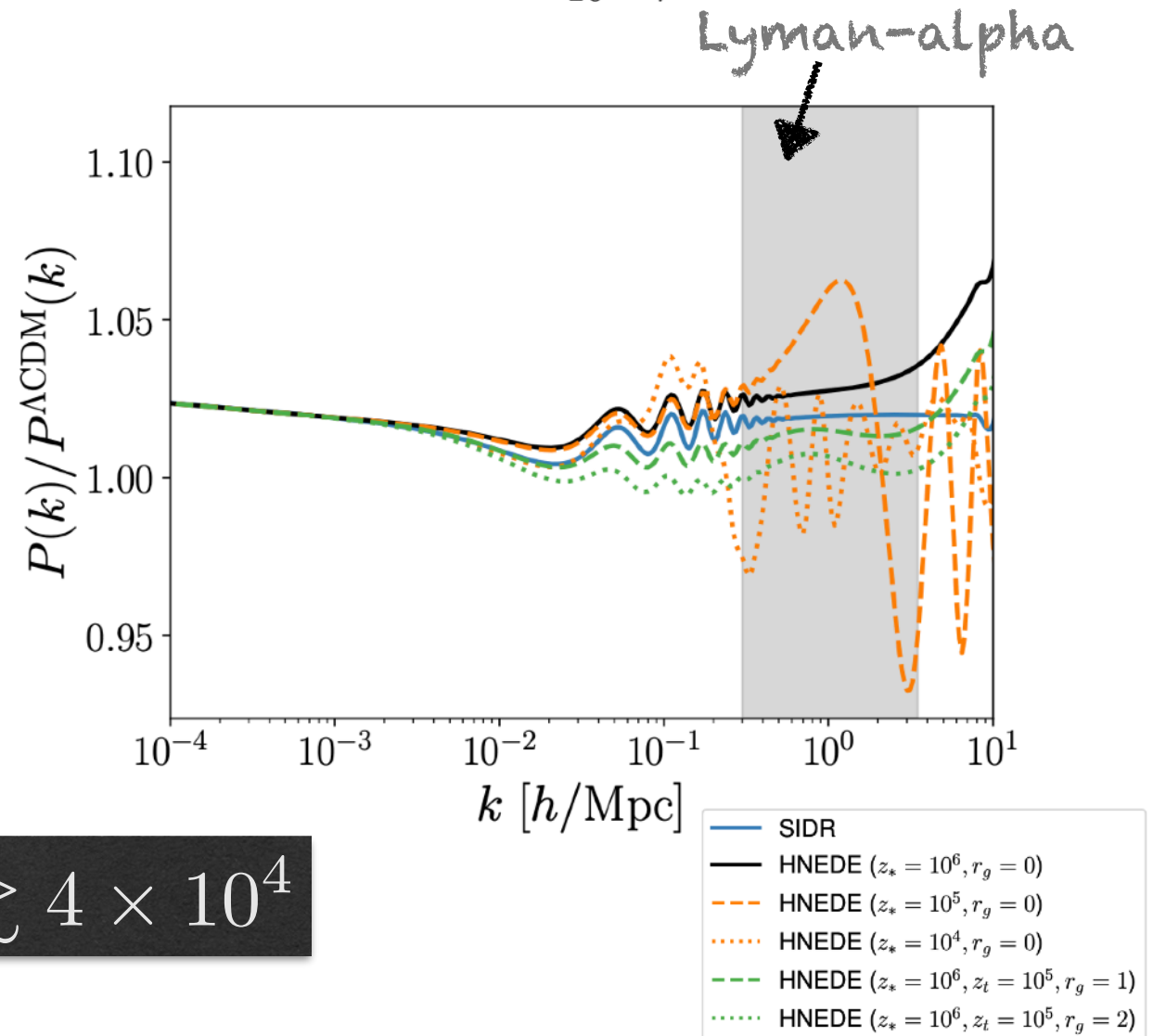
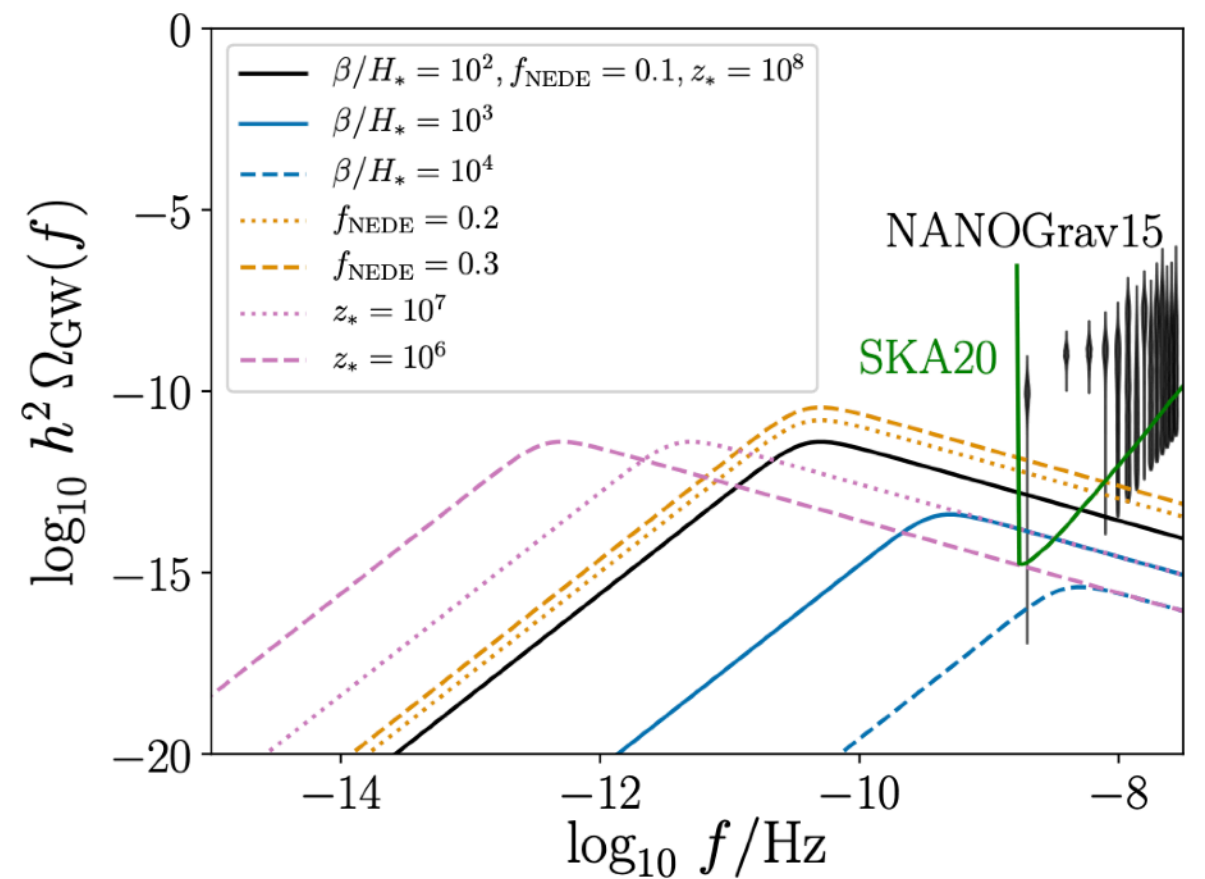
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► **CMB anisotropies**

- ◆ CMB provides lower bound on redshift
- ◆ Difference with cold NEDE due to character of trigger field.
- ◆ Acoustic oscillations in post p.t. fluid stronger in Hot NEDE.

$$z_* \gtrsim 4 \times 10^4$$



Summary

- ▶ The **Hubble tension** calls for new physics operative during the CMB epoch.
- ▶ Opportunity to probe new fundamental physics **above (but close to) the eV scale!**
- ▶ A strong first-order phase transition offers a simple microscopic scenario.
- ▶ **Cold New Early Dark Energy** relies on a triggered vacuum phase transition to bring the tension down to 2 sigma (challenges: describe post p.t. fluid, keep testing against new data).
- ▶ **Hot New Early Dark Energy** relies on a supercooled p.t. to produce DR **after BBN**.
- ▶ Simplest model: Brings tension below 3 sigma...  **Now: Exploit benefits of having microscopic scenario.**
- ▶ ... with **unique signatures** in matter power spectrum (+ PTAs).
- ▶ Take home:
 - Exciting times in cosmology as constraining power of cosmological probes is increasing.
 - Phase transitions are a simple playground for (early) dark energy / dark radiation physics.
 - Time to go beyond model-independent parametrizations and use **both** theory and data constraints.
 - **Invitation:** Many ideas wait to be explored!

Hot New Early Dark Energy

- Dark non-Abelian Higgs model with radiative breaking of conformal symmetry à la Coleman-Weinberg (CW)

$$\mathcal{L} = |D\Psi|^2 - V(\Psi) - \frac{1}{2}\text{tr}(F_{\mu\nu}F^{\mu\nu})$$

gauge coupling: $g \ll 1$

Ψ dark Higgs; transforms in fundamental

$$V(\psi; T_d) = V_0 + B\psi^4 \left(\ln \frac{\psi^2}{v^2} - \frac{1}{2} \right) - \frac{\mu_{\text{eff}}^2}{2} \psi^2 \left(1 - \frac{\psi^2}{2v^2} \right) + \Delta V_{\text{thermal}}(\psi; T_d),$$

$\psi = \sqrt{2}|\Psi|$
 $B \sim g^4$

CW 1-loop result

radiative symmetry breaking due to dim. transmutation

$$SU(N) \rightarrow SU(N-1) \rightarrow \psi \rightarrow v$$

(i) $\mu_{\text{eff}} \rightarrow 0$: $V'(0) = V''(0) = 0$

nucleation inhibited for $T_d > T_d^*|_{\text{CW}} (\ll v)$

strong supercooling

(ii) $\mu_{\text{eff}} > 0$: soft breaking of conformal symmetry

nucleation possible for $T_d \gg T_d^*|_{\text{CW}}$

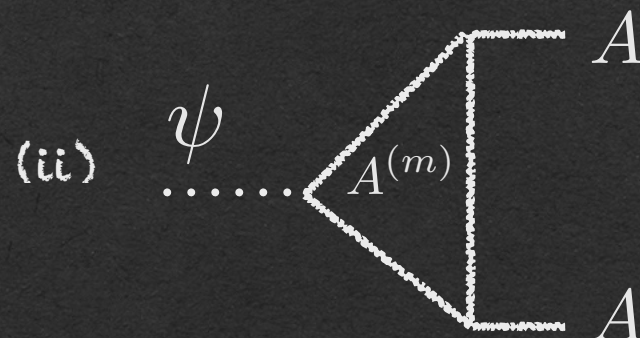
controlled supercooling

$$\mu_{\text{eff}}^2 \ll v^2 g^4$$

$$\frac{N_{\text{eff}}^{\text{after}} - N_{\text{eff}}^{\text{before}}}{N_{\text{eff}}^{\text{before}}} \propto 1/\mu_{\text{eff}}^4 \leftarrow \text{controls energy injection}$$

check assumptions:

(i) $\beta/H_* \propto g^{-2}$



efficient decay in massless gauge bosons

$$\frac{\Gamma_{\psi \rightarrow AA}^{(\text{cm})}}{H_*} = \mathcal{O}(1) \times \frac{g^9 f_{\text{NEDE}}^{1/4}}{1+z_*} 10^{24}$$

viable window $\rightarrow 0.1 \gtrsim g \gtrsim 0.01$